

Compensated Integrability and Conservation Laws

Denis SERRE

UMPA, UMR 5669 CNRS



International PDE Conference 2022, Oxford, UK, July 20th–23rd

Plan of the talk

- What** are the objects under consideration?
Div-BV or Div-free symmetric tensors.
- Why** are they ubiquitous?
Because of Noether's Theorem.
- How** do we treat them?
With Compensated Integrability.
- Which** results do we get?
Dispersion (Strichartz-like) estimates.

WHAT? Div-free tensors

$U \subset \mathbb{R}^n$ an open domain.

Definition 1

A symmetric tensor over U is an $n \times n$ symmetric matrix A whose entries a_{jk} are distributions over $U \subset \mathbb{R}^n$.

WHAT? Div-free tensors

$U \subset \mathbb{R}^n$ an open domain.

Definition 1

A symmetric tensor over U is an $n \times n$ symmetric matrix A whose entries a_{jk} are distributions over $U \subset \mathbb{R}^n$.

Its (row-wise) Divergence is a vector of distributions :

$$(\text{Div } A)_j = \sum_{k=1}^n \partial_k a_{jk}.$$

WHAT? Div-free tensors

$U \subset \mathbb{R}^n$ an open domain.

Definition 1

A symmetric tensor over U is an $n \times n$ symmetric matrix A whose entries a_{jk} are distributions over $U \subset \mathbb{R}^n$.

Its (row-wise) Divergence is a vector of distributions :

$$(\text{Div } A)_j = \sum_{k=1}^n \partial_k a_{jk}.$$

We often assume positive semi-definiteness ; whence the distributions are Radon measures.

WHAT? Div-free tensors

$U \subset \mathbb{R}^n$ an open domain.

Definition 1

A symmetric tensor over U is an $n \times n$ symmetric matrix A whose entries a_{jk} are distributions over $U \subset \mathbb{R}^n$.

Its (row-wise) Divergence is a vector of distributions :

$$(\text{Div } A)_j = \sum_{k=1}^n \partial_k a_{jk}.$$

We often assume positive semi-definiteness ; whence the distributions are Radon measures.

Definition 2

The tensor A is Div-free if its entries are Radon measures, and $\text{Div } A \equiv 0$.

Two examples

1. Cf. R. Kupferman & A. Shachar : A geometric perspective on the Piola identity in Riemannian settings (arXiv 2018).

Two examples

Diagonal tensors ($U = I_1 \times \cdots \times I_n$).

-
1. Cf. R. Kupferman & A. Shachar : A geometric perspective on the Piola identity in Riemannian settings (arXiv 2018).

Two examples

Diagonal tensors ($U = I_1 \times \cdots \times I_n$).

Given n functions of $n - 1$ variables $f_j = f_j(\widehat{x}_j)$,

$$A := \text{diag}(f_1, \dots, f_n).$$

Since $\partial_j f_j = 0$, A is Div-free.

1. Cf. R. Kupferman & A. Shachar : A geometric perspective on the Piola identity in Riemannian settings (arXiv 2018).

Two examples

Diagonal tensors ($U = I_1 \times \cdots \times I_n$).

Given n functions of $n - 1$ variables $f_j = f_j(\widehat{x}_j)$,

$$A := \text{diag}(f_1, \dots, f_n).$$

Since $\partial_j f_j = 0$, A is Div-free.

Special tensors (from Piola's identity¹).

1. Cf. R. Kupferman & A. Shachar : A geometric perspective on the Piola identity in Riemannian settings (arXiv 2018).

Two examples

Diagonal tensors ($U = I_1 \times \cdots \times I_n$).

Given n functions of $n - 1$ variables $f_j = f_j(\widehat{x}_j)$,

$$A := \text{diag}(f_1, \dots, f_n).$$

Since $\partial_j f_j = 0$, A is Div-free.

Special tensors (from Piola's identity¹).

Given a potential $\theta : U \rightarrow \mathbb{R}$, the matrix of cofactors

$$A = \widehat{D^2 \theta}$$

is Div-free.

1. Cf. R. Kupferman & A. Shachar : A geometric perspective on the Piola identity in Riemannian settings (arXiv 2018).

Two examples

Diagonal tensors ($U = I_1 \times \cdots \times I_n$).

Given n functions of $n - 1$ variables $f_j = f_j(\widehat{x}_j)$,

$$A := \text{diag}(f_1, \dots, f_n).$$

Since $\partial_j f_j = 0$, A is Div-free.

Special tensors (from Piola's identity¹).

Given a potential $\theta : U \rightarrow \mathbb{R}$, the matrix of cofactors

$$A = \widehat{D^2 \theta}$$

is Div-free.

If $n = 2$, every Div-free tensor is special.

1. Cf. R. Kupferman & A. Shachar : A geometric perspective on the Piola identity in Riemannian settings (arXiv 2018).

Compensated Integrability will involve the total mass $\|\operatorname{Div} A\|_{\mathcal{M}}$.
Whence the

Definition 3

A tensor $A : \mathbb{R}^n \rightarrow \mathbf{Sym}_n$ is Div-BV if its entries a_{jk}

Compensated Integrability will involve the total mass $\|\operatorname{Div} A\|_{\mathcal{M}}$.
Whence the

Definition 3

*A tensor $A : \mathbb{R}^n \rightarrow \mathbf{Sym}_n$ is Div-BV if its entries a_{jk} **and** the coordinates $(\operatorname{Div} A)_j$ are finite measures.*

Compensated Integrability will involve the total mass $\|\operatorname{Div} A\|_{\mathcal{M}}$.
Whence the

Definition 3

A tensor $A : \mathbb{R}^n \rightarrow \mathbf{Sym}_n$ is Div-BV if its entries a_{jk} **and** the coordinates $(\operatorname{Div} A)_j$ are finite measures.

Definition 4

A tensor $A : U \rightarrow \mathbf{Sym}_n$ is Div-BV if its extension by 0_n to U^c is Div-BV over \mathbb{R}^n .

Compensated Integrability will involve the total mass $\|\operatorname{Div} A\|_{\mathcal{M}}$.
Whence the

Definition 3

A tensor $A : \mathbb{R}^n \rightarrow \mathbf{Sym}_n$ is Div-BV if its entries a_{jk} **and** the coordinates $(\operatorname{Div} A)_j$ are finite measures.

Definition 4

A tensor $A : U \rightarrow \mathbf{Sym}_n$ is Div-BV if its extension by 0_n to U^c is Div-BV over \mathbb{R}^n .

Equivalently, $A : U \rightarrow \mathbf{Sym}_n$ is Div-BV if its entries and the coordinates $(\operatorname{Div} A)_j$,

Compensated Integrability will involve the total mass $\|\text{Div } A\|_{\mathcal{M}}$.
Whence the

Definition 3

A tensor $A : \mathbb{R}^n \rightarrow \mathbf{Sym}_n$ is Div-BV if its entries a_{jk} **and** the coordinates $(\text{Div } A)_j$ are finite measures.

Definition 4

A tensor $A : U \rightarrow \mathbf{Sym}_n$ is Div-BV if its extension by 0_n to U^c is Div-BV over \mathbb{R}^n .

Equivalently, $A : U \rightarrow \mathbf{Sym}_n$ is Div-BV if its entries and the coordinates $(\text{Div } A)_j$, **and** its (well-defined) normal trace $A\vec{\nu}$ are finite measures.

