

①

$$(a) \quad y'''' + p_2 y'' + p_1 y' + p_0 y = 0$$

i)  $y_1, y_2, y_3$  are lin. dependent if 1

$$c_1 y_1(x) + c_2 y_2(x) + c_3 y_3(x) \equiv 0 \quad \text{for a non-trivial solution}$$

$c_1, c_2, c_3$

$$\Rightarrow c_1 y_1' + c_2 y_2' + c_3 y_3' \equiv 0$$

$$\Rightarrow c_1 y_1'' + c_2 y_2'' + c_3 y_3'' \equiv 0$$

which can be written

$$\begin{pmatrix} y_1 & y_2 & y_3 \\ y_1' & y_2' & y_3' \\ y_1'' & y_2'' & y_3'' \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \\ c_3 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \quad \text{2}$$

This has a non-trivial soln if the determinant = 0,

$$\therefore \text{if } W(y_1, y_2, y_3) = y_1 (y_2' y_3'' - y_2'' y_3') - y_2 (y_1' y_3'' - y_1'' y_3') + y_3 (y_1' y_2'' - y_1'' y_2') \equiv 0 \quad \text{2}$$

Wronskian

Hence if  $W \neq 0$ , the  $y_i(x)$  are lin. independent

$$ii) \quad y''' + 4y' = 0$$

$$\text{Let } w = y' \Rightarrow w'' + 4w = 0 \Rightarrow w = \tilde{c}_1 \cos 2x + \tilde{c}_2 \sin 2x$$

$$\Rightarrow y(x) = c_0 + c_1 \cos 2x + c_2 \sin 2x \quad \boxed{2}$$

$$\text{Here } y_1 = 1, y_2 = \cos 2x, y_3 = \sin 2x$$

$$\begin{aligned} \text{So } W(y_1, y_2, y_3) &= 1 \cdot (-2 \sin 2x) \cdot (-4 \sin 2x) - (-4 \cos 2x) \cdot (2 \cos 2x) \\ &= 8 \quad \checkmark \quad \boxed{2} \end{aligned}$$

(b)

$$(*) \begin{cases} Ly \equiv y''(x) + y(x) = f(x), & 0 < x < a \\ y(0) = 1, & y(a) = 0 \end{cases}$$

(i) (\*) will have a unique soln if there is no zero eig.-value, by FAT, we if

$$\begin{cases} y'' + y = 0 \\ y(0) = 0, y(a) = 0 \end{cases} \quad \text{has only trivial soln. (except)} \quad \boxed{2}$$

$$\begin{aligned} y &= A \cos x + B \sin x, & y(0) &= A = 0 \\ y(a) &= B \sin a = 0, & \text{so non-trivial soln} \end{aligned}$$

$$\text{if } a = n\pi, \quad n \in \mathbb{N}$$

Thus, if  $a \neq \pi, 2\pi, 3\pi, \dots$  then (\*) will have a unique soln.  $\boxed{2}$

(ii) For  $a = \pi$ ,  $y_0 = \sin x$  solves the homog. problem.

Fredholm Alt Thm: no unique soln  $\boxed{1}$

To get solvability cond, multiply (\*) by  $y_0$  and

$$\text{integrate: } \int_0^\pi (y'' + y) y_0 dx = \int_0^\pi f(x) y_0 dx$$

$$\text{LHS} = y' y_0 - y y_0' \Big|_0^\pi + \int_0^\pi y (y_0'' + y_0) dx \quad \boxed{1}$$

$$= y''(0) y_0'(0) = \cos 0 = 1$$

Thus, solvability cond is  $\int_0^{\pi} f(x) \sin x \, dx = 1$  [1]

(iii)  $a=1 \Rightarrow \exists$  unique soln.

