

Plateau flow  
alias  
The heat flow for half-harmonic maps

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## Half-harmonic maps

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$$E_{1/2}(u) = \frac{1}{2} \int_{S^1} |(-\Delta)^{1/4}u|^2 d\phi$$

of a map  $u \in H^{1/2}(S^1; N)$ , where

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Numerous results on regularity, “bubbling”, and quantization have been obtained by **Da Lio-Rivière**, **Da Lio-Martinazzi-Rivière**, **R. Moser**, **Schikorra**, and others, in particular, when  $N = S^k$ .

## Half-energy and Dirichlet energy

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of the half-Laplacian as the Dirichlet-to-Neumann map. Thus, the representation of the operators  $(-\Delta)^{1/2}$  and  $(-\Delta)^{1/4}$  in Fourier space with symbols  $|\xi|$ ,  $\sqrt{|\xi|}$ , and Parseval's identity give

$$\begin{aligned} \int_B |\nabla u|^2 dz &= \int_{S^1} u \partial_r u d\phi \\ &= \int_{S^1} u (-\Delta)^{1/2} u d\phi = \int_{S^1} |(-\Delta)^{1/4} u|^2 d\phi. \end{aligned}$$



## Connection with minimal surfaces

The characterization  $(-\Delta)^{1/2}u = \partial_r u$  of the half-Laplacian shows that a half-harmonic map  $u \in H^{1/2}(S^1; \mathbb{R}^n)$  induces a minimal surface with free boundary on  $N$  (Da Lio-Rivière (2011)).

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Indeed, the harmonic extension  $u \in H^1(B; \mathbb{R}^n)$  of  $u$  satisfies

$$d\pi(u)\partial_r u = d\pi(u)((-\Delta)^{1/2}u) = 0.$$

Thus,  $u_r \cdot u_\phi = 0$  on  $\partial B = S^1$ , and the (analytic) Hopf differential

$$\Phi = z^2 u_z^2 = r^2 |u_r|^2 - |u_\phi|^2 - 2iru_r \cdot u_\phi$$

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Moreover, from the condition  $d\pi(u)\partial_r u = 0$  we infer that  $u$  meets  $N$  vertically.

## The heat flow for half-harmonic maps

Wettstein (2021) studied the heat flow

$$d\pi(u)(u_t + (-\Delta)^{1/2}u) = 0 \quad \text{on } S^1 \times [0, \infty[, \quad (2)$$

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for a family of maps  $u = u(t) \in H^{1/2}(S^1; N)$  with harmonic extension  $u \in H^1(B; \mathbb{R}^n)$ . Observing that

$$\begin{aligned} \frac{d}{dt}E(u) &= \int_B \nabla u \nabla u_t \, dx = \int_{\partial B} u_r \cdot u_t \, d\phi \\ &= - \int_{\partial B} |d\pi(u)u_r|^2 \, d\phi = - \int_{\partial B} |u_t|^2 \, d\phi, \end{aligned}$$

we view (3) as heat flow for  $E$  with a free boundary constraint.

## The heat flow for half-harmonic maps

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**Theorem (S.(2021)).** Let  $N \subset \mathbb{R}^n$  be a closed, smooth sub-manifold of  $\mathbb{R}^n$ , and suppose that the normal bundle  $T^\perp N$  is parallelizable. For any  $u_0 \in H^{1/2}(S^1; N)$  there exists a unique global weak solution  $u = u(t)$  of (3) (hence (2)) with  $u(0) = u_0$ , whose energy is non-increasing and which is smooth away from finitely many points in space-time where non-trivial half-harmonic maps “bubble off”.

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## Remarks

- We exploit the surprising regularity properties of the normal component  $d\pi_N^\perp(u)\partial_r u$  of the harmonic extension of  $u$ , likely related to the commutator estimates for the normal projection in the work of **Da Lio-Rivière** and others.

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- A simple identity, and the use of the Dirichlet-to-Neumann map for the harmonic extension  $u: B \rightarrow \mathbb{R}^n$  of  $u$  instead of the half-Laplacian  $(-\Delta)^{1/2}u$  allow to perform the analysis completely avoiding fractional calculus.

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- Even though (3) appears to be degenerate, within our framework we are able to obtain similar smoothing properties as in the case of the harmonic map heat flow of surfaces.
- When  $N$  is a smoothly embedded, oriented closed curve  $\Gamma \subset \mathbb{R}^n$  the flow (3) may be viewed as an alternative gradient flow for the Plateau problem of disc-type minimal surfaces, the **Plateau flow**.

## Remarks on the Plateau flow

- It is not clear whether the **classical** Plateau boundary condition (requiring monotonicity) is preserved along the flow (3), even when  $\Gamma$  is a strictly convex planar curve. But also without monotonicity energy quantization holds (see below).



