

Complex Geometry and Gravitational Instantons

Bernardo Araneda

School of Mathematics, Edinburgh
& Max Planck Institute, Potsdam

From Good Cuts to Celestial Holography

Oxford, July 2025

Joint work with Lars Andersson

Introduction

Newman was one of the pioneers in uncovering connections between complex geometry and gravitation,

Introduction

Newman was one of the pioneers in uncovering connections between complex geometry and gravitation, especially when considering asymptotically flat spaces, via the heavenly construction.

Introduction

Newman was one of the pioneers in uncovering connections between complex geometry and gravitation, especially when considering asymptotically flat spaces, via the heavenly construction.

Usually these connections involve self-dual gravity, particularly from the twistor perspective.

Introduction

Newman was one of the pioneers in uncovering connections between complex geometry and gravitation, especially when considering asymptotically flat spaces, via the heavenly construction.

Usually these connections involve self-dual gravity, particularly from the twistor perspective. But here we will focus on the strictly non-self-dual case.

Introduction

Newman was one of the pioneers in uncovering connections between complex geometry and gravitation, especially when considering asymptotically flat spaces, via the heavenly construction.

Usually these connections involve self-dual gravity, particularly from the twistor perspective. But here we will focus on the strictly non-self-dual case. New connections between asymptotic flatness and complex structures seem to emerge,

Introduction

Newman was one of the pioneers in uncovering connections between complex geometry and gravitation, especially when considering asymptotically flat spaces, via the heavenly construction.

Usually these connections involve self-dual gravity, particularly from the twistor perspective. But here we will focus on the strictly non-self-dual case. New connections between asymptotic flatness and complex structures seem to emerge, via a Witten-like identity.

Introduction

Newman was one of the pioneers in uncovering connections between complex geometry and gravitation, especially when considering asymptotically flat spaces, via the heavenly construction.

Usually these connections involve self-dual gravity, particularly from the twistor perspective. But here we will focus on the strictly non-self-dual case. New connections between asymptotic flatness and complex structures seem to emerge, via a Witten-like identity.

This will be the main message of this talk.

Introduction

Newman was one of the pioneers in uncovering connections between complex geometry and gravitation, especially when considering asymptotically flat spaces, via the heavenly construction.

Usually these connections involve self-dual gravity, particularly from the twistor perspective. But here we will focus on the strictly non-self-dual case. New connections between asymptotic flatness and complex structures seem to emerge, via a Witten-like identity.

This will be the main message of this talk. But the motivation is a recent counterexample to Euclidean Black Hole Uniqueness,

Introduction

Newman was one of the pioneers in uncovering connections between complex geometry and gravitation, especially when considering asymptotically flat spaces, via the heavenly construction.

Usually these connections involve self-dual gravity, particularly from the twistor perspective. But here we will focus on the strictly non-self-dual case. New connections between asymptotic flatness and complex structures seem to emerge, via a Witten-like identity.

This will be the main message of this talk. But the motivation is a recent counterexample to Euclidean Black Hole Uniqueness, i.e. a new AF gravitational instanton different from Kerr.

Introduction

Newman was one of the pioneers in uncovering connections between complex geometry and gravitation, especially when considering asymptotically flat spaces, via the heavenly construction.

Usually these connections involve self-dual gravity, particularly from the twistor perspective. But here we will focus on the strictly non-self-dual case. New connections between asymptotic flatness and complex structures seem to emerge, via a Witten-like identity.

This will be the main message of this talk. But the motivation is a recent counterexample to Euclidean Black Hole Uniqueness, i.e. a new AF gravitational instanton different from Kerr. We will be interested in integrable deformations of these spaces.

Gravitational Instantons

Gravitational instanton (M, g) : 4-dim, complete, non-singular, Ricci-flat Riemannian manifold, either compact or “asymptotically flat”

Gravitational Instantons

Gravitational instanton (M, g) : 4-dim, complete, non-singular, Ricci-flat Riemannian manifold, either compact or “asymptotically flat”

Name	Infinity	Examples
Compact	—	T^4 , $K3$
AE	\mathbb{E}^4	only \mathbb{E}^4
ALE	\mathbb{E}^4/Γ	Eguchi-Hanson $T^*\mathbb{CP}^1$
AF	$S^1 \times S^2$	Schw., Kerr $\mathbb{R}^2 \times S^2$
ALF	S^1 bundle over S^2	Taub-NUT \mathbb{R}^4 , Taub-bolt $\mathbb{CP}^2 \setminus \text{pt}$

(also ALG, ALH, ALG*, ALH*, Kasner,...)