Div-free/BV tensors are ubiquitous

Div-free/BV tensors are ubiquitous

First example : d -**dimensional Gas dynamics**.

Div-free/BV tensors are ubiquitous

First example : *d*-**dimensional Gas dynamics**.

The conservation of mass and momentum writes

$$\begin{aligned}\partial_t \rho + \operatorname{div}_y(\rho u) &= 0, \\ \partial_t(\rho u) + \operatorname{Div}_y(\rho u \otimes u) &= \operatorname{Div}_y \Sigma,\end{aligned}$$

where $\Sigma(t, y) \in \mathbf{Sym}_d$ is the stress tensor.

Div-free/BV tensors are ubiquitous

First example : d -**dimensional Gas dynamics**.

The conservation of mass and momentum writes

$$\begin{aligned}\partial_t \rho + \operatorname{div}_y(\rho u) &= 0, \\ \partial_t(\rho u) + \operatorname{Div}_y(\rho u \otimes u) &= \operatorname{Div}_y \Sigma,\end{aligned}$$

where $\Sigma(t, y) \in \mathbf{Sym}_d$ is the stress tensor.

This can be recast as $\operatorname{Div}_{t,y} A = 0$ with

$$A = \begin{pmatrix} \rho & \rho u^T \\ \rho u & \rho u \otimes u - \Sigma \end{pmatrix} \in \mathbf{Sym}_{1+d}.$$

Second example : **Relativistic GD.**

Second example : **Relativistic GD**.

Warning : the tensor involves the conservation law of energy, instead of that of the mass :

Second example : **Relativistic GD**.

Warning : the tensor involves the conservation law of energy, instead of that of the mass :

$$\partial_t \left(\frac{\rho c^2 + p}{c^2 - |v|^2} - \frac{p}{c^2} \right) + \operatorname{div}_y \left(\frac{\rho c^2 + p}{c^2 - |v|^2} v \right) = 0,$$

Second example : **Relativistic GD.**

Warning : the tensor involves the conservation law of energy, instead of that of the mass :

$$\begin{aligned} \partial_t \left(\frac{\rho c^2 + p}{c^2 - |v|^2} - \frac{p}{c^2} \right) + \operatorname{div}_y \left(\frac{\rho c^2 + p}{c^2 - |v|^2} v \right) &= 0, \\ \partial_t \left(\frac{\rho c^2 + p}{c^2 - |v|^2} v \right) + \operatorname{Div}_y \left(\frac{\rho c^2 + p}{c^2 - |v|^2} v \otimes v \right) + \nabla_y p &= 0. \end{aligned}$$

Second example : **Relativistic GD.**

Warning : the tensor involves the conservation law of energy, instead of that of the mass :

$$\begin{aligned}\partial_t \left(\frac{\rho c^2 + p}{c^2 - |v|^2} - \frac{p}{c^2} \right) + \operatorname{div}_y \left(\frac{\rho c^2 + p}{c^2 - |v|^2} v \right) &= 0, \\ \partial_t \left(\frac{\rho c^2 + p}{c^2 - |v|^2} v \right) + \operatorname{Div}_y \left(\frac{\rho c^2 + p}{c^2 - |v|^2} v \otimes v \right) + \nabla_y p &= 0.\end{aligned}$$

The (symmetric !) energy-momentum tensor

$$A = \begin{pmatrix} \frac{\rho c^2 + p}{c^2 - |v|^2} - \frac{p}{c^2} & \frac{\rho c^2 + p}{c^2 - |v|^2} v \\ \frac{\rho c^2 + p}{c^2 - |v|^2} v & \frac{\rho c^2 + p}{c^2 - |v|^2} v \otimes v + pI_3 \end{pmatrix}$$

is Div-free.

Why? Variational Principles!

Why? Variational Principles!

Many physical processes obey to a Variational Principle for a Lagrangian

$$\mathcal{L}[\alpha] = \int_U L(\alpha) dx,$$

α being a closed differential form :

$$d\alpha = 0.$$

Why? Variational Principles!

Many physical processes obey to a Variational Principle for a Lagrangian

$$\mathcal{L}[\alpha] = \int_U L(\alpha) dx,$$

α being a closed differential form :

$$d\alpha = 0.$$

Warning : the VP

$$\left. \frac{d}{d\epsilon} \right|_{\epsilon=0} \mathcal{L}[\alpha_\epsilon] = 0$$

Why? Variational Principles!

Many physical processes obey to a Variational Principle for a Lagrangian

$$\mathcal{L}[\alpha] = \int_U L(\alpha) dx,$$

α being a closed differential form :

$$d\alpha = 0.$$

Warning : the VP

$$\left. \frac{d}{d\epsilon} \right|_{\epsilon=0} \mathcal{L}[\alpha_\epsilon] = 0$$

involves admissible variations that are **not** additive ($\alpha_\epsilon \neq \alpha + \epsilon\beta$),

Why? Variational Principles!

Many physical processes obey to a Variational Principle for a Lagrangian

$$\mathcal{L}[\alpha] = \int_U L(\alpha) dx,$$

α being a closed differential form :

$$d\alpha = 0.$$

Warning : the VP

$$\left. \frac{d}{d\epsilon} \right|_{\epsilon=0} \mathcal{L}[\alpha_\epsilon] = 0$$

involves admissible variations that are **not** additive ($\alpha_\epsilon \neq \alpha + \epsilon\beta$), but **are** “compositional” :

$$\alpha_\epsilon = \underbrace{\phi_\epsilon^* \alpha}_{\text{pullback}}, \quad (\phi_\epsilon)_{\epsilon \in \mathbb{R}} \text{ a flow in } U.$$

Symmetry group :

$$L(M^* \alpha) = L(\alpha), \quad \forall \alpha, \quad \forall M \in G.$$

Symmetry group :

$$L(M^* \alpha) = L(\alpha), \quad \forall \alpha, \quad \forall M \in G.$$

Translations, either rotations or Lorentz transformations ...

Symmetry group :

$$L(M^* \alpha) = L(\alpha), \quad \forall \alpha, \quad \forall M \in G.$$

Translations, either rotations or Lorentz transformations ...

Theorem 1 (E. Noether 1918.)

The VP associates to every one-parameter group of symmetries, a conservation law.

Symmetry group :

$$L(M^* \alpha) = L(\alpha), \quad \forall \alpha, \quad \forall M \in G.$$

Translations, either rotations or Lorentz transformations ...

Theorem 1 (E. Noether 1918.)

The VP associates to every one-parameter group of symmetries, a conservation law.

Applications

Time invariance \longrightarrow conservation of energy,

Symmetry group :

$$L(M^* \alpha) = L(\alpha), \quad \forall \alpha, \quad \forall M \in G.$$

Translations, either rotations or Lorentz transformations ...

Theorem 1 (E. Noether 1918.)

The VP associates to every one-parameter group of symmetries, a conservation law.

Applications

Time invariance \longrightarrow conservation of energy,

Space invariance \longrightarrow conservation of momentum,

whence a Div-free tensor T .

Symmetry group :

$$L(M^* \alpha) = L(\alpha), \quad \forall \alpha, \quad \forall M \in G.$$

Translations, either rotations or Lorentz transformations ...

Theorem 1 (E. Noether 1918.)

The VP associates to every one-parameter group of symmetries, a conservation law.