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- It should be straightforward to extend our results to the case when the disc  $B$  is replaced by a surface  $S$  of higher genus, if for given initial data  $u_0 \in H^{1/2}(S^1; N)$  we consider a family  $u = u(t)$  in  $H^{1/2}(S^1; N)$  solving the equation (3), that is,

$$u_t + d\pi(u)\partial_\nu u = 0$$

instead of (2), where  $\nu$  is the outward unit normal along  $\partial S = S^1$  and where for each time we harmonically extend  $u(t)$  to  $S$ ; see **Da Lio-Pigati** (2020) in the time-independent case.

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- In these cases, in order to obtain minimal surfaces one would also need to flow the conformal type of the domain, similar to **Rupflin-Topping** for harmonic map flow.

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$$\nu(u) \partial_r (\text{dist}_N(u)) = \nu(u) \nu(u) \cdot u_r =: d\pi_N^\perp(u) u_r$$

on  $\partial B = S^1$ , where  $d\pi_N^\perp(p) = 1 - d\pi_N(p)$  for each  $p \in N$  is the orthogonal projection to  $T_p^\perp N$ .



## A regularity estimate

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**Proposition.** There exists  $\delta > 0$  such that for any smooth solution  $u \in H^{1/2}(S^1; N)$  with harmonic extension  $u \in H^1(B, \mathbb{R}^n)$  of

$$d\pi_N(u)\partial_r u = f \quad \text{on } \partial B = S^1,$$

with  $E(u) \leq \delta^2$  there holds

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**Proof.** We have the orthogonal decomposition

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Extend  $\nu(p) := \nabla \text{dist}_N(p)$ . Note that  $\text{dist}_N(u) \in H_0^1(B)$  with

$$\|\nabla(\text{dist}_N(u))\|_{L^2(B)}^2 \leq C \|\nabla u\|_{L^2(B)}^2 \leq C\delta^2.$$

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Moreover,  $dist_N(u) \in H_0^1(B)$  satisfies the equation

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The basic  $L^2$ -theory for the Laplace equation and Sobolev's embedding  $H^{1/2}(B) \hookrightarrow L^4(B)$  then give the estimate

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Fourier expansion shows  $\|\nabla u\|_{H^{1/2}(B)}^2 = C \|\partial_r u\|_{L^2(S^1)}^2$ ;



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$$\begin{aligned} \|\partial_r(dist_N(u))\|_{L^2(S^1)}^2 &= \int_{\partial B} |z \cdot \nabla(dist_N(u))|^2 d\phi \\ &\leq C \|\nabla(dist_N(u))\|_{H^1(B)} \|\nabla(dist_N(u))\|_{L^2(B)} \leq C\delta \|\nabla u\|_{H^{1/2}(B)}^2. \end{aligned}$$

Fourier expansion shows  $\|\nabla u\|_{H^{1/2}(B)}^2 = C \|\partial_r u\|_{L^2(S^1)}^2$ ; thus

$$\begin{aligned} \|\partial_\phi u\|_{L^2(S^1)}^2 &= \|\partial_r u\|_{L^2(S^1)}^2 = \|f\|_{L^2(S^1)}^2 + \|\partial_r(dist_N(u))\|_{L^2(S^1)}^2 \\ &\leq \|f\|_{L^2(S^1)}^2 + C\delta \|\nabla u\|_{H^{1/2}(B)}^2 \leq \|f\|_{L^2(S^1)}^2 + C\delta \|\partial_r u\|_{L^2(S^1)}^2. \end{aligned}$$

## A threshold result

**Corollary.** (i) For any non-constant, smooth solution  $u \in H^{1/2}(S^1; N)$  with harmonic extension  $u \in H^1(B, \mathbb{R}^n)$  of

$$d\pi_N(u)\partial_r u = 0 \quad \text{on } \partial B = S^1, \quad (5)$$

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**Proof.** Let  $\gamma: S^1 \rightarrow \Gamma$  be an arc-length reference parametrization of  $\Gamma$ , and let  $u = \gamma \circ \xi$  on  $\partial B = S^1$  for some smooth  $\xi: S^1 \rightarrow S^1$ . Then at any  $p \in S^1$  where  $|\partial_\theta u| = |\partial_r u| \neq 0$  with some  $\lambda \neq 0$  we have

$$|d\pi_\Gamma(u)\partial_r u| = |\gamma'(\xi) \cdot \partial_r u| = \lambda |\partial_\theta u \cdot \partial_r u| = 0,$$

and  $u$  is conformal iff  $u$  solves (5). The claim then follows from (i).