Gravitational Instantons

Gravitational instanton (M, g) : 4-dim, complete, non-singular, Ricci-flat Riemannian manifold, either compact or “asymptotically flat”

Name	Infinity	Examples
Compact	—	T^4 , $K3$
AE	\mathbb{E}^4	only \mathbb{E}^4
ALE	\mathbb{E}^4/Γ	Eguchi-Hanson $T^*\mathbb{CP}^1$
AF	$S^1 \times S^2$	Schw., Kerr $\mathbb{R}^2 \times S^2$
ALF	S^1 bundle over S^2	Taub-NUT \mathbb{R}^4 , Taub-bolt $\mathbb{CP}^2 \setminus \text{pt}$

(also ALG, ALH, ALG*, ALH*, Kasner,...)

Some conjectures

- Besse (1987): “All compact instantons are hyperkähler”

Gravitational Instantons

Gravitational instanton (M, g) : 4-dim, complete, non-singular, Ricci-flat Riemannian manifold, either compact or “asymptotically flat”

Name	Infinity	Examples
Compact	—	T^4 , $K3$
AE	\mathbb{E}^4	only \mathbb{E}^4
ALE	\mathbb{E}^4/Γ	Eguchi-Hanson $T^*\mathbb{CP}^1$
AF	$S^1 \times S^2$	Schw., Kerr $\mathbb{R}^2 \times S^2$
ALF	S^1 bundle over S^2	Taub-NUT \mathbb{R}^4 , Taub-bolt $\mathbb{CP}^2 \backslash \text{pt}$

(also ALG, ALH, ALG*, ALH*, Kasner,...)

Some conjectures

- Besse (1987): “All compact instantons are hyperkähler”

open

Gravitational Instantons

Gravitational instanton (M, g) : 4-dim, complete, non-singular, Ricci-flat Riemannian manifold, either compact or “asymptotically flat”

Name	Infinity	Examples
Compact	—	T^4 , $K3$
AE	\mathbb{E}^4	only \mathbb{E}^4
ALE	\mathbb{E}^4/Γ	Eguchi-Hanson $T^*\mathbb{CP}^1$
AF	$S^1 \times S^2$	Schw., Kerr $\mathbb{R}^2 \times S^2$
ALF	S^1 bundle over S^2	Taub-NUT \mathbb{R}^4 , Taub-bolt $\mathbb{CP}^2 \setminus \text{pt}$

(also ALG, ALH, ALG*, ALH*, Kasner,...)

Some conjectures

- Besse (1987): “All compact instantons are hyperkähler”
- Nakajima (1990): “All ALE instantons are hyperkähler”

open

Gravitational Instantons

Gravitational instanton (M, g) : 4-dim, complete, non-singular, Ricci-flat Riemannian manifold, either compact or “asymptotically flat”

Name	Infinity	Examples
Compact	—	T^4 , $K3$
AE	\mathbb{E}^4	only \mathbb{E}^4
ALE	\mathbb{E}^4/Γ	Eguchi-Hanson $T^*\mathbb{CP}^1$
AF	$S^1 \times S^2$	Schw., Kerr $\mathbb{R}^2 \times S^2$
ALF	S^1 bundle over S^2	Taub-NUT \mathbb{R}^4 , Taub-bolt $\mathbb{CP}^2 \setminus \text{pt}$

(also ALG, ALH, ALG*, ALH*, Kasner,...)

Some conjectures

- Besse (1987): “All compact instantons are hyperkähler”
- Nakajima (1990): “All ALE instantons are hyperkähler”

open

open

Gravitational Instantons

Gravitational instanton (M, g) : 4-dim, complete, non-singular, Ricci-flat Riemannian manifold, either compact or “asymptotically flat”

Name	Infinity	Examples
Compact	—	T^4 , $K3$
AE	\mathbb{E}^4	only \mathbb{E}^4
ALE	\mathbb{E}^4/Γ	Eguchi-Hanson $T^*\mathbb{CP}^1$
AF	$S^1 \times S^2$	Schw., Kerr $\mathbb{R}^2 \times S^2$
ALF	S^1 bundle over S^2	Taub-NUT \mathbb{R}^4 , Taub-bolt $\mathbb{CP}^2 \setminus \text{pt}$

(also ALG, ALH, ALG*, ALH*, Kasner,...)