Applications

Time invariance \longrightarrow conservation of energy,

Space invariance \longrightarrow conservation of momentum,

whence a Div-free tensor T .

Other invariances \longrightarrow If G contains $O(q)$ for a non-degenerate quadratic form $q(x) = x^T Sx$

Symmetry group :

$$L(M^* \alpha) = L(\alpha), \quad \forall \alpha, \quad \forall M \in G.$$

Translations, either rotations or Lorentz transformations ...

Theorem 1 (E. Noether 1918.)

The VP associates to every one-parameter group of symmetries, a conservation law.

Applications

Time invariance \rightarrow conservation of energy,

Space invariance \rightarrow conservation of momentum,

whence a Div-free tensor T .

Other invariances \rightarrow If G contains $O(q)$ for a non-degenerate quadratic form $q(x) = x^T Sx$ ($q(x) = |x|^2$ or $q(t, y) = c^2 t^2 - |y|^2$),

Symmetry group :

$$L(M^* \alpha) = L(\alpha), \quad \forall \alpha, \quad \forall M \in G.$$

Translations, either rotations or Lorentz transformations ...

Theorem 1 (E. Noether 1918.)

The VP associates to every one-parameter group of symmetries, a conservation law.

Applications

Time invariance \rightarrow conservation of energy,

Space invariance \rightarrow conservation of momentum,

whence a Div-free tensor T .

Other invariances \rightarrow If G contains $O(q)$ for a non-degenerate quadratic form $q(x) = x^T S x$ ($q(x) = |x|^2$ or $q(t, y) = c^2 t^2 - |y|^2$), then $A := S^{-1} T$ is **symmetric** and still Div-free.

Examples with Lorentz group

Examples with Lorentz group

- Relativistic GD (see above).

Examples with Lorentz group

- Relativistic GD (see above).
- (Maxwell system in vacuum) The **Electro-magnetic field**

Examples with Lorentz group

- Relativistic GD (see above).
- (Maxwell system in vacuum) The **Electro-magnetic field** is a 2-form

$$\alpha = (\vec{E} \cdot dy) \wedge dt + B_1 dy_2 \wedge dy_3 + B_2 dy_3 \wedge dy_1 + B_3 dy_1 \wedge dy_2.$$

Examples with Lorentz group

- Relativistic GD (see above).
- (Maxwell system in vacuum) The **Electro-magnetic field** is a 2-form

$$\alpha = (\vec{E} \cdot dy) \wedge dt + B_1 dy_2 \wedge dy_3 + B_2 dy_3 \wedge dy_1 + B_3 dy_1 \wedge dy_2.$$

Its closedness expresses the Gauß–Faraday law

$$\partial_t \vec{B} + \text{curl } \vec{E} = 0, \quad \text{div } \vec{B} = 0.$$

The Div-free energy-momentum tensor :

$$A = \begin{pmatrix} L - \vec{E} \cdot \vec{D} & \vec{H} \times \vec{E} \\ \vec{D} \times \vec{B} & (L + \vec{B} \cdot \vec{H}) I_3 - \vec{E} \otimes \vec{D} - \vec{H} \otimes \vec{B} \end{pmatrix},$$

where $\vec{D} = \frac{\partial L}{\partial \vec{E}}$, and $\vec{H} = -\frac{\partial L}{\partial \vec{B}}$ are the electric/magnetic inductions.

The Div-free energy-momentum tensor :

$$A = \begin{pmatrix} L - \vec{E} \cdot \vec{D} & \vec{H} \times \vec{E} \\ \vec{D} \times \vec{B} & (L + \vec{B} \cdot \vec{H}) I_3 - \vec{E} \otimes \vec{D} - \vec{H} \otimes \vec{B} \end{pmatrix},$$

where $\vec{D} = \frac{\partial L}{\partial \vec{E}}$, and $\vec{H} = -\frac{\partial L}{\partial \vec{B}}$ are the electric/magnetic inductions.

The symmetry requires the identity

$$\vec{H} \times \vec{E} = \vec{D} \times \vec{B},$$

The Div-free energy-momentum tensor :

$$A = \begin{pmatrix} L - \vec{E} \cdot \vec{D} & \vec{H} \times \vec{E} \\ \vec{D} \times \vec{B} & (L + \vec{B} \cdot \vec{H}) I_3 - \vec{E} \otimes \vec{D} - \vec{H} \otimes \vec{B} \end{pmatrix},$$

where $\vec{D} = \frac{\partial L}{\partial \vec{E}}$, and $\vec{H} = -\frac{\partial L}{\partial \vec{B}}$ are the electric/magnetic inductions.

The symmetry requires the identity

$$\vec{H} \times \vec{E} = \vec{D} \times \vec{B},$$

which is equivalent to Lorentz invariance :

$$L = L \left(\vec{E} \cdot \vec{B}, \frac{c^2 |\vec{B}|^2 - |\vec{E}|^2}{2} \right).$$

HOW do we treat Div-BV tensor?

HOW do we treat Div-BV tensor?

Div-BV mimics the space

$$BV(\mathbb{R}^n) = \{f \in \mathcal{M}(\mathbb{R}^n) \mid \nabla f \in \mathcal{M}(\mathbb{R}^n)\},$$

HOW do we treat Div-BV tensor?

Div-BV mimics the space

$$BV(\mathbb{R}^n) = \{f \in \mathcal{M}(\mathbb{R}^n) \mid \nabla f \in \mathcal{M}(\mathbb{R}^n)\},$$

for which we have

$$BV(\mathbb{R}^n) \subset L^{\frac{n}{n-1}}(\mathbb{R}^n),$$

HOW do we treat Div-BV tensor?

Div-BV mimics the space

$$BV(\mathbb{R}^n) = \{f \in \mathcal{M}(\mathbb{R}^n) \mid \nabla f \in \mathcal{M}(\mathbb{R}^n)\},$$

for which we have

$$BV(\mathbb{R}^n) \subset L^{\frac{n}{n-1}}(\mathbb{R}^n),$$

with a functional inequality (Gagliardo–Nirenberg–Sobolev)

$$\|f\|_{\frac{n}{n-1}} \leq c_n \|\nabla f\|_{\mathcal{M}}.$$

HOW do we treat Div-BV tensor?

Div-BV mimics the space

$$BV(\mathbb{R}^n) = \{f \in \mathcal{M}(\mathbb{R}^n) \mid \nabla f \in \mathcal{M}(\mathbb{R}^n)\},$$

for which we have

$$BV(\mathbb{R}^n) \subset L^{\frac{n}{n-1}}(\mathbb{R}^n),$$

with a functional inequality (Gagliardo–Nirenberg–Sobolev)

$$\|f\|_{\frac{n}{n-1}} \leq c_n \|\nabla f\|_{\mathcal{M}}.$$

But ∇ is elliptic, while Div is not !!

HOW do we treat Div-BV tensor?

Div-BV mimics the space

$$BV(\mathbb{R}^n) = \{f \in \mathcal{M}(\mathbb{R}^n) \mid \nabla f \in \mathcal{M}(\mathbb{R}^n)\},$$

for which we have

$$BV(\mathbb{R}^n) \subset L^{\frac{n}{n-1}}(\mathbb{R}^n),$$

with a functional inequality (Gagliardo–Nirenberg–Sobolev)

$$\|f\|_{\frac{n}{n-1}} \leq c_n \|\nabla f\|_{\mathcal{M}}.$$

But ∇ is elliptic, while Div is not !! In the spirit of Compensated Compactness, we expect that some non-linear quantity $D(A)$ behaves better than the entries a_{ij} do individually ...

