## Solving the KPZ equation

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## Introduction

Object of study: KPZ equation of surface growth:

$$\partial_t h = \partial_x^2 h + \lambda (\partial_x h)^2 + \xi - \infty$$

#### with either $x \in \mathbf{R}$ or $x \in S^1$ , and $\xi$ is space-time white noise.

1. Model for interface fluctuations (metastable  $\Rightarrow$  stable).

2. Free energy for polymer models.

3. Scaling limit of time-dependent parabolic Anderson model.

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- Expected to be universal for fluctuations of "weakly asymmetric" growing interfaces (see Quastel & Corwin's talks). Shown rigorously only for SOS/WASEP model, see Bertini-Giacomin 1997.
- 2. Interpolates between universality classes: Edwards-Wilkinson (Gaussian;  $t^{1/4}$ ) at small scales to KPZ (Tracy-Widom;  $t^{1/3}$ ) at large scales.
- 3. Interesting large time / space behaviour: fluctuation exponent  $t^{1/3}$ .
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## The SOS model

Simplest possible model of surface growth. Surface modelled by graph with slope  $\pm 1{:}$ 



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**Theorem (Bertini & Giacomin, 1997)**:  $\varepsilon^{\frac{1}{2}}h(\frac{x}{\varepsilon}, \frac{t}{\varepsilon^2}) - \varepsilon^{-1}$  converges to KPZ.

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## Cole-Hopf solution

#### Trick introduced by Cole and Hopf in the 50's. Write

$$h = \lambda^{-1} \log Z$$
 ,

#### then Z solves

$$\partial_t Z = \partial_x^2 Z + \lambda Z \,\xi \,. \tag{(\star)}$$

**Idea:** Take this as definition of solution, where (\*) is interpreted in the Itô sense. Work by Bertini-Giacomin shows that this is the physically relevant solution.

Write  $h = S_{CH}(h_0, \omega)$ , taking values in  $C(\mathbf{R}_+, C)$ .

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### Properties of Cole-Hopf

Mollify W, so  $W_{\varepsilon,k} = \varphi(\varepsilon k) W_k$  for cutoff  $\varphi$ , and set

$$dZ_{\varepsilon} = \partial_x^2 Z_{\varepsilon} \, dt + \lambda Z_{\varepsilon} \, dW_{\varepsilon}$$
,  $h_{\varepsilon} = \lambda^{-1} \, \log Z_{\varepsilon}$ .

Then  $h_{\varepsilon}$  solves

$$\partial_t h_{\varepsilon} = \partial_x^2 h_{\varepsilon} + \lambda \left( (\partial_x h_{\varepsilon})^2 - C_{\varepsilon} \right) + \xi_{\varepsilon} , \quad C_{\varepsilon} \approx \frac{1}{\varepsilon} \int \varphi^2 .$$

#### Problems with this notion of solution:

- 1. Not satisfactory at the formal level.
- 2. Lack of robustness: no good approximation theory for other modifications (hyperviscosity, time-smoothing, etc).
- 3. Properties of solutions do not always transform well (regularity of difference for example).

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Then  $h_{\varepsilon}$  solves

$$\partial_t h_{arepsilon} = \partial_x^2 h_{arepsilon} + \lambda ig( (\partial_x h_{arepsilon})^2 - C_{arepsilon} ig) + \xi_{arepsilon}$$
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- 1. Consider product as Wick product  $\partial_x h \diamond \partial_x h$  (Øksendal & Al 1995): wrong notion of solution ( $\neq S_{CH}$ ). Also wrong scaling properties (Chan 2000).
- 2. Formulate as martingale problem (Assing 2002): no well-posedness, "generator" not shown to be closable.
- Apply "standard" renormalisation techniques inspired by QFT (Da Prato, Debussche, Tubaro 2007): only works for a regularised equation.
- 4. Define nonlinearity on some distributional space (Gonçalves, Jara 2010, Assing 2011): potentially very powerful but no uniqueness as of now.

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#### A robustness result

**Theorem (H. 2011):** For  $\alpha > 0$  arbitrary one can build the following objects:

$$\begin{array}{c} \mathcal{X} \times \mathcal{C}^{\alpha} \xrightarrow{\mathcal{S}_{R}} \mathcal{C}(\mathbf{R}_{+}, \mathcal{C}^{\alpha}) \\ \Psi \uparrow \qquad \uparrow \qquad \downarrow \\ \Omega \ \times \mathcal{C}^{\alpha} \xrightarrow{\mathcal{S}_{CH}} \mathcal{C}(\mathbf{R}_{+}, \mathcal{C}^{\alpha}) \end{array}$$

where  $S_R$  is jointly continuous, but  $\Psi$  is only measurable. (Slight cheat: solutions are only local in general.)