Some conjectures

- Besse (1987): “All compact instantons are hyperkähler”
- Nakajima (1990): “All ALE instantons are hyperkähler”
- BH Uniqueness (1980): “All AF instantons are in the Kerr family”

open

open

Gravitational Instantons

Gravitational instanton (M, g) : 4-dim, complete, non-singular, Ricci-flat Riemannian manifold, either compact or “asymptotically flat”

Name	Infinity	Examples
Compact	—	T^4 , $K3$
AE	\mathbb{E}^4	only \mathbb{E}^4
ALE	\mathbb{E}^4/Γ	Eguchi-Hanson $T^*\mathbb{CP}^1$
AF	$S^1 \times S^2$	Schw., Kerr $\mathbb{R}^2 \times S^2$
ALF	S^1 bundle over S^2	Taub-NUT \mathbb{R}^4 , Taub-bolt $\mathbb{CP}^2 \setminus \text{pt}$

(also ALG, ALH, ALG*, ALH*, Kasner,...)

Some conjectures

- Besse (1987): “All compact instantons are hyperkähler” open
- Nakajima (1990): “All ALE instantons are hyperkähler” open
- BH Uniqueness (1980): “All AF instantons are in the Kerr family” false!

The Chen-Teo instanton

Physics Letters B 703 (2011) 359–362

Contents lists available at [ScienceDirect](#)

Physics Letters B

www.elsevier.com/locate/physletb



A new AF gravitational instanton

Yu Chen, Edward Teo*

Department of Physics, National University of Singapore, Singapore 119260, Singapore

The Chen-Teo instanton

Physics Letters B 703 (2011) 359–362

Contents lists available at [ScienceDirect](#)

Physics Letters B

www.elsevier.com/locate/physletb



A new AF gravitational instanton

Yu Chen, Edward Teo*

Department of Physics, National University of Singapore, Singapore 119260, Singapore

- 2-parameter family
- Non-self-dual
- $M = \mathbb{CP}^2 \setminus S^1$
- Toric

The Chen-Teo instanton

Physics Letters B 703 (2011) 359–362

Contents lists available at [ScienceDirect](#)

Physics Letters B

www.elsevier.com/locate/physletb



A new AF gravitational instanton

Yu Chen, Edward Teo*

Department of Physics, National University of Singapore, Singapore 119260, Singapore

- 2-parameter family
- Non-self-dual
- $M = \mathbb{CP}^2 \setminus S^1$
- Toric

Remarks

- Given mass & ang. momentum, can't distinguish Kerr from Chen-Teo

The Chen-Teo instanton

Physics Letters B 703 (2011) 359–362

Contents lists available at ScienceDirect

Physics Letters B

www.elsevier.com/locate/physletb



A new AF gravitational instanton

Yu Chen, Edward Teo*

Department of Physics, National University of Singapore, Singapore 119260, Singapore

- 2-parameter family
- Non-self-dual
- $M = \mathbb{CP}^2 \setminus S^1$
- Toric

Remarks

- Given mass & ang. momentum, can't distinguish Kerr from Chen-Teo
- Constructed with Belinski-Zakharov method w/3 solitons

The Chen-Teo instanton

Physics Letters B 703 (2011) 359–362

Contents lists available at ScienceDirect

Physics Letters B

www.elsevier.com/locate/physletb



A new AF gravitational instanton

Yu Chen, Edward Teo*

Department of Physics, National University of Singapore, Singapore 119260, Singapore

- 2-parameter family
- Non-self-dual
- $M = \mathbb{CP}^2 \setminus S^1$
- Toric

Remarks

- Given mass & ang. momentum, can't distinguish Kerr from Chen-Teo
- Constructed with Belinski-Zakharov method w/3 solitons
- BZ method generates toric solutions with arbitrary number of solitons

⇒ there might be ∞ family of unknown AF instantons!