Which quantity ?

Which quantity ?

The examination of either the **diagonal** (Gagliardo inequality), or the **special** cases in a periodic setting

Which quantity ?

The examination of either the **diagonal** (Gagliardo inequality), or the **special** cases in a periodic setting

$$\int (\det A)^{\frac{1}{n-1}} dx = \int \underbrace{\det D^2\theta}_{\text{null-Lagr.}} dx = \left(\det \int A dx \right)^{\frac{1}{n-1}}$$

Which quantity ?

The examination of either the **diagonal** (Gagliardo inequality), or the **special** cases in a periodic setting

$$\int (\det A)^{\frac{1}{n-1}} dx = \int \underbrace{\det D^2\theta}_{\text{null-Lagr.}} dx = \left(\det \int A dx \right)^{\frac{1}{n-1}}$$

suggests that this nice quantity is

$$A \xrightarrow{D} (\det A)^{\frac{1}{n-1}}.$$

Which quantity ?

The examination of either the **diagonal** (Gagliardo inequality), or the **special** cases in a periodic setting

$$\int (\det A)^{\frac{1}{n-1}} dx = \int \underbrace{\det D^2\theta}_{\text{null-Lagr.}} dx = \left(\det \int A dx \right)^{\frac{1}{n-1}}$$

suggests that this nice quantity is

$$A \xrightarrow{D} (\det A)^{\frac{1}{n-1}}.$$

This is where we need an extra assumption :

positive semi-definiteness : $A(x) \succ 0_n$.

Which quantity ?

The examination of either the **diagonal** (Gagliardo inequality), or the **special** cases in a periodic setting

$$\int (\det A)^{\frac{1}{n-1}} dx = \int \underbrace{\det D^2\theta}_{\text{null-Lagr.}} dx = \left(\det \int A dx \right)^{\frac{1}{n-1}}$$

suggests that this nice quantity is

$$A \xrightarrow{D} (\det A)^{\frac{1}{n-1}}.$$

This is where we need an extra assumption :

$$\underline{\text{positive semi-definiteness}} : A(x) \succ 0_n.$$

For special tensors, this means θ convex.

Main result : Compensated Integrability

Main result : Compensated Integrability

Observe that $(\det A)^{\frac{1}{n}}$ is a well-defined measure,

$$0_n \leq (\det A)^{\frac{1}{n}} \leq \frac{1}{n} \operatorname{Tr} A.$$

Main result : Compensated Integrability

Observe that $(\det A)^{\frac{1}{n}}$ is a well-defined measure,

$$0_n \leq (\det A)^{\frac{1}{n}} \leq \frac{1}{n} \operatorname{Tr} A.$$

Theorem 2 (Comp. Int. in \mathbb{R}^n (D.S., JMPA 2019).)

Let $A \succ 0_n$ be a Div-BV tensor over \mathbb{R}^n .

Main result : Compensated Integrability

Observe that $(\det A)^{\frac{1}{n}}$ is a well-defined measure,

$$0_n \leq (\det A)^{\frac{1}{n}} \leq \frac{1}{n} \operatorname{Tr} A.$$

Theorem 2 (Comp. Int. in \mathbb{R}^n (D.S., JMPA 2019).)

Let $A \succ 0_n$ be a Div-BV tensor over \mathbb{R}^n . Then $(\det A)^{\frac{1}{n}} \in L^{\frac{n}{n-1}}(\mathbb{R}^n)$

Main result : Compensated Integrability

Observe that $(\det A)^{\frac{1}{n}}$ is a well-defined measure,

$$0_n \leq (\det A)^{\frac{1}{n}} \leq \frac{1}{n} \operatorname{Tr} A.$$

Theorem 2 (Comp. Int. in \mathbb{R}^n (D.S., JMPA 2019).)

Let $A \succ 0_n$ be a Div-BV tensor over \mathbb{R}^n . Then $(\det A)^{\frac{1}{n}} \in L^{\frac{n}{n-1}}(\mathbb{R}^n)$ and we have

$$\int_{\mathbb{R}^n} (\det A)^{\frac{1}{n-1}} dx \leq c_n \|\operatorname{Div} A\|_{\mathcal{M}}.$$

Main result : Compensated Integrability

Observe that $(\det A)^{\frac{1}{n}}$ is a well-defined measure,

$$0_n \leq (\det A)^{\frac{1}{n}} \leq \frac{1}{n} \operatorname{Tr} A.$$

Theorem 2 (Comp. Int. in \mathbb{R}^n (D.S., JMPA 2019).)

Let $A \succ 0_n$ be a Div-BV tensor over \mathbb{R}^n . Then $(\det A)^{\frac{1}{n}} \in L^{\frac{n}{n-1}}(\mathbb{R}^n)$ and we have

$$\int_{\mathbb{R}^n} (\det A)^{\frac{1}{n-1}} dx \leq c_n \|\operatorname{Div} A\|_{\mathcal{M}}.$$

Dual structure : the “2nd” BVP for the Monge-Ampère equation

$$\det D^2 u = f \quad (> 0, u \text{ convex}).$$

Main result : Compensated Integrability

Observe that $(\det A)^{\frac{1}{n}}$ is a well-defined measure,

$$0_n \leq (\det A)^{\frac{1}{n}} \leq \frac{1}{n} \operatorname{Tr} A.$$

Theorem 2 (Comp. Int. in \mathbb{R}^n (D.S., JMPA 2019).)

Let $A \succ 0_n$ be a Div-BV tensor over \mathbb{R}^n . Then $(\det A)^{\frac{1}{n}} \in L^{\frac{n}{n-1}}(\mathbb{R}^n)$ and we have

$$\int_{\mathbb{R}^n} (\det A)^{\frac{1}{n-1}} dx \leq c_n \|\operatorname{Div} A\|_{\mathcal{M}}.$$

Dual structure : the “2nd” BVP for the Monge-Ampère equation

$$\det D^2 u = f \quad (> 0, u \text{ convex}).$$

The proof exploits Brenier’s theorem in Optimal Transport.

- The constant c_n is explicit and sharp! Equality happens when $A = \chi_B I_n$ and B is a ball.

- The constant c_n is explicit and sharp! Equality happens when $A = \chi_B I_n$ and B is a ball.
- For general domains Ω , the choice $A = \chi_\Omega I_n$ yields the Isoperimetric Inequality

$$\frac{\text{Vol}(\Omega)}{\text{Vol}(B_n)} \leq \left(\frac{\text{Area}(\partial\Omega)}{\text{Area}(\partial B_n)} \right)^{\frac{n}{n-1}}.$$

- The constant c_n is explicit and sharp! Equality happens when $A = \chi_B I_n$ and B is a ball.
- For general domains Ω , the choice $A = \chi_\Omega I_n$ yields the Isoperimetric Inequality

$$\frac{\text{Vol}(\Omega)}{\text{Vol}(B_n)} \leq \left(\frac{\text{Area}(\partial\Omega)}{\text{Area}(\partial B_n)} \right)^{\frac{n}{n-1}}.$$

- With $A = f(x)I_n$, one recovers the Sobolev embedding

$$BV(\mathbb{R}^n) \subset L^{\frac{n}{n-1}}(\mathbb{R}^n).$$

Theorem 3 (D.S., Ann. IHP 2018.)

