

Decay estimates for large velocities in the Boltzmann equation without cut-off

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joint work with C. Mouhot (Cambridge) and L. Silvestre (Chicago)

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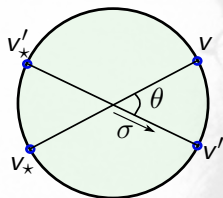


The Boltzmann equation

$$\underbrace{\partial_t f + v \cdot \nabla_x f}_{\text{free transport operator}} = \underbrace{Q(f, f)}_{\text{collision kernel}} = \int (f'_* f' - f f_*) B(r, \cos \theta) d v_* d \sigma$$

where $f'_* = f(v'_*)$, $f' = f(v')$, $f_* = f(v_*)$

Pre- and post-collisional velocities ($v, v_*/v', v'_*$)



$$\begin{cases} v + v_* = v' + v'_* \\ |v|^2 + |v_*|^2 = |v'|^2 + |v'_*|^2 \end{cases}$$

$$\begin{cases} r = |v - v_*| \\ \cos \theta := \frac{v - v_*}{|v - v_*|} \cdot \frac{v' - v'_*}{|v' - v'_*|} \end{cases}$$

Deviation angle θ

Cross section B : the way "particles" interact

$$Q(f, f) = \int (f'_* f' - f f_*) B(r, \cos \theta) \, d v_* \, d \sigma$$

$$B(r, \cos \theta) = r^\gamma b(\cos \theta) \quad \text{with} \quad -d < \gamma \leq 2$$

Cut-off $\int_{S^{d-1}} b(\cos \theta) \, d \sigma < \boxed{+\infty}$

Non cut-off $b(\cos \theta) \simeq \boxed{\frac{1}{|\theta|^{d-1+2s}}}$ with $s \in (0, 1)$

Moment estimates (cut-off)

Moment of order k : $\int_{\mathbb{R}^d} f(t, x, v) |v|^k dv$

{ classical & important theme in kinetic theory
first step in the theory of the Boltzmann equation
mostly for $\gamma > 0$ and for the homogeneous case

Moment estimates (cut-off)

Hard potentials ($\gamma > 0$) and homogeneous

- Povzner '62, Elmroth '83
generation of moments by Povzner's identities
- Desvilletes '93
if k -moment with $k > 2$ initially, then all moments are generated
- Wennberg '97
not even necessary to have an initial > 2 moment
- Bobylev '97
 L^1 estimates with exponential weight

Moderately soft potentials ($-2 < \gamma < 0$) and homogeneous

- Desvilletes '93
propagation of moments

Pointwise decay cannot be generated (cut-off)

Recall $B(r, \cos \theta) = |v_* - v|^\gamma b(\cos \theta)$

$\cos \theta = k \cdot \sigma$ with $k = (v_* - v)/|v_* - v|$

Gain and loss terms

$$\begin{aligned} Q(f, f) &= \int (f'_* f' - f f_*) B(r, \cos \theta) \, dv_* \, d\sigma \\ &= \underbrace{\int f'_* f' B(r, \cos \theta) \, dv_* \, d\sigma}_{Q_+(f, f)} - \underbrace{\int f f_* B(r, \cos \theta) \, dv_* \, d\sigma}_{Q_-(f, f)} \end{aligned}$$

Pointwise decay cannot be generated (cut-off)

Recall $Q(f, f) = Q_+(f, f) - Q_-(f, f)$

Loss term

$$\begin{aligned} Q_-(f, f) &= f(v) \int f(v_*) |v - v_*|^\gamma \, dv_* \underbrace{\left(\int b(\cos \theta) \, d\sigma \right)}_{C_b} \\ &= C_b f(f *_{v} |\cdot|^\gamma) \\ &\leq C_0 f(1 + |v|^\gamma) \end{aligned}$$

for some constant C_0 only depending on $\int f_0(1 + |v|^2) \, dv$.