The Chen-Teo instanton

Physics Letters B 703 (2011) 359–362

Contents lists available at ScienceDirect

Physics Letters B

www.elsevier.com/locate/physletb



A new AF gravitational instanton

Yu Chen, Edward Teo*

Department of Physics, National University of Singapore, Singapore 119260, Singapore

- 2-parameter family
- Non-self-dual
- $M = \mathbb{CP}^2 \setminus S^1$
- Toric

Remarks

- Given mass & ang. momentum, can't distinguish Kerr from Chen-Teo
- Constructed with Belinski-Zakharov method w/3 solitons
- BZ method generates toric solutions with arbitrary number of solitons

⇒ there might be ∞ family of unknown AF instantons!

open

(cf. [Kunduri-Lucietti 2021])

The Chen-Teo instanton

Physics Letters B 703 (2011) 359–362

Contents lists available at ScienceDirect

Physics Letters B

www.elsevier.com/locate/physletb



A new AF gravitational instanton

Yu Chen, Edward Teo*

Department of Physics, National University of Singapore, Singapore 119260, Singapore

- 2-parameter family
- Non-self-dual
- $M = \mathbb{CP}^2 \setminus S^1$
- Toric

Remarks

- Given mass & ang. momentum, can't distinguish Kerr from Chen-Teo
- Constructed with Belinski-Zakharov method w/3 solitons
- BZ method generates toric solutions with arbitrary number of solitons

⇒ there might be ∞ family of unknown AF instantons!

open

(cf. [Kunduri-Lucietti 2021])

- Need to have better understanding of Moduli Spaces

Moduli Spaces

$$\mathcal{E}(M) = \{\text{Einstein metrics on } M\} / \text{Diff}(M)$$

Moduli Spaces

$$\mathcal{E}(M) = \{\text{Einstein metrics on } M\} / \text{Diff}(M)$$

More is known if M compact – T^4 , $K3$, del Pezzos, ... (cf. [LeBrun](#))

Moduli Spaces

$$\mathcal{E}(M) = \{\text{Einstein metrics on } M\} / \text{Diff}(M)$$

More is known if M compact – T^4 , $K3$, del Pezzos, ... (cf. [LeBrun](#))

Example: $M = \mathbb{R}^2 \times S^2$

Moduli Spaces

$$\mathcal{E}(M) = \{\text{Einstein metrics on } M\} / \text{Diff}(M)$$

More is known if M compact – T^4 , $K3$, del Pezzos, ... (cf. [LeBrun](#))

Example: $M = \mathbb{R}^2 \times S^2$

- M admits the Schwarzschild instanton g_m .

Moduli Spaces

$$\mathcal{E}(M) = \{\text{Einstein metrics on } M\} / \text{Diff}(M)$$

More is known if M compact – T^4 , $K3$, del Pezzos, ... (cf. [LeBrun](#))

Example: $M = \mathbb{R}^2 \times S^2$

- M admits the Schwarzschild instanton g_m . Are there others?

Moduli Spaces

$$\mathcal{E}(M) = \{\text{Einstein metrics on } M\} / \text{Diff}(M)$$

More is known if M compact – T^4 , $K3$, del Pezzos, ... (cf. [LeBrun](#))

Example: $M = \mathbb{R}^2 \times S^2$

- M admits the Schwarzschild instanton g_m . Are there others?
- g_m has ALF **infinitesimal deformations** (soln to linearised Einstein).

Moduli Spaces

$$\mathcal{E}(M) = \{\text{Einstein metrics on } M\} / \text{Diff}(M)$$

More is known if M compact – T^4 , $K3$, del Pezzos, ... (cf. [LeBrun](#))

Example: $M = \mathbb{R}^2 \times S^2$

- M admits the Schwarzschild instanton g_m . Are there others?
- g_m has ALF **infinitesimal deformations** (soln to linearised Einstein). These integrate to a curve $g_{m(s),a(s)} \in \mathcal{E}(M) \leadsto$ Kerr instanton

Moduli Spaces

$$\mathcal{E}(M) = \{\text{Einstein metrics on } M\} / \text{Diff}(M)$$

More is known if M compact – T^4 , $K3$, del Pezzos, ... (cf. [LeBrun](#))

Example: $M = \mathbb{R}^2 \times S^2$

- M admits the Schwarzschild instanton g_m . Are there others?
- g_m has ALF **infinitesimal deformations** (soln to linearised Einstein). These integrate to a curve $g_{m(s),a(s)} \in \mathcal{E}(M) \leadsto$ Kerr instanton

For general M :