Let $A \succ 0_n$ be a periodic Div-free tensor.

Theorem 3 (D.S., Ann. IHP 2018.)

Let $A \succ 0_n$ be a periodic Div-free tensor. Then

$$\int (\det A)^{\frac{1}{n-1}} dx \leq \left(\det \int A dx \right)^{\frac{1}{n-1}} .$$

Theorem 3 (D.S., Ann. IHP 2018.)

Let $A \succ 0_n$ be a periodic Div-free tensor. Then

$$\int (\det A)^{\frac{1}{n-1}} dx \leq \left(\det \int A dx \right)^{\frac{1}{n-1}}.$$

Looks like Jensen's Inequality ...

Theorem 3 (D.S., Ann. IHP 2018.)

Let $A \succ 0_n$ be a periodic Div-free tensor. Then

$$\int (\det A)^{\frac{1}{n-1}} dx \leq \left(\det \int A dx \right)^{\frac{1}{n-1}}.$$

Looks like Jensen's Inequality ...

but $\det^{\frac{1}{n-1}}$ is **not** concave over \mathbf{Sym}_n^+ .

This is Div-quasi-concavity (terminology of Fonseca, Müller, de Philippis)

This is Div-quasi-concavity (terminology of Fonseca, Müller, de Philippis) whence a weak-star upper semi-continuity result :

Theorem 4 (L. de Rosa & R. Tione, JFA 2020.)

Let $A_m \succ 0_n$ be a sequence of Div-BV tensors, such that $\text{Div } A_m$ is bounded in $\mathcal{M}(U)$ and $A_m \xrightarrow{} A$ in L^p with $p > \frac{n}{n-1}$. Then up to a subsequence*

$$* \lim_{m \rightarrow \infty} (\det A_m)^{\frac{1}{n-1}} \leq (\det A)^{\frac{1}{n-1}}.$$

This is Div-quasi-concavity (terminology of Fonseca, Müller, de Philippis) whence a weak-star upper semi-continuity result :

Theorem 4 (L. de Rosa & R. Tione, JFA 2020.)

Let $A_m \succ 0_n$ be a sequence of Div-BV tensors, such that $\text{Div } A_m$ is bounded in $\mathcal{M}(U)$ and $A_m \xrightarrow{} A$ in L^p with $p > \frac{n}{n-1}$. Then up to a subsequence*

$$* \lim_{m \rightarrow \infty} (\det A_m)^{\frac{1}{n-1}} \leq (\det A)^{\frac{1}{n-1}}.$$

See also

- Skipper & Wiedemann (2021),
- Guerra, Raiță & Schrecker (2021, 2022).

Evolution problems

Here $n = 1 + d$ and $x = (t, y)$.

Evolution problems

Here $n = 1 + d$ and $x = (t, y)$. A splits accordingly

$$A = \begin{pmatrix} \rho & m^T \\ m & \frac{1}{\rho} m \otimes m + \sigma \end{pmatrix}, \quad \det A = \rho \det \sigma.$$

Evolution problems

Here $n = 1 + d$ and $x = (t, y)$. A splits accordingly

$$A = \begin{pmatrix} \rho & m^T \\ m & \frac{1}{\rho} m \otimes m + \sigma \end{pmatrix}, \quad \det A = \rho \det \sigma.$$

The positiveness of A amounts to that of ρ and σ .

Evolution problems

Here $n = 1 + d$ and $x = (t, y)$. A splits accordingly

$$A = \begin{pmatrix} \rho & m^T \\ m & \frac{1}{\rho} m \otimes m + \sigma \end{pmatrix}, \quad \det A = \rho \det \sigma.$$

The positiveness of A amounts to that of ρ and σ . Denote

$$M \equiv \int \rho(t, y) dy.$$

Evolution problems

Here $n = 1 + d$ and $x = (t, y)$. A splits accordingly

$$A = \begin{pmatrix} \rho & m^T \\ m & \frac{1}{\rho} m \otimes m + \sigma \end{pmatrix}, \quad \det A = \rho \det \sigma.$$

The positiveness of A amounts to that of ρ and σ . Denote

$$M \equiv \int \rho(t, y) dy.$$

Theorem 5 (D.S. 2021.)

Let $A \succ 0_n$ be a Div-free tensor over $(0, T) \times \mathbb{R}^d$. Then

$$\int_0^T dt \int_{\mathbb{R}^d} (\rho \det \sigma)^{\frac{1}{d}} dx \leq_d M^{\frac{1}{d}} (\|m(0, \cdot)\|_{\mathcal{M}} + \|m(T, \cdot)\|_{\mathcal{M}}).$$

WHICH results do we get ?

C.I. applies to models that involve a **positive** Div-BV (often Div-free) tensor :

WHICH results do we get ?

C.I. applies to models that involve a **positive** Div-BV (often Div-free) tensor :

- Compressible Euler (the pressure being ≥ 0),

WHICH results do we get ?

C.I. applies to models that involve a **positive** Div-BV (often Div-free) tensor :

- Compressible Euler (the pressure being ≥ 0),
- Boltzmann equation,

WHICH results do we get ?

C.I. applies to models that involve a **positive** Div-BV (often Div-free) tensor :

- Compressible Euler (the pressure being ≥ 0),
- Boltzmann equation,
- Particle dynamics or mean field models, under a radial, repulsive interaction force,

WHICH results do we get ?

C.I. applies to models that involve a **positive** Div-BV (often Div-free) tensor :

- Compressible Euler (the pressure being ≥ 0),
- Boltzmann equation,
- Particle dynamics or mean field models, under a radial, repulsive interaction force,
- Hard spheres dynamics,

WHICH results do we get ?

C.I. applies to models that involve a **positive** Div-BV (often Div-free) tensor :

- Compressible Euler (the pressure being ≥ 0),
- Boltzmann equation,
- Particle dynamics or mean field models, under a radial, repulsive interaction force,
- Hard spheres dynamics,
- Multi-D scalar conservation laws.

WHICH results do we get ?

C.I. applies to models that involve a **positive** Div-BV (often Div-free) tensor :

- Compressible Euler (the pressure being ≥ 0),
- Boltzmann equation,
- Particle dynamics or mean field models, under a radial, repulsive interaction force,
- Hard spheres dynamics,
- Multi-D scalar conservation laws.

It does not (?) if the Div-free tensor is indefinite (or can be so) :

WHICH results do we get ?

C.I. applies to models that involve a **positive** Div-BV (often Div-free) tensor :

- Compressible Euler (the pressure being ≥ 0),
- Boltzmann equation,
- Particle dynamics or mean field models, under a radial, repulsive interaction force,
- Hard spheres dynamics,
- Multi-D scalar conservation laws.

It does not (?) if the Div-free tensor is indefinite (or can be so) :

- Navier-Stokes system,

WHICH results do we get ?

C.I. applies to models that involve a **positive** Div-BV (often Div-free) tensor :

- Compressible Euler (the pressure being ≥ 0),
- Boltzmann equation,
- Particle dynamics or mean field models, under a radial, repulsive interaction force,
- Hard spheres dynamics,
- Multi-D scalar conservation laws.