$My = \int_0^1 g(x, \xi) y(\xi) \, d\xi$ . We want  $\{\mu, y(x)\}$  st

$$My = \mu y.$$

Consider  $Ly = \lambda y \Rightarrow \begin{cases} y'' + (1-\lambda)y = 0 \\ y(0) = y(1) = 0 \end{cases}$

1.1:  $y = A \cos \sqrt{1-\lambda} x + B \sin \sqrt{1-\lambda} x$  [3]

$$y(0) = A = 0 \quad y(1) = 0 \Rightarrow \sqrt{1-\lambda} = n\pi \Rightarrow \lambda_n = 1 - n^2\pi^2$$

are eig vals of  $L$ , w/  $y_n = \sin n\pi x$

Now, in Green's fn approach,  $Ly = \lambda y$  has soln (where  $\lambda y = "f(x)"$ )

$$y(x) = \int_0^1 g(x, \xi) \lambda y \, d\xi = \lambda \int_0^1 g(x, \xi) y(\xi) \, d\xi \quad [3]$$

Thus,  $My = \mu y$  has solns  $\mu_n = \frac{1}{\lambda_n} = \frac{1}{1 - n^2\pi^2}$ . [3]  
w/ same eig fns.

(2)

$$(a) \quad y'' - 2xy' + 2ny = 0$$

(i) For Sturm-Liouville form, multiply by  $\mu(x)$ , st

$$\mu y'' - 2x\mu y' = \frac{d}{dx}(\mu y')$$

$$\Rightarrow \text{need } \mu' = -2x\mu \Rightarrow \mu = e^{-x^2}$$

[3]

$$\text{SL form is thus } \frac{d}{dx} \left( e^{-x^2} y' \right) + 2ne^{-x^2} y = 0$$

(ii) Define  $Ly \equiv \frac{d}{dx} \left( e^{-x^2} y' \right)$  so that  $y_n$

is an eig. fn of  $Ly = -2ne^{-x^2} y$  w/ eig. val.  $n$ .

By properties of Sturm-Liouville, eig. fns  $y_n, y_m$  w/

$n \neq m$  will be orthogonal. However, since no

bdy cond. given, we need to establish a domain.

Supp.  $a < x < b$ , and consider

$$\int_a^b \underbrace{Ly_n}_{-2ne^{-x^2} y_n} \cdot y_m dx = \left. e^{-x^2} y_n'(x) y_m(x) \right|_a^b - \left. e^{-x^2} y_n' y_m \right|_a^b$$

$$+ \int_a^b y_n \underbrace{Ly_m}_{-2me^{-x^2} y_m} dx$$

[4]

Since  $y_n, y_m$  are polynomials, the bdy terms will vanish if  $a \rightarrow -\infty, b \rightarrow \infty$ . Then we

have  $2(n-m) \int_{-\infty}^{\infty} y_n(x) y_m(x) e^{-x^2} dx = 0$ . 2

Thus for  $n \neq m$  we have the orthog. relation

$$\int_{\mathbb{R}} y_n(x) y_m(x) e^{-x^2} dx = 0$$

$$(b) \quad \sin x \, y'' + 2(x-1)y' + \frac{2}{x}y = 0$$

i) Singular pts are  $x = n\pi$ ,  $n \in \mathbb{Z}$

The pt  $x_0$  is regular if the functions

$$P(x) = \frac{2(x-1)}{\sin x} \cdot (x-x_0), \quad Q(x) = \frac{2}{\sin x \cdot x} \cdot (x-x_0)^2$$

are analytic near  $x = x_0$ .

Near  $x = n\pi$ ,  $\sin x = (-1)^n (x - n\pi) + O((x - n\pi)^3)$  [3]

Hence  $P(x)$  is analytic for all  $x = n\pi$ , and so

$$\text{is } Q(x) \quad \left[ \begin{array}{l} P(x) \sim 2(x-1) + O((x - \frac{x_0}{\sin x})^2) \\ Q(x) \sim \frac{2(x-x_0)}{x} + O((x-x_0)^3) \end{array} \right]$$

as  $x \rightarrow x_0$

$\therefore$  The sing. pts are all regular.

Now check  $x = \infty$ : let  $t = \frac{1}{x}$ ,  $w(t) = y(x)$

$$\frac{d}{dx} = -t^2 \frac{d}{dt}, \quad \frac{d^2}{dx^2} = t^4 \frac{d^2}{dt^2} + 2t^3 \frac{d}{dt}$$

$$\Rightarrow \sin\left(\frac{1}{t}\right) \left( t^4 w'' + 2t^3 w' \right) + 2\left(\frac{1}{t} - 1\right) (-t^2 w') + 2t w = 0$$

Here  $t=0$  is an irreg. singular point, since eq

$$\frac{2t \cdot t^2}{\sin\left(\frac{1}{t}\right) t^4} \text{ is not analytic near } t=0$$

$\therefore x = \infty$  irreg. sing. pt

(iii) One soln given by  $y_1(x) = \sum_{k=0}^{\infty} a_k x^{k+\alpha_1} = \sum_{k=0}^{\infty} a_k x^{k+\alpha_2}$

For 2<sup>nd</sup> soln, since the  $\alpha_i$  differ by an integer, there may or may not be another Frobenius soln.