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### A deterministic result

Consider solutions  $h_{\varepsilon}$  to

$$\partial_t h_\varepsilon = \partial_x^2 h_\varepsilon + (\partial_x h_\varepsilon)^2 + \varepsilon^{-3/2} g(\varepsilon^{-1} x - \varepsilon^{-2} t) - K_\varepsilon ,$$

for centred periodic g and suitable large constants  $K_{\varepsilon} = \frac{C_0}{\varepsilon} + \frac{C_1}{\sqrt{\varepsilon}} + \overline{K}$ . There is K such that  $h_{\varepsilon} \to h$  with

$$\partial_t h = \partial_x^2 h + (\partial_x h)^2 + K \partial_x h \; .$$

**Proof:** Just show that  $\Psi(g_{\varepsilon})$  converges to a limit in  $\mathcal{X}$ ... **Warning:** Two potential sources of logarithmic divergencies!

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### Ideas of technique

Idea: Perform Wild expansion of solution: define

$$\partial_t Y^{\bullet}_{\varepsilon} = \partial_x^2 Y^{\bullet}_{\varepsilon} + \xi_{\varepsilon}$$

For any binary tree  $au = [ au_1, au_2]$ , define  $Y_{arepsilon}^{ au}$  recursively by

$$\partial_t Y_{\varepsilon}^{\tau} = \partial_x^2 Y_{\varepsilon}^{\tau} + \partial_x Y_{\varepsilon}^{\tau_1} \, \partial_x Y_{\varepsilon}^{\tau_2} - C_{\varepsilon}^{\tau} \; .$$

Formal calculation shows that

$$h_arepsilon(t) = \sum_{ au} \lambda^{| au|-1} Y_arepsilon^{ au}(t)$$
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#### A convergence result

**Theorem:** For every  $\tau,$  there is a choice of  $\alpha_\tau$  and  $C_\varepsilon^\tau$  such that

$$Y_{\varepsilon}^{ au} o Y^{ au}$$
 (Independent of  $\varphi$ .)

in probability in  $\mathcal{C}(\mathbf{R}, \mathcal{C}^{\alpha}) \cap \mathcal{C}^{\beta}(\mathbf{R}, \mathcal{C})$  for  $\alpha < \alpha_{\tau}$  and  $\beta < \frac{1}{2}$ .

Optimal choice:  $\alpha_{\bullet} = \frac{1}{2}$ ,  $\alpha_{V} = 1$ ,  $\alpha_{\tau} = (\alpha_{\tau_{1}} \wedge \alpha_{\tau_{2}}) + 1$ .

$$\begin{split} C_{\varepsilon}^{\mathsf{V}} &= C_{\varepsilon} \sim \frac{1}{\varepsilon} \int_{\mathbf{R}} \varphi^2(x) \, dx \, ,\\ C_{\varepsilon}^{\mathsf{V}} &= \frac{4\pi}{\sqrt{3}} |\log \varepsilon| - C(\varphi) \, ,\\ C_{\varepsilon}^{\mathsf{V}} &= -\frac{1}{4} C_{\varepsilon}^{\mathsf{V}} \, . \end{split}$$

#### Truncated expansion

Idea: Write  $h_{\varepsilon}$  as

$$h_{\varepsilon} = \sum_{\tau \in \mathcal{T}} \lambda^{|\tau|-1} Y_{\varepsilon}^{\tau} + u_{\varepsilon}$$
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for a finite set  $\mathcal{T}$ , derive an equation for  $u_{\varepsilon}$ , and pass to limit. Minimal working choice:  $\mathcal{T} = \{\cdot, v, \forall, \forall\}$ . One obtains

$$\partial_t u_\varepsilon = \partial_x^2 u_\varepsilon + 2\lambda \, \partial_x u_\varepsilon \, \partial_x Y_\varepsilon^{\bullet} + \text{``I.o.t.''} \, .$$

Would like to make sense of

$$\partial_t u = \partial_x^2 u + 2\lambda \,\partial_x u \,\partial_x Y^{\bullet} \,.$$

"Theorem:" There exists no pair of Banach spaces containing u and Y such that the right-hand side makes sense. (Very different from DiPerna - Lions, closer to Flandoli - Russo.)

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#### How to solve that equation?

Writing  $v = \partial_x u$ , recall we want to solve

$$\partial_t v = \partial_x^2 v + 2\lambda \,\partial_x \big( v \,\partial_x Y^{\bullet} \big) \;.$$

If v were constant on the right hand side, then one would expect v to "look locally like"  $2\lambda v\Phi$ , where

$$\partial_t \Phi = \partial_x^2 \Phi + \partial_x^2 Y^{\bullet} \,.$$

**Idea:** Set up fixed point argument in space of functions that "look like  $\Phi$ " and use the fact that one can define  $\Phi \partial_x Y^{\bullet}$  "by hand".

Resulting space is a non-linear algebraic variety embedded in a larger Banach space. Uses controlled rough paths à la Gubinelli-Lyons.

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Take-away message: Nonlinear spaces are required to solve some rough equations pathwise.

- Convergence of microscopic models (for example lattice KPZ of Sasamoto & Spohn) to KPZ. See work with J. Maas and H. Weber.
- Extension to other equations in similar class (for example nonlinearities like  $g(h) (\partial_x h)^2$  or noises like  $f(h)\xi$ ).
- Encoding renormalisation procedures in higher dimensions?
- Deal with dynamically created singularities??

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