It does not (?) if the Div-free tensor is indefinite (or can be so) :

- Navier-Stokes system,
- Maxwell's equations,

WHICH results do we get ?

C.I. applies to models that involve a **positive** Div-BV (often Div-free) tensor :

- Compressible Euler (the pressure being ≥ 0),
- Boltzmann equation,
- Particle dynamics or mean field models, under a radial, repulsive interaction force,
- Hard spheres dynamics,
- Multi-D scalar conservation laws.

It does not (?) if the Div-free tensor is indefinite (or can be so) :

- Navier-Stokes system,
- Maxwell's equations,
- Attractive particle dynamics.

Estimate 1

Hypotheses : physical domain \mathbb{R}^d , finite mass M and initial energy E_0 .

Estimate 1

Hypotheses : physical domain \mathbb{R}^d , finite mass M and initial energy E_0 .

We have

$$\int_0^{+\infty} dt \int_{\mathbb{R}^d} \rho^{\frac{1}{d}} p \, dy \leq_d M^{\frac{1}{d}} \sqrt{ME_0} .$$

Estimate 1

*Hypotheses : physical domain \mathbb{R}^d , finite mass M and initial energy E_0 .
We have*

$$\int_0^{+\infty} dt \int_{\mathbb{R}^d} \rho^{\frac{1}{d}} p \, dy \leq_d M^{\frac{1}{d}} \sqrt{ME_0} .$$

A refinement of C.I., involving the action of the *projective group*, yields an improved dispersion :

Estimate 2 (mono-atomic gas.)

Suppose $p = \frac{2}{d} \rho e$ or $p = \rho^{1+\frac{2}{d}}$.

Estimate 1

*Hypotheses : physical domain \mathbb{R}^d , finite mass M and initial energy E_0 .
We have*

$$\int_0^{+\infty} dt \int_{\mathbb{R}^d} \rho^{\frac{1}{d}} p \, dy \leq_d M^{\frac{1}{d}} \sqrt{ME_0} .$$

A refinement of C.I., involving the action of the *projective group*, yields an improved dispersion :

Estimate 2 (mono-atomic gas.)

Suppose $p = \frac{2}{d} \rho e$ or $p = \rho^{1+\frac{2}{d}}$. Denoting $I_0 = \int \rho(0, x) \frac{|x|^2}{2} \, dx$ the moment of inertia at initial time,

Estimate 1

*Hypotheses : physical domain \mathbb{R}^d , finite mass M and initial energy E_0 .
We have*

$$\int_0^{+\infty} dt \int_{\mathbb{R}^d} \rho^{\frac{1}{d}} p \, dy \leq_d M^{\frac{1}{d}} \sqrt{ME_0} .$$

A refinement of C.I., involving the action of the *projective group*, yields an improved dispersion :

Estimate 2 (mono-atomic gas.)

Suppose $p = \frac{2}{d} \rho e$ or $p = \rho^{1+\frac{2}{d}}$. Denoting $I_0 = \int \rho(0, x) \frac{|x|^2}{2} \, dx$ the moment of inertia at initial time, we have

$$\int_0^{+\infty} t \, dt \int_{\mathbb{R}^d} \rho^{\frac{1}{d}} p \, dy \leq_d M^{\frac{1}{d}} \sqrt{MI_0} .$$

The time-space integrals do not depend upon the choice of the Galilean frame.

The time-space integrals do not depend upon the choice of the Galilean frame.

The right-hand sides do ... Take the infimum !

The time-space integrals do not depend upon the choice of the Galilean frame.

The right-hand sides do ... Take the infimum !

This lets us replace

$$ME_0 \longmapsto \frac{1}{4} \iint_{\mathbb{R}^d \times \mathbb{R}^d} \rho_0(y) \rho_0(z) |u_0(z) - u_0(y)|^2 dz dy \\ + M \int_{\mathbb{R}^d} \rho_0 e_0 dy,$$

The time-space integrals do not depend upon the choice of the Galilean frame.

The right-hand sides do ... Take the infimum !

This lets us replace

$$ME_0 \longmapsto \frac{1}{4} \iint_{\mathbb{R}^d \times \mathbb{R}^d} \rho_0(y) \rho_0(z) |u_0(z) - u_0(y)|^2 dz dy \\ + M \int_{\mathbb{R}^d} \rho_0 e_0 dy,$$

and

$$MI_0 \longmapsto \frac{1}{4} \iint_{\mathbb{R}^d \times \mathbb{R}^d} \rho_0(y) \rho_0(z) |z - y|^2 dz dy.$$

- The estimates do not assume an entropy condition.

- The estimates do not assume an entropy condition. They involve however the decay of the mechanical energy $t \mapsto E(t)$.

- The estimates do not assume an entropy condition. They involve however the decay of the mechanical energy $t \mapsto E(t)$.
- Say that the gas is barotropic ($p(\rho) = \rho^\gamma$ for $\gamma > 1$). Then

$$\rho \in \underbrace{L_t^\infty(L_y^1)}_{\text{mass}}$$

- The estimates do not assume an entropy condition. They involve however the decay of the mechanical energy $t \mapsto E(t)$.
- Say that the gas is barotropic ($p(\rho) = \rho^\gamma$ for $\gamma > 1$). Then

$$\rho \in \underbrace{L_t^\infty(L_y^1)}_{\text{mass}} \cap \underbrace{L_t^\infty(L_y^\gamma)}_{\text{energy}}$$

- The estimates do not assume an entropy condition. They involve however the decay of the mechanical energy $t \mapsto E(t)$.
- Say that the gas is barotropic ($p(\rho) = \rho^\gamma$ for $\gamma > 1$). Then

$$\rho \in \underbrace{L_t^\infty(L_y^1)}_{\text{mass}} \cap \underbrace{L_t^\infty(L_y^\gamma)}_{\text{energy}} \cap \underbrace{L_{t,y}^{\gamma + \frac{1}{d}}}_{\text{C.I.}}.$$

- The estimates do not assume an entropy condition. They involve however the decay of the mechanical energy $t \mapsto E(t)$.
- Say that the gas is barotropic ($p(\rho) = \rho^\gamma$ for $\gamma > 1$). Then

$$\rho \in \underbrace{L_t^\infty(L_y^1)}_{\text{mass}} \cap \underbrace{L_t^\infty(L_y^\gamma)}_{\text{energy}} \cap \underbrace{L_{t,y}^{\gamma + \frac{1}{d}}}_{\text{C.I.}}.$$

The internal energy may not concentrate.

- The estimates do not assume an entropy condition. They involve however the decay of the mechanical energy $t \mapsto E(t)$.
- Say that the gas is barotropic ($p(\rho) = \rho^\gamma$ for $\gamma > 1$). Then

$$\rho \in \underbrace{L_t^\infty(L_y^1)}_{\text{mass}} \cap \underbrace{L_t^\infty(L_y^\gamma)}_{\text{energy}} \cap \underbrace{L_{t,y}^{\gamma + \frac{1}{d}}}_{\text{C.I.}}.$$

The internal energy may not concentrate.

- We lack a companion estimate for the velocity.

- The estimates do not assume an entropy condition. They involve however the decay of the mechanical energy $t \mapsto E(t)$.
- Say that the gas is barotropic ($p(\rho) = \rho^\gamma$ for $\gamma > 1$). Then

$$\rho \in \underbrace{L_t^\infty(L_y^1)}_{\text{mass}} \cap \underbrace{L_t^\infty(L_y^\gamma)}_{\text{energy}} \cap \underbrace{L_{t,y}^{\gamma+\frac{1}{d}}}_{\text{C.I.}}.$$

The internal energy may not concentrate.