To check, we try  $y_2(x) = \sum_{k=0}^{\infty} b_k x^{k+1} = b_0 x + b_1 x^2 + \dots$

and see if there is a contradiction at  $b_1$ , assuming  $b_0 \neq 0$ .

Plugging in:  $(x - \frac{x^3}{6} \dots) (b_0 x + b_1 x^2 \dots)'' + (2x-2)(b_0 x + b_1 x^2 \dots)' + \frac{2}{x} (b_0 x + b_1 x^2 \dots) = 0$  2

original ODE, eq)

At  $\underline{x^0}$ :  $0 = 0 \checkmark$

At  $\underline{x^1}$ :  $2b_1 + 2b_0 - 4b_1 + 2b_1 = 0 \Rightarrow 2b_0 = 0 \rightarrow \leftarrow$

Contradiction  $\Rightarrow$  2<sup>nd</sup> soln is of form (Case II  $b_i$  in notes):

$$y_2(x) = y_1(x) \ln x + \sum_{k=0}^{\infty} c_k x^{k+1}$$



(ii) To find the indicial eqn for  $x=0$ , use

$$\lim_{x \rightarrow 0} \frac{2(x-1)}{\sin x} \cdot x = -2, \quad \lim_{x \rightarrow 0} \frac{2}{\sin x \cdot x} \cdot x^2 = 2$$

Thus, writing  $P(x) = \sum_{i=0}^{\infty} P_i x^i$ ,  $Q(x) = \sum_{i=0}^{\infty} Q_i x^i$

we have  $P_0 = -2$ ,  $Q_0 = 2$

The ODE is of form  $x^2 y'' + xP(x)y' + Q(x)y = 0$

Hence  $y = x^\alpha$  gives at lowest order

$$\alpha(\alpha-1) + P_0 \alpha + Q_0 = 0 \quad \leftarrow \text{Indicial Eqn}$$

$$\Rightarrow \alpha^2 - 3\alpha + 2 = (\alpha-2)(\alpha-1) = 0$$

$$\Rightarrow \alpha_1 = 2, \alpha_2 = 1 \quad \leftarrow \text{Indicial Exponents}$$

[3]

$$\textcircled{3} \quad (a) \quad \tan(\pi + \varepsilon) = \frac{\sin(\pi + \varepsilon)}{\cos(\pi + \varepsilon)}$$

$$\sin(\pi + \varepsilon) \sim -\varepsilon + \frac{\varepsilon^3}{6} + O(\varepsilon^5)$$

$$\cos(\pi + \varepsilon) \sim -1 + \frac{\varepsilon^2}{2} + O(\varepsilon^4) \Rightarrow \cos(\pi + \varepsilon)^{-1} \sim -1 - \frac{\varepsilon^2}{2} + O(\varepsilon^4)$$

$$\Rightarrow \tan(\pi + \varepsilon) \sim \left(-\varepsilon + \frac{\varepsilon^3}{6} + O(\varepsilon^5)\right) \left(-1 - \frac{\varepsilon^2}{2} + O(\varepsilon^4)\right) \quad \boxed{3}$$

$$= \varepsilon + \left(\frac{1}{2} - \frac{1}{6}\right)\varepsilon^3 + O(\varepsilon^5) = \varepsilon + \frac{1}{3}\varepsilon^3 + O(\varepsilon^5)$$

$\boxed{2}$

$$(b) \quad \begin{cases} xy' + y = \epsilon y^2, & x > 1 \\ y(1) = 1 + \epsilon^{\frac{1}{2}} \end{cases}$$