Moduli Spaces

$$\mathcal{E}(M) = \{\text{Einstein metrics on } M\} / \text{Diff}(M)$$

More is known if M compact – T^4 , $K3$, del Pezzos, ... (cf. [LeBrun](#))

Example: $M = \mathbb{R}^2 \times S^2$

- M admits the Schwarzschild instanton g_m . Are there others?
- g_m has ALF **infinitesimal deformations** (soln to linearised Einstein). These integrate to a curve $g_{m(s),a(s)} \in \mathcal{E}(M) \leadsto$ Kerr instanton

For general M :

- **Open problem:** determine whether infinitesimal deformations are integrable

Moduli Spaces

$$\mathcal{E}(M) = \{\text{Einstein metrics on } M\} / \text{Diff}(M)$$

More is known if M compact – T^4 , $K3$, del Pezzos, ... (cf. [LeBrun](#))

Example: $M = \mathbb{R}^2 \times S^2$

- M admits the Schwarzschild instanton g_m . Are there others?
- g_m has ALF **infinitesimal deformations** (soln to linearised Einstein). These integrate to a curve $g_{m(s),a(s)} \in \mathcal{E}(M) \leadsto$ Kerr instanton

For general M :

- **Open problem:** determine whether infinitesimal deformations are integrable
- $\mathcal{E}(M)$ may be a non-smooth variety

Moduli Spaces

$$\mathcal{E}(M) = \{\text{Einstein metrics on } M\} / \text{Diff}(M)$$

More is known if M compact – T^4 , $K3$, del Pezzos, ... (cf. [LeBrun](#))

Example: $M = \mathbb{R}^2 \times S^2$

- M admits the Schwarzschild instanton g_m . Are there others?
- g_m has ALF **infinitesimal deformations** (soln to linearised Einstein). These integrate to a curve $g_{m(s),a(s)} \in \mathcal{E}(M) \leadsto$ Kerr instanton

For general M :

- **Open problem:** determine whether infinitesimal deformations are integrable
- $\mathcal{E}(M)$ may be a non-smooth variety
- $\mathcal{E}(M)$ is **integrable** at g_0 if any infinitesimal deformation integrates to a curve of Einstein metrics in $\mathcal{E}(M)$

Integrability

Examples

- Calabi-Yau (cf. [Besse 1987](#)); self-dual Yang-Mills ([Atiyah, Hitchin, Singer 1978](#)); twistor theory & deformations of complex structures ([Penrose, Kodaira](#))



Integrability

Examples

- Calabi-Yau (cf. Besse 1987); self-dual Yang-Mills (Atiyah, Hitchin, Singer 1978); twistor theory & deformations of complex structures (Penrose, Kodaira) ✓
- $\mathbb{CP}^1 \times \mathbb{CP}^{2n}$ has infinitesimal deformations, but they are not tangent to curves in $\mathcal{E}(M)$ (Koiso 1982) ✗

Integrability

Examples

- Calabi-Yau (cf. Besse 1987); self-dual Yang-Mills (Atiyah, Hitchin, Singer 1978); twistor theory & deformations of complex structures (Penrose, Kodaira) ✓
- $\mathbb{CP}^1 \times \mathbb{CP}^{2n}$ has infinitesimal deformations, but they are not tangent to curves in $\mathcal{E}(M)$ (Koiso 1982) ✗
- If g_0 = vacuum space-time with compact Cauchy surface and Killing field $\Rightarrow \mathcal{E}(M)$ is a quadratic variety near g_0 (Fischer, Marsden, Moncrief 1980) ✗

Integrability

Examples

- Calabi-Yau (cf. Besse 1987); self-dual Yang-Mills (Atiyah, Hitchin, Singer 1978); twistor theory & deformations of complex structures (Penrose, Kodaira) ✓
- $\mathbb{CP}^1 \times \mathbb{CP}^{2n}$ has infinitesimal deformations, but they are not tangent to curves in $\mathcal{E}(M)$ (Koiso 1982) ✗
- If g_0 = vacuum space-time with compact Cauchy surface and Killing field $\Rightarrow \mathcal{E}(M)$ is a quadratic variety near g_0 (Fischer, Marsden, Moncrief 1980) ✗