## Non-concentration of energy gives regularity

Let  $u$  be a smooth solution of (3) on  $[0, T_0[$ .

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Hence the flow (3) can only become singular at time  $T_0$  if energy of size at least  $\delta > 0$  concentrates as  $t \uparrow T_0$ . (Energy concentration can then be analyzed via the usual rescaling procedure.)

## $H^2$ -bounds

**Lemma.** With a constant  $C > 0$  depending only on  $N$  there holds

$$\frac{d}{dt} \left( \int_{\partial B} |u_\phi|^2 d\phi \right) + \int_B |\nabla u_\phi|^2 dz \leq C \int_B |\nabla u|^2 |u_\phi|^2 dz.$$

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**Proof.** Writing  $d\pi_N(u) = 1 - \nu(u) \times \nu(u)$  we compute

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \left( \int_{\partial B} |u_\phi|^2 d\phi \right) &= \int_{\partial B} u_\phi \cdot u_{\phi,t} d\phi = - \int_{\partial B} u_{\phi\phi} \cdot u_t d\phi \\ &= \int_{\partial B} u_{\phi\phi} \cdot d\pi_N(u) u_r d\phi = - \int_{\partial B} (u_\phi \cdot u_{r\phi} - u_\phi \cdot \partial_\phi(\nu(u) \nu(u) \cdot u_r)) d\phi \\ &= -\frac{1}{2} \int_{\partial B} \partial_r(|u_\phi|^2) d\phi - \int_{\partial B} u_\phi \cdot d\nu(u) u_\phi \nu(u) \cdot u_r d\phi \\ &= - \int_B |\nabla u_\phi|^2 dz - \int_B \nabla u \cdot \nabla(\nu(u) u_\phi \cdot d\nu(u) u_\phi) dz. \end{aligned}$$

## Integral $H^2$ -bounds

Let  $u$  be a smooth solution of (3) on  $[0, T_0[$ , and for any  $\delta > 0$ , any  $T < T_0$  let  $R > 0$  such that (6) holds on  $[0, T]$ .

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From the previous lemma, with the multiplicative inequality

$$\int_B |\nabla u|^4 \varphi_{z_i, R}^2 dz \leq C\delta \int_{B_R(z_i)} |\nabla^2 u|^2 dz + C\delta R^{-2} \int_{B_R(z_i)} |\nabla u|^2 dz,$$

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**Proposition.** There exist constants  $\delta > 0$  and  $C > 0$  such that for any  $T < T_0$  there holds

$$\begin{aligned} \sup_{0 < t < T} \int_{\partial B} |u_\phi(t)|^2 d\phi + \int_0^T \int_B |\nabla u_\phi|^2 dx dt \\ \leq \int_{\partial B} |u_{0, \phi}|^2 d\phi + CTR^{-2} E(u_0), \end{aligned}$$

where  $R > 0$  is as in (6).

## Uniform $H^2$ - and higher bounds

With similar reasoning we also obtain a uniform  $H^2$ -bound.

**Proposition.** For any smooth  $u_0 \in H^{1/2}(S^1; N)$  and any  $T < T_0$  as specified above, with a constant  $C > 0$  possibly depending also on  $T$  there holds

$$\sup_{0 < t < T} \int_B |\nabla u_\phi(t)|^2 dx + \int_0^T \int_{\partial B} |u_{\phi r}|^2 d\phi dt \leq C.$$



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Higher bounds can then be obtained with two lemmas like the following.

**Lemma** For any  $k \geq 2$ , with a constant  $C > 0$  depending only on  $k$  and  $\Gamma$ , for the solution  $u = u(t)$  to (2) for any  $0 < t < T_0$  there holds

$$\begin{aligned} \frac{d}{dt} (\|\nabla \partial_\phi^k u\|_{L^2(B)}^2) + \|\partial_\phi^k u_r\|_{L^2(\partial B)}^2 \\ \leq C \sum_{1 \leq j_i \leq k+1, \sum j_i = k+2} \|\nabla \partial_\phi^k u\|_{L^2(B)} \|\Pi_i \nabla^{j_i} u\|_{L^2(B)}. \end{aligned}$$

## Local existence, blow-up and asymptotics

- Local existence can be obtained with a contraction-mapping argument for the map  $\Phi_\varepsilon: v \rightarrow u = \Phi_\varepsilon(v)$ , the solution of

$$u_t + (\varepsilon + d\pi(v))\partial_r u = 0, \quad u|_{t=0} = u_0,$$

using bounds similar to the previous lemma, which are uniform in  $0 < \varepsilon \ll 1$ .

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**Reference:** <http://arxiv.org/abs/2202.02083>

Thank you for your attention