Seek expansion  $y \sim y_0 + \epsilon^{\frac{1}{2}} y_1 + \epsilon y_2 + \dots$

$$\underline{O(1)} \quad xy_0' + y_0 = 0 \quad \checkmark \quad y_0(1) = 1$$

$$\Rightarrow \ln y_0 = -\ln x + c \rightarrow y_0 = \frac{1}{x} \quad \boxed{2}$$

$$\underline{O(\epsilon^{\frac{1}{2}})} \quad xy_1' + y_1 = 0$$

$$y_1(1) = 1$$

$$\text{Same system} \rightarrow y_1 = \frac{1}{x} \quad \boxed{2}$$

$$\Rightarrow y \sim \frac{1}{x} + \epsilon^{\frac{1}{2}} \cdot \frac{1}{x} + O(\epsilon)$$

$$\underline{O(\epsilon)} \quad y_2 \text{ satisfies } \begin{cases} xy_2' + y_2 = y_0^2 = \frac{1}{x^2} \\ y_2(1) = 0 \end{cases} \quad \boxed{2}$$

$$\underline{O(\epsilon^{\frac{3}{2}})} \quad y_3 \text{ satisfies } \begin{cases} xy_3' + y_3 = 2y_0y_1 = \frac{2}{x^2} \\ y_3(1) = 0 \end{cases} \quad \boxed{2}$$

$$(c) \quad \begin{cases} \epsilon y'' + p(x)y' + y = 0 \\ y(-1) = 0, \quad y(1) = 1 \end{cases}$$

For outer soln, try  $y \sim y_0 + \epsilon y_1 + \dots$

$$\rightarrow p(x)y_0' + y_0 = 0 \Rightarrow y_0 = c e^{-\int \frac{1}{p(x)} dx}$$

Need to know where bdy layer is. Supp. layer is at  $x = x_0$  ( $x_0 = \pm 1$  TBD)

Set  $x = x_0 + \epsilon^\alpha X$ , define  $Y(X) = y(x)$

$$\rightarrow \epsilon^{1-2\alpha} Y'' + \underbrace{\left( p(x_0) + \epsilon^\alpha X p'(x_0) + \dots \right)}_{\text{[expanding } p \text{ about } x_0]} \epsilon^{-\alpha} Y' + Y = 0$$

To balance, we need  $1-2\alpha = -\alpha \Rightarrow \alpha = 1$

Then expanding  $Y \sim Y_0 + \epsilon Y_1 + \dots$  gives

$$\text{So ODE reads } Y'' + p(x_0)Y' + \epsilon \left( X p'(x_0)Y' + Y \right) + \dots = 0$$

$$\Rightarrow Y_0'' + p(x_0)Y_0' = 0 \Rightarrow Y_0 = c_1 + c_2 e^{-p(x_0)X}$$

If  $x_0 = 1$ , then matching implies  $X \rightarrow -\infty$

and  $Y_0 \rightarrow \infty$  since  $p(1) = 1$

So matching won't work.

If  $\underline{k_0 = -1}$ , matching as  $X \rightarrow \infty$ , so

$$e^{-p(x)X} \rightarrow 0$$

Hence bdy layer is at  $k_0 = -1$

$\Rightarrow$  impose BC  $[x = -1 \Rightarrow X = 0]$ : [2]

$$Y_0(0) = 0 \Rightarrow Y_0 = c_1 (1 - e^{-X})$$

Outer: impose  $y_0(1) = 1 \Rightarrow \cancel{c} = e^{\int \frac{1}{p(s)} ds} \Big|_{x=1}$

$$\Rightarrow y_0 = e^{-\int_1^x \frac{1}{p(s)} ds}$$

$\leftarrow$  Better route:  $\frac{1}{y_0} dy_0 = -\frac{1}{p(x)} dx$   
 $\int dx$  both sides  $\rightarrow$

$$\ln y_0 \Big|_1^x = -\int_1^x \frac{1}{p} ds$$

$$\rightarrow y_0 = e^{-\int_1^x \frac{1}{p} ds} \quad ]$$

To find  $c_1$ , require  $\lim_{X \rightarrow \infty} Y_0 = \lim_{x \rightarrow -1} y_0$

$$\Rightarrow c_1 = e^{\int_{-1}^1 \frac{1}{p(s)} ds}$$

[4]

$$\text{So } Y_{\text{inner}}(x) = e^{-\int_1^x \frac{1}{p} ds} \left( 1 - e^{-\frac{x+1}{e}} \right)$$