A lower bound for f

$$\partial_t f = Q(f, f) = Q_+(f, f) - Q_-(f, f) \geq -C_0(1 + |v|^\gamma) f.$$

$$f \geq e^{-tC_0(1+|v|^\gamma)} f_0$$

$f_0 =$ initial condition

Generation of pointwise decay (non cut-off)

Non-homogeneous

- **Generation** of polynomial decay
- **Arbitrary** polynomial decay for $\gamma > 0$
- **Restricted** polynomial decay for $\gamma < 0$ (moderately soft)

Conditional to hydrodynamical bounds

- $0 \leq m_0 \leq \int f(t, x, v) dv \leq M_0$
- $\int f(t, x, v) |v|^2 dv \leq E_0$
- $\int f \ln f(t, x, v) dv \leq H_0$

Theorem (CI, Mouhot, Silvestre)

Let $\gamma + 2s \in [0, 2]$ and f be a solution of the Boltzmann equation such that • mass, energy and entropy are controlled •. Then

① (Generation) $\boxed{\text{If } \gamma > 0}$ then $\boxed{\forall q > 0}$, $\exists N, \beta > 0$

$$f(t, x, v) \leq N(1 + t^{-\beta}) \min(1, |v|^{-q})$$

② (Generation) $\boxed{\text{If } \gamma \leq 0}$, then

$$f(t, x, v) \leq N(1 + t^{-\frac{d}{2s}}) \min \left(1, |v|^{-\boxed{1 + \frac{d(\gamma+2s)}{2s}}} \right)$$

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③ (Propagation) $\boxed{\text{If } \gamma \leq 0}$, then $\forall q \geq \exists q_0$, $f_{in} \leq C \min(1, |v|^{-q})$

$$f(t, x, v) \leq N \min(1, |v|^{-q})$$

Conditional C^∞ smoothing effect (non-cutoff)

Reduce well-posedness of Boltzmann non-cutoff to hydro controls

This important equation is well-posed as long as mass and energy do not explode.

Conditional regularity for Boltzmann

- L^∞ bound (Silvestre)
- Local Hölder estimate (CI, Silvestre)
- Decay estimates for large velocities (CI, Mouhot, Silvestre)
- Schauder theory (CI, Silvestre)
- Global Hölder estimate and bootstrap (CI, Silvestre)

Why decay estimates?

- Boltzmann collision operator = non-local
- Control of coefficients only if control of tails

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Why decay estimates?

- Boltzmann collision operator = **non-local**
- Control of **coefficients** only if control of tails

Rewriting Boltzmann (non-cutoff)

$$Q(f, f)(v) = \int_{\mathbb{R}^n} \int_{\mathbb{S}^{d-1}} \left(f(v'_*)f(v') - f(v_*)f(v) \right) B(|v - v_*|, \theta) d\sigma dv_*.$$

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Rewriting Boltzmann (non-cutoff)

$$Q(f, f)(v) = \int_{\mathbb{R}^n} \left(f(v') - f(v) \right) K_f(v' - v) \, d v' + f(v)(c|\cdot|^\gamma * f).$$

Rewriting Boltzmann (non-cutoff)

$$Q(f, f)(v) = \underbrace{\int_{\mathbb{R}^n} \left(f(v') - f(v) \right) K_f(v' - v) dv'}_{\text{kernel}} + f(v)(c|\cdot|^\gamma * f).$$

- K_f is a kernel depending on f and the cross section B . It can be computed after a somewhat **difficult change of variables** in the integral.

Rewriting Boltzmann (non-cutoff)

$$Q(f, f)(v) = \underbrace{\int_{\mathbb{R}^n} \left(f(v') - f(v) \right) K_f(v' - v) \, dv'}_{\text{blue bracket}} + \underbrace{f(v)(c|\cdot|^\gamma * f)}_{\text{red bracket}}.$$

- K_f is a kernel depending on f and the cross section B . It can be computed after a somewhat **difficult change of variables** in the integral.
- c is a constant depending on the cross section B only. In particular, $f(c|\cdot|^\gamma * f)$ is bounded if $f \in L^1 \cap L^\infty$.

Rewriting Boltzmann (non-cutoff)

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Boltzmann rewritten

$$(\partial_t + v \cdot \nabla_x) f = \underbrace{L_v f}_{\text{blue}} + \underbrace{h}_{\text{red}}$$

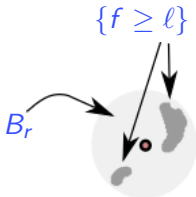
$$\text{with } L_v f = \int \left(f(v') - f(v) \right) K_f^{t,x,v}(v' - v) \, dv'$$

Non-degeneracy from hydro controls

$$Q(f, f) = \int (f(v') - f(v)) K_f(v' - v) dv' + h$$

$$K_f(v' - v) \simeq \frac{1}{|v' - v|^{d+2s}} \int_{w \perp v' - v} f(v + w) |w|^{\gamma+2s+1} dw$$