This talk

- Integrability of $\mathcal{E}(M)$ is **open** for general Einstein manifolds

Integrability

Examples

- Calabi-Yau (cf. Besse 1987); self-dual Yang-Mills (Atiyah, Hitchin, Singer 1978); twistor theory & deformations of complex structures (Penrose, Kodaira) ✓
- $\mathbb{CP}^1 \times \mathbb{CP}^{2n}$ has infinitesimal deformations, but they are not tangent to curves in $\mathcal{E}(M)$ (Koiso 1982) ✗
- If g_0 = vacuum space-time with compact Cauchy surface and Killing field $\Rightarrow \mathcal{E}(M)$ is a quadratic variety near g_0 (Fischer, Marsden, Moncrief 1980) ✗

This talk

- Integrability of $\mathcal{E}(M)$ is **open** for general Einstein manifolds
- Is the Moduli Space of ALF instantons integrable?

Integrability

Examples

- Calabi-Yau (cf. Besse 1987); self-dual Yang-Mills (Atiyah, Hitchin, Singer 1978); twistor theory & deformations of complex structures (Penrose, Kodaira) ✓
- $\mathbb{CP}^1 \times \mathbb{CP}^{2n}$ has infinitesimal deformations, but they are not tangent to curves in $\mathcal{E}(M)$ (Koiso 1982) ✗
- If g_0 = vacuum space-time with compact Cauchy surface and Killing field $\Rightarrow \mathcal{E}(M)$ is a quadratic variety near g_0 (Fischer, Marsden, Moncrief 1980) ✗

This talk

- Integrability of $\mathcal{E}(M)$ is **open** for general Einstein manifolds
- Is the Moduli Space of ALF instantons integrable?
- We will be able to answer this using complex structures

Complex structures and black holes

Special geometry of black holes

- [Newman & Janis \(1965\)](#): Kerr is a “complex translation” of Schwarzschild

Complex structures and black holes

Special geometry of black holes

- [Newman & Janis \(1965\)](#): Kerr is a “complex translation” of Schwarzschild
- [Flaherty \(1974, 1976\)](#): Kerr has Lorentzian Hermitian structure

Complex structures and black holes

Special geometry of black holes

- [Newman & Janis \(1965\)](#): Kerr is a “complex translation” of Schwarzschild
- [Flaherty \(1974, 1976\)](#): Kerr has Lorentzian Hermitian structure
- [Walker & Penrose \(1970\)](#): Kerr has non-degenerate valence-2 Killing spinor

Complex structures and black holes

Special geometry of black holes

- [Newman & Janis \(1965\)](#): Kerr is a “complex translation” of Schwarzschild
- [Flaherty \(1974, 1976\)](#): Kerr has Lorentzian Hermitian structure
- [Walker & Penrose \(1970\)](#): Kerr has non-degenerate valence-2 Killing spinor

Riemannian versions

- Non-degenerate valence-2 Killing spinor \Leftrightarrow conformally Kähler
[\[Dunajski & Tod 2009; Pontecorvo 1992\]](#)

Complex structures and black holes

Special geometry of black holes

- [Newman & Janis \(1965\)](#): Kerr is a “complex translation” of Schwarzschild
- [Flaherty \(1974, 1976\)](#): Kerr has Lorentzian Hermitian structure
- [Walker & Penrose \(1970\)](#): Kerr has non-degenerate valence-2 Killing spinor

Riemannian versions

- Non-degenerate valence-2 Killing spinor \Leftrightarrow conformally Kähler
[\[Dunajski & Tod 2009; Pontecorvo 1992\]](#)
- Conformally Kähler \Rightarrow Hermitian

Complex structures and black holes

Special geometry of black holes

- [Newman & Janis \(1965\)](#): Kerr is a “complex translation” of Schwarzschild
- [Flaherty \(1974, 1976\)](#): Kerr has Lorentzian Hermitian structure
- [Walker & Penrose \(1970\)](#): Kerr has non-degenerate valence-2 Killing spinor

Riemannian versions

- Non-degenerate valence-2 Killing spinor \Leftrightarrow conformally Kähler
[\[Dunajski & Tod 2009; Pontecorvo 1992\]](#)
- Conformally Kähler \Rightarrow Hermitian
- Ricci-flat + algebraically special \Rightarrow conformally Kähler

Complex structures and instantons

- [Aksteiner & Andersson \(2021\)](#): Chen-Teo is Hermitian (non-Kähler), and it has no Lorentzian sections.