- We lack a companion estimate for the velocity.
- (Mono-atomic) Compare

$$\int_0^{+\infty} t \, dt \int_{\mathbb{R}^d} \rho^{\frac{1}{d}} p \, dy \leq_d M^{\frac{1}{d}} \sqrt{MI_0}$$

with (J.-Y. Chemin, 1990)

$$t^2 \int_{\mathbb{R}^d} p \, dy \leq \frac{2}{d} I_0.$$

Kinetic equations (Boltzman)

$f(t, y, \xi)$ the distribution of mass, ξ the velocity of particles.

Kinetic equations (Boltzman)

$f(t, y, \xi)$ the distribution of mass, ξ the velocity of particles.

Essentially the same estimates,

Kinetic equations (Boltzman)

$f(t, y, \xi)$ the distribution of mass, ξ the velocity of particles.

Essentially the same estimates, but $\rho^{\frac{1}{d}} p$ is replaced by $(\det \Xi)^{\frac{1}{d}}$ where

$$\Xi(t, x) = \int_{\mathbb{R}^d} f(t, y, \xi) \begin{pmatrix} 1 & \xi^T \\ \xi & \xi \otimes \xi \end{pmatrix} d\xi.$$

Kinetic equations (Boltzman)

$f(t, y, \xi)$ the distribution of mass, ξ the velocity of particles.

Essentially the same estimates, but $\rho^{\frac{1}{d}} p$ is replaced by $(\det \Xi)^{\frac{1}{d}}$ where

$$\Xi(t, x) = \int_{\mathbb{R}^d} f(t, y, \xi) \begin{pmatrix} 1 & \xi^T \\ \xi & \xi \otimes \xi \end{pmatrix} d\xi.$$

That is

$$\det \Xi = \frac{d!}{d+1} \int_{\mathbb{R}^d}^{\otimes(1+d)} f(\xi_0) \cdots f(\xi_d) V(\xi_0, \dots, \xi_d)^2 d\xi_0 \cdots d\xi_d,$$

Kinetic equations (Boltzman)

$f(t, y, \xi)$ the distribution of mass, ξ the velocity of particles.

Essentially the same estimates, but $\rho^{\frac{1}{d}} p$ is replaced by $(\det \Xi)^{\frac{1}{d}}$ where

$$\Xi(t, x) = \int_{\mathbb{R}^d} f(t, y, \xi) \begin{pmatrix} 1 & \xi^T \\ \xi & \xi \otimes \xi \end{pmatrix} d\xi.$$

That is

$$\det \Xi = \frac{d!}{d+1} \int_{\mathbb{R}^d}^{\otimes(1+d)} f(\xi_0) \cdots f(\xi_d) V(\xi_0, \dots, \xi_d)^2 d\xi_0 \cdots d\xi_d,$$

where $V(\xi_0, \dots, \xi_d)$ is the volume of the simplex whose vertices are ξ_0, \dots, ξ_d .

Kinetic equations (Boltzman)

$f(t, y, \xi)$ the distribution of mass, ξ the velocity of particles.

Essentially the same estimates, but $\rho^{\frac{1}{d}} p$ is replaced by $(\det \Xi)^{\frac{1}{d}}$ where

$$\Xi(t, x) = \int_{\mathbb{R}^d} f(t, y, \xi) \begin{pmatrix} 1 & \xi^T \\ \xi & \xi \otimes \xi \end{pmatrix} d\xi.$$

That is

$$\det \Xi = \frac{d!}{d+1} \int_{\mathbb{R}^d}^{\otimes(1+d)} f(\xi_0) \cdots f(\xi_d) V(\xi_0, \dots, \xi_d)^2 d\xi_0 \cdots d\xi_d,$$

where $V(\xi_0, \dots, \xi_d)$ is the volume of the simplex whose vertices are ξ_0, \dots, ξ_d .

1-D estimate known by J.-M. Bony (1987)

Kinetic equations (Boltzman)

$f(t, y, \xi)$ the distribution of mass, ξ the velocity of particles.

Essentially the same estimates, but $\rho^{\frac{1}{d}} p$ is replaced by $(\det \Xi)^{\frac{1}{d}}$ where

$$\Xi(t, x) = \int_{\mathbb{R}^d} f(t, y, \xi) \begin{pmatrix} 1 & \xi^T \\ \xi & \xi \otimes \xi \end{pmatrix} d\xi.$$

That is

$$\det \Xi = \frac{d!}{d+1} \int_{\mathbb{R}^d}^{\otimes(1+d)} f(\xi_0) \cdots f(\xi_d) V(\xi_0, \dots, \xi_d)^2 d\xi_0 \cdots d\xi_d,$$

where $V(\xi_0, \dots, \xi_d)$ is the volume of the simplex whose vertices are ξ_0, \dots, ξ_d .

1-D estimate known by J.-M. Bony (1987); used by C. Cercignani (2005) to prove that the DiPerna–Lions' renormalized solutions are distributional.

Hard spheres dynamics

Large number of spherical particles $B_\alpha(t)$, $\alpha \in \llbracket 1, N \rrbracket$.

Hard spheres dynamics

Large number of spherical particles $B_\alpha(t)$, $\alpha \in \llbracket 1, N \rrbracket$. Elastic collisions.

Hard spheres dynamics

Large number of spherical particles $B_\alpha(t)$, $\alpha \in \llbracket 1, N \rrbracket$. Elastic collisions.
Total mass $M = Nm$.

Initial data : positions/velocities. Yields conserved quantities

$$\text{energy} \quad E_0 = \frac{m}{2} \sum |u_\alpha(0)|^2,$$

Hard spheres dynamics

Large number of spherical particles $B_\alpha(t)$, $\alpha \in \llbracket 1, N \rrbracket$. Elastic collisions.
Total mass $M = Nm$.

Initial data : positions/velocities. Yields conserved quantities

energy	$E_0 = \frac{m}{2} \sum u_\alpha(0) ^2,$
standard deviation of velocity	$\bar{u}.$

Hard spheres dynamics

Large number of spherical particles $B_\alpha(t)$, $\alpha \in \llbracket 1, N \rrbracket$. Elastic collisions.
Total mass $M = Nm$.

Initial data : positions/velocities. Yields conserved quantities

energy	$E_0 = \frac{m}{2} \sum u_\alpha(0) ^2,$
standard deviation of velocity	$\bar{u}.$

Theorem 6 (R. K. Alexander 1975.)

Global existence, pairwise collisions only, for almost every initial data.

Hard spheres dynamics

Large number of spherical particles $B_\alpha(t)$, $\alpha \in \llbracket 1, N \rrbracket$. Elastic collisions.
Total mass $M = Nm$.

Initial data : positions/velocities. Yields conserved quantities

$$\begin{array}{ll} \text{energy} & E_0 = \frac{m}{2} \sum |u_\alpha(0)|^2, \\ \text{standard deviation of velocity} & \bar{u}. \end{array}$$

Theorem 6 (R. K. Alexander 1975.)

Global existence, pairwise collisions only, for almost every initial data.

Ya. Sinai's question :

Is the number K of collisions finite ? If so, how does it behave with N ?