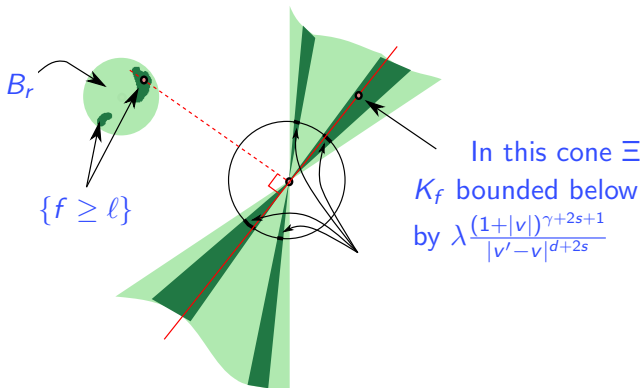
Hydro controls yields $\{f \geq \ell\}$ has positive measure close to 0



Non-degeneracy from hydro controls

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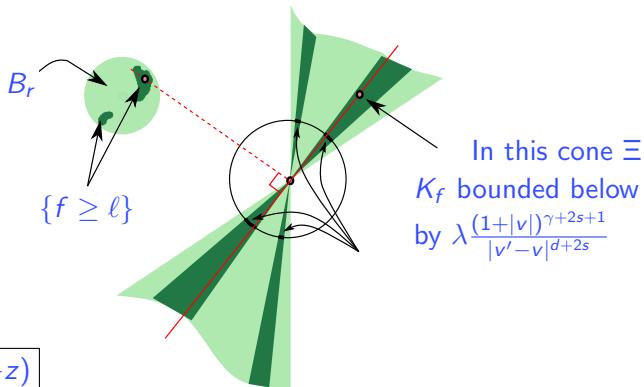
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Non-degeneracy from hydro controls

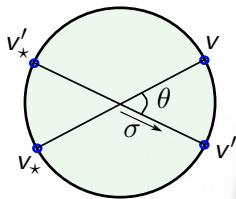
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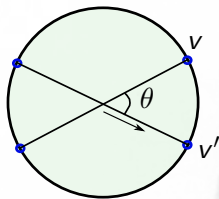
$$K_f(z) = K_f(-z)$$

The non-divergence form for Boltzmann



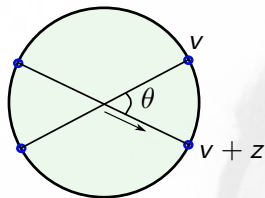
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The non-divergence form for Boltzmann



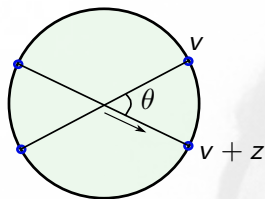
$$\begin{aligned} Q(f, f) &= \iint (f'_* f' - f_* f) B dv_* d\sigma \\ &= \int (f(v') - f(v)) K_f(v' - v) dv' + h \end{aligned}$$

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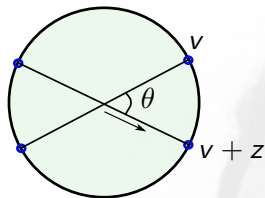


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Recall

$$K_f(z) = K_f(-z)$$

The non-divergence form for Boltzmann



$$\begin{aligned} Q(f, f) &= \iint (f'_* f' - f_* f) B dv_* d\sigma \\ &= \int (f(v') - f(v)) K_f(v' - v) dv' + h \\ &= \int (f(v + z) - f(v)) K_f(z) dz + h \end{aligned}$$

Recall

$$K_f(z) = K_f(-z)$$

$$= \frac{1}{2} \int (f(v + z) + f(v - z) - 2f(v)) K_f(v, z) dz + h$$

Barriers

$$Q(f, f) = \frac{1}{2} \int \left(f(v+z) + f(v-z) - 2f(v) \right) K_f(v, z) dz + h$$

See Silvestre '15, Schwab-Silvestre '16 (also CI-Silvestre '16)

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$$Q(f, f) = \frac{1}{2} \int \left(f(v+z) + f(v-z) - 2f(v) \right) K_f(v, z) dz + h$$