Complex structures and instantons

- [Aksteiner & Andersson \(2021\)](#): Chen-Teo is Hermitian (non-Kähler), and it has no Lorentzian sections.
- [Biquard & Gauduchon \(2021\)](#): The only Hermitian (non-Kähler), ALF, toric instantons are

Taub-NUT, **Taub-bolt**, **Kerr**, **Chen-Teo**.

Complex structures and instantons

- Aksteiner & Andersson (2021): Chen-Teo is Hermitian (non-Kähler), and it has no Lorentzian sections.
- Biquard & Gauduchon (2021): The only Hermitian (non-Kähler), ALF, toric instantons are

Taub-NUT, Taub-bolt, Kerr, Chen-Teo.

- Biquard, Gauduchon & LeBrun (2023): If (M, g_0) is ALF Hermitian and toric, and g is another Ricci-flat metric on M sufficiently C_1^3 -close to g_0 , then g is also ALF Hermitian.

(N.B: This does not cover infinitesimal deformations)

Complex structures and instantons

- Aksteiner & Andersson (2021): Chen-Teo is Hermitian (non-Kähler), and it has no Lorentzian sections.
- Biquard & Gauduchon (2021): The only Hermitian (non-Kähler), ALF, toric instantons are

Taub-NUT, Taub-bolt, Kerr, Chen-Teo.

- Biquard, Gauduchon & LeBrun (2023): If (M, g_0) is ALF Hermitian and toric, and g is another Ricci-flat metric on M sufficiently C_1^3 -close to g_0 , then g is also ALF Hermitian.

(N.B: This does not cover infinitesimal deformations)

- Li (2023): ALF Hermitian instantons are toric.

Main results

To look for further solutions, we want to see if the instantons can be smoothly deformed, and if these deformations are integrable.

Main results

To look for further solutions, we want to see if the instantons can be smoothly deformed, and if these deformations are integrable.

Theorem A (Stability)

[B.A, L. Andersson & M. Dahl, 2024]

Let (M, g) be an ALF Hermitian non-Kähler gravitational instanton. Then:

- Infinitesimal deformations satisfy the Teukolsky equation.*
- There are no non-trivial ALF solutions to the Teukolsky eqn.*

Main results

To look for further solutions, we want to see if the instantons can be smoothly deformed, and if these deformations are integrable.

Theorem A (Stability)

[B.A, L. Andersson & M. Dahl, 2024]

Let (M, g) be an ALF Hermitian non-Kähler gravitational instanton. Then:

- Infinitesimal deformations satisfy the Teukolsky equation.*
- There are no non-trivial ALF solutions to the Teukolsky eqn.*

Theorem B (Integrability)

[B.A & L. Andersson, 2025]

Let (M, g) be an ALF Hermitian non-Kähler gravitational instanton. Then:

- Infinitesimal deformations are conformally Kähler.*
- The Moduli Space $\mathcal{E}(M)$ is integrable at g .*

Main results

To look for further solutions, we want to see if the instantons can be smoothly deformed, and if these deformations are integrable.

Theorem A (Stability)

[B.A, L. Andersson & M. Dahl, 2024]

Let (M, g) be an ALF Hermitian non-Kähler gravitational instanton. Then:

- *Infinitesimal deformations satisfy the Teukolsky equation.*
- *There are no non-trivial ALF solutions to the Teukolsky eqn.*

Theorem B (Integrability)

[B.A & L. Andersson, 2025]

Let (M, g) be an ALF Hermitian non-Kähler gravitational instanton. Then:

- *Infinitesimal deformations are conformally Kähler.*
- *The Moduli Space $\mathcal{E}(M)$ is integrable at g .*

We will focus on **Theorem B**, using “Witten-like” identity (in the sense of PMT)

Notation

Signature (+ + ++)

Notation

Signature $(+ + + +)$

- Spin group $SU(2) \times SU(2)$, symplectic spin bundles (\mathbb{S}, ϵ) , (\mathbb{S}', ϵ')

$$\phi^A \in \mathbb{S}, \quad \psi^{A'} \in \mathbb{S}', \quad \phi_B = \epsilon_{AB} \phi^A, \quad \psi_{B'} = \epsilon_{A'B'} \psi^{A'}$$