Answers :

- **Yes** (Vaserstein 1979, Illner 1989),

Answers :

- **Yes** (Vaserstein 1979, Illner 1989),
- $\log K = O(N^2 \log N)$ (Burago & al. 1998),

Answers :

- **Yes** (Vaserstein 1979, Illner 1989),
- $\log K = O(N^2 \log N)$ (Burago & al. 1998),
- $\log K = O(N \log N)$ (Burdzy 2022),

Answers :

- **Yes** (Vaserstein 1979, Illner 1989),
- $\log K = O(N^2 \log N)$ (Burago & al. 1998),
- $\log K = O(N \log N)$ (Burdzy 2022),
- For some configuration, $\log K \sim \frac{N}{2} \log 2$ (Burago & Ivanov 2018).

Answers :

- **Yes** (Vaserstein 1979, Illner 1989),
- $\log K = O(N^2 \log N)$ (Burago & al. 1998),
- $\log K = O(N \log N)$ (Burdzy 2022),
- For some configuration, $\log K \sim \frac{N}{2} \log 2$ (Burago & Ivanov 2018).

The above estimates don't involve Functional Analysis. They provide huge, useless, upper bounds.

Answers :

- **Yes** (Vaserstein 1979, Illner 1989),
- $\log K = O(N^2 \log N)$ (Burago & al. 1998),
- $\log K = O(N \log N)$ (Burdzy 2022),
- For some configuration, $\log K \sim \frac{N}{2} \log 2$ (Burago & Ivanov 2018).

The above estimates don't involve Functional Analysis. They provide huge, useless, upper bounds.

Mind that K may be exponentially large !

Weighted estimate

1st step. Construct a Div-free tensor encoding the dynamics.

1st step. Construct a Div-free tensor encoding the dynamics.

Obstacle : The support is a graph \implies the rank is ≤ 1 a.e. Hence

$$(\det A)^{\frac{1}{n}} \equiv 0.$$

C. I. seems useless.

1st step. Construct a Div-free tensor encoding the dynamics.

Obstacle : The support is a graph \implies the rank is ≤ 1 a.e. Hence

$$(\det A)^{\frac{1}{n}} \equiv 0.$$

C. I. seems useless.

2nd step. Apply a modified version of C.I., adapted to singular supports :

1st step. Construct a Div-free tensor encoding the dynamics.

Obstacle : The support is a graph \implies the rank is ≤ 1 a.e. Hence

$$(\det A)^{\frac{1}{n}} \equiv 0.$$

C. I. seems useless.

2nd step. Apply a modified version of C.I., adapted to singular supports : $(\det A)^{\frac{1}{n-1}}$ is a set of Dirac masses at the nodes of the graph.

1st step. Construct a Div-free tensor encoding the dynamics.

Obstacle : The support is a graph \implies the rank is ≤ 1 a.e. Hence

$$(\det A)^{\frac{1}{n}} \equiv 0.$$

C. I. seems useless.

2nd step. Apply a modified version of C.I., adapted to singular supports : $(\det A)^{\frac{1}{n-1}}$ is a set of Dirac masses at the nodes of the graph.

Related to Minkowski's problem, solved by Pogorelov (1978).

Theorem 7 (D.S., ARMA 2021.)

Then

$$\sum_{\text{coll.}} \underbrace{|u_{\text{out}} - u_{\text{in}}|}_{\text{weight}} \leq_d N^2 \bar{u}.$$

Way better than N^N ...!

Theorem 7 (D.S., ARMA 2021.)

Then

$$\sum_{\text{coll.}} \underbrace{|u_{\text{out}} - u_{\text{in}}|}_{\text{weight}} \leq_d N^2 \bar{u}.$$

Way better than N^N ...!

In other words ($q_\alpha = mu_\alpha$ the linear momenta)

$$\text{mean} [TV(t \mapsto q_\alpha(t))] \leq_d \sqrt{ME_0}.$$

Entropy solutions of

$$\partial_t u + \operatorname{div}_y \vec{f}(u) = 0.$$

Entropy solutions of

$$\partial_t u + \operatorname{div}_y \vec{f}(u) = 0.$$

Analysis based upon the tensor $A = T \circ u$,

$$T(v) = \int_0^v \begin{pmatrix} 1 \\ \vec{f}'(s) \end{pmatrix} \otimes \begin{pmatrix} 1 \\ \vec{f}'(s) \end{pmatrix} d\mu(s),$$

Entropy solutions of

$$\partial_t u + \operatorname{div}_y \vec{f}(u) = 0.$$

Analysis based upon the tensor $A = T \circ u$,

$$T(v) = \int_0^v \begin{pmatrix} 1 \\ \vec{f}'(s) \end{pmatrix} \otimes \begin{pmatrix} 1 \\ \vec{f}'(s) \end{pmatrix} d\mu(s),$$

whose rows represent entropy-flux pairs.

Their divergence is controlled :

Entropy solutions of

$$\partial_t u + \operatorname{div}_y \vec{f}(u) = 0.$$

Analysis based upon the tensor $A = T \circ u$,

$$T(v) = \int_0^v \begin{pmatrix} 1 \\ \vec{f}'(s) \end{pmatrix} \otimes \begin{pmatrix} 1 \\ \vec{f}'(s) \end{pmatrix} d\mu(s),$$

whose rows represent entropy-flux pairs.

Their divergence is controlled : If $u_0, \vec{f} \circ u_0 \in L^1(\mathbb{R}^d)$, then

A is Div-BV.

Ex. : Burgers equation

$$\partial_t u + \partial_1 \frac{u^2}{2} + \cdots + \partial_d \frac{u^{d+1}}{d+1} = 0.$$

Ex. : Burgers equation

$$\partial_t u + \partial_1 \frac{u^2}{2} + \cdots + \partial_d \frac{u^{d+1}}{d+1} = 0.$$

The non-degeneracy implies a dispersion :

Ex. : Burgers equation

$$\partial_t u + \partial_1 \frac{u^2}{2} + \cdots + \partial_d \frac{u^{d+1}}{d+1} = 0.$$

The non-degeneracy implies a dispersion :

Theorem 8 (D. S. & L. Silvestre, ARMA 2019.)

There are exponents $\alpha(d, p), \beta(d, p) > 0$ such that the solution of the Cauchy problem with $u_0 \in L^1(\mathbb{R}^d)$ exists and satisfies for every $1 < p \leq +\infty$ and every $t > 0$

$$\|u(t)\|_p \leq_{d,p} \|u_0\|_1^\alpha t^{-\beta}.$$

Ex. : Burgers equation

$$\partial_t u + \partial_1 \frac{u^2}{2} + \cdots + \partial_d \frac{u^{d+1}}{d+1} = 0.$$

The non-degeneracy implies a dispersion :

Theorem 8 (D. S. & L. Silvestre, ARMA 2019.)

There are exponents $\alpha(d, p), \beta(d, p) > 0$ such that the solution of the Cauchy problem with $u_0 \in L^1(\mathbb{R}^d)$ exists and satisfies for every $1 < p \leq +\infty$ and every $t > 0$

$$\|u(t)\|_p \leq_{d,p} \|u_0\|_1^\alpha t^{-\beta}.$$

Corollary 1 (M. Crandall's question (1972).)

The PDE and the entropy inequalities are satisfied in the distributional sense.

Thank you for your attention !