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- **Important property:** If f is smooth, then $Q(f, f)$ is bounded

$$Q(f, f) \leq \frac{1}{2} \|D^2 f\|_\infty \int_{|z| \leq 1} |z|^2 K_f(v, z) dz + \|f\|_\infty \int_{|z| \geq 1} K_f(v, z) dz + h$$

- **Barriers:** If $f \leq g$ with g smooth and $f(v_0) = g(v_0)$, then

$$Q(f, f) \leq \frac{1}{2} \int \left(g(v+z) + g(v-z) - 2g(v) \right) K_f(v, z) dz + h_g$$

at $v = v_0$

Maximum principle

Strategy

Goal: Prove $f(t, x, v) \leq \underbrace{N(t)(1 + |v|)^{-q}}_{g(v)}$

Method: consider the barrier g and use the maximum principle

- Choose N and q such that $f \leq g$ at initial time

- Consider the first time t_* where $f \leq g$ fails

$$\implies (\partial_t + v \nabla_x) f \geq (\partial_t + v \nabla_x) g \text{ at some } (t_*, x_*, v_*)$$

$$\implies N'(t_*)(1 + |v_*|)^{-q} \leq Q(f, f)$$

- contradict this inequality by using

$$\left| \begin{array}{l} f \leq g \\ f(t_*, x_*, v_*) = g(t_*, x_*, v_*) \end{array} \right.$$

The good term

$$Q(f, f) \simeq \int \frac{f(v') - f(v)}{|v' - v|^{d+2s}} \left\{ \int_{w \perp v' - v} f(v + w) |w|^{\gamma+2s+1} dw \right\} dv + h$$
$$= \mathcal{G}(f, f) + (\text{bad terms}) + h$$

$$\mathcal{G}(f, f) \simeq \int_{|v+w| \leq c_q |v|} f(v+w) |w|^{\gamma+2s} \left\{ \int_{(v'-v) \perp w} (f(v') - f(v)) \frac{dv'}{|v' - v|^{d-1+2s}} \right\} dw$$

$\mathcal{G}(f, f)$ = The kernel adds $|v|^{\gamma+2s}$ & the operator differentiates $2s$ times

The good term

$$Q(f, f) \simeq \int \frac{f(v') - f(v)}{|v' - v|^{d+2s}} \left\{ \int_{w \perp v' - v} f(v + w) |w|^{\gamma+2s+1} dw \right\} dv + h$$
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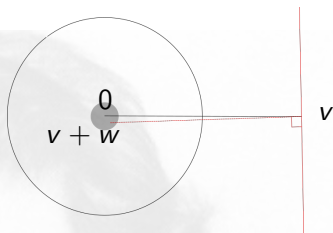
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The good term

Pick $c_q \ll 1$

- Grey circle of radius $c_q|v|$
- Circle of radius $\frac{1}{2}|v|$



$$\mathcal{G}(f, f) \simeq \int_{|v+w| \leq c_q|v|} f(v+w)|w|^{\gamma+2s} \left\{ \int_{(v'-v)_{\perp w}} (f(v') - f(v)) \frac{dv'}{|v' - v|^{d-1+2s}} \right\} dw$$

$$\lesssim \int_{|v+w| \leq c_q|v|} f(v+w)|w|^{\gamma+2s} \left\{ \int_{(v'-v)_{\perp w}} (g(v') - g(v)) \frac{dv'}{|v' - v|^{d-1+2s}} \right\} dw$$

if $v' \simeq v$, the function g is concave

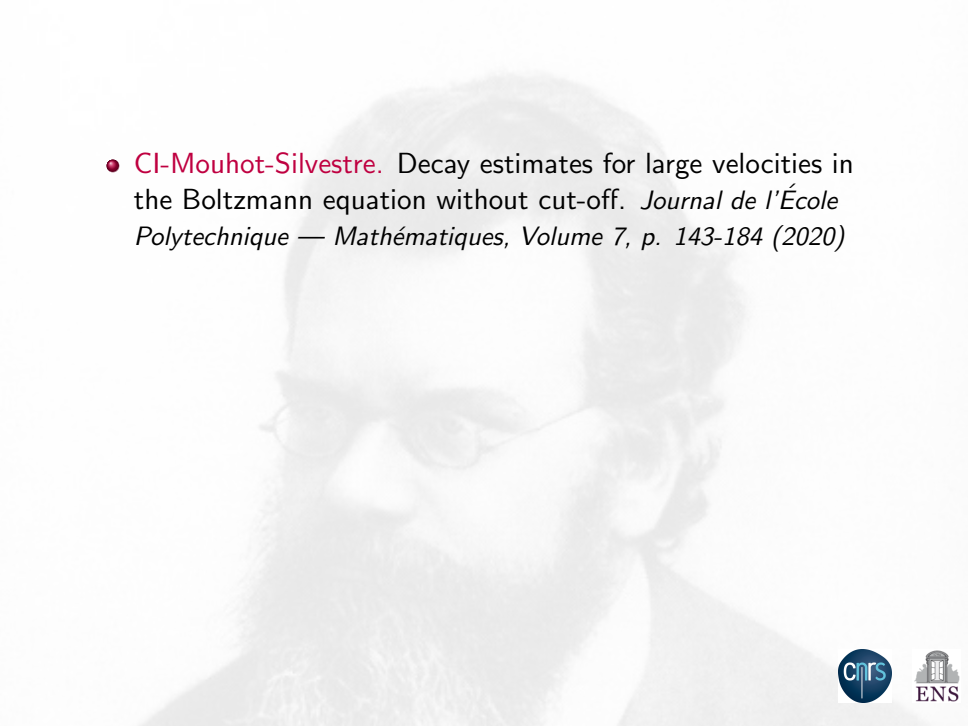
&

if not, use its decay

$$\lesssim \int_{|v+w| \leq c_q|v|} f(v+w)|w|^{\gamma+2s} \left\{ \int_{\substack{(v'-v)_{\perp w} \\ |v'-v| > r_q|v|}} (g(v') - g(v)) \frac{dv'}{|v' - v|^{d-1+2s}} \right\} dw$$

$$\lesssim -q^s |v|^{\gamma-q}$$

Recall the lower bound

- 
- **CI-Mouhot-Silvestre.** Decay estimates for large velocities in the Boltzmann equation without cut-off. *Journal de l'École Polytechnique — Mathématiques*, Volume 7, p. 143-184 (2020)

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Thank you for your attention

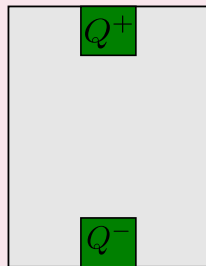
The Weak Harnack inequality

Theorem (CI-Silvestre)

Let f be a non-negative supersolution of the linear kinetic equation in $(-1, 0] \times B_1 \times B_1$. Assume that the kernel satisfies the *coercivity*, *upper bound*, *cancellation* and *non-degeneracy* conditions. Then

$$\|f\|_{L^\varepsilon(Q^-)} \leq C \left(\inf_{Q^+} f + \|h\|_\infty \right)$$

where ε and C only depends on dimension, λ and Λ .



Comments

- Weak Harnack implies Hölder

Weak Harnack inequality: comments

- **Kinetic parabolic equations:** $\partial_t f + v \cdot \nabla_x f = \nabla_v (A \nabla_v f)$
 - L^∞ bound: Pascucci, Polidoro (2004)
 - Hölder continuity: Wang, Zhang (2008, 2009, 2011)
 - Harnack inequality: Golse, CI, Mouhot, Vasseur (2015/2016)
- **Fact:** | the full Harnack inequality may not hold true
for integro-differential equations,
see for instance Bogdan, Sztonyk (2005)
- **Main ideas involved in the proof**
 - Use ideas from **de Giorgi** (first lemma)
 - **Transfer of regularity** from v to x (**kinetic theory**)
 - Use ideas from **non-div form equations** (Krylov, Safonov)

Weak Harnack inequality: further comments

- **First results** for **kinetic** integro-differential equations

Weak Harnack inequality: further comments

- **First results** for **kinetic** integro-differential equations
- **New regularity results** for integro-differential equations in divergence form, *even in the homogeneous case*
 - **very mild assumptions on the kernel**
 - nonsymmetric kernels with cancellation conditions

Weak Harnack inequality: further comments

- **First results** for **kinetic** integro-differential equations
- **New regularity results** for integro-differential equations in divergence form, *even in the homogeneous case*
 - **very mild assumptions on the kernel**
 - nonsymmetric kernels with cancellation conditions
- Our results apply to the Boltzmann equation without cut-off
 - This is the reason why we are forced to work with so mild assumptions on the nonsymmetric kernel