Notation

Signature $(+ + + +)$

- Spin group $SU(2) \times SU(2)$, symplectic spin bundles (\mathbb{S}, ϵ) , (\mathbb{S}', ϵ')

$$\phi^A \in \mathbb{S}, \quad \psi^{A'} \in \mathbb{S}', \quad \phi_B = \epsilon_{AB} \phi^A, \quad \psi_{B'} = \epsilon_{A'B'} \psi^{A'}$$

- Tangent bundle $TM \cong \mathbb{S} \otimes \mathbb{S}'$, with metric $g_{ab} = \epsilon_{AB} \epsilon_{A'B'}$

Notation

Signature $(+ + + +)$

- Spin group $SU(2) \times SU(2)$, symplectic spin bundles (\mathbb{S}, ϵ) , (\mathbb{S}', ϵ')

$$\phi^A \in \mathbb{S}, \quad \psi^{A'} \in \mathbb{S}', \quad \phi_B = \epsilon_{AB} \phi^A, \quad \psi_{B'} = \epsilon_{A'B'} \psi^{A'}$$

- Tangent bundle $TM \cong \mathbb{S} \otimes \mathbb{S}'$, with metric $g_{ab} = \epsilon_{AB} \epsilon_{A'B'}$
- Self-dual 2-forms $\Lambda_+^2 \cong \mathbb{S} \odot \mathbb{S} = \text{span}(\varphi_{AB}^i)$, with L-C connection

$$\nabla_a \varphi_{BC}^i = \Gamma_a^i{}_j \varphi_{BC}^j, \quad i, j = 1, 2, 3$$

Notation

Signature $(+ + + +)$

- Spin group $SU(2) \times SU(2)$, symplectic spin bundles (\mathbb{S}, ϵ) , (\mathbb{S}', ϵ')

$$\phi^A \in \mathbb{S}, \quad \psi^{A'} \in \mathbb{S}', \quad \phi_B = \epsilon_{AB} \phi^A, \quad \psi_{B'} = \epsilon_{A'B'} \psi^{A'}$$

- Tangent bundle $TM \cong \mathbb{S} \otimes \mathbb{S}'$, with metric $g_{ab} = \epsilon_{AB} \epsilon_{A'B'}$
- Self-dual 2-forms $\Lambda_+^2 \cong \mathbb{S} \odot \mathbb{S} = \text{span}(\varphi_{AB}^i)$, with L-C connection

$$\nabla_a \varphi_{BC}^i = \Gamma_a^i{}_j \varphi_{BC}^j, \quad i, j = 1, 2, 3$$

- If $|\varphi|^2 = 1$, then $J^a{}_b = \sqrt{2} \varphi_B^A \delta_{B'}^{A'}$ is almost-complex structure

Notation

Signature $(+ + + +)$

- Spin group $SU(2) \times SU(2)$, symplectic spin bundles (\mathbb{S}, ϵ) , (\mathbb{S}', ϵ')

$$\phi^A \in \mathbb{S}, \quad \psi^{A'} \in \mathbb{S}', \quad \phi_B = \epsilon_{AB} \phi^A, \quad \psi_{B'} = \epsilon_{A'B'} \psi^{A'}$$

- Tangent bundle $TM \cong \mathbb{S} \otimes \mathbb{S}'$, with metric $g_{ab} = \epsilon_{AB} \epsilon_{A'B'}$
- Self-dual 2-forms $\Lambda_+^2 \cong \mathbb{S} \odot \mathbb{S} = \text{span}(\varphi_{AB}^i)$, with L-C connection

$$\nabla_a \varphi_{BC}^i = \Gamma_a^i{}_j \varphi_{BC}^j, \quad i, j = 1, 2, 3$$

- If $|\varphi|^2 = 1$, then $J^a{}_b = \sqrt{2} \varphi_B^A \delta_B^{A'}$ is almost-complex structure
- (M, g) is Kähler iff $\nabla_a \varphi_{BC} = 0$

Notation

Signature $(+ + + +)$

- Spin group $SU(2) \times SU(2)$, symplectic spin bundles (\mathbb{S}, ϵ) , (\mathbb{S}', ϵ')

$$\phi^A \in \mathbb{S}, \quad \psi^{A'} \in \mathbb{S}', \quad \phi_B = \epsilon_{AB} \phi^A, \quad \psi_{B'} = \epsilon_{A'B'} \psi^{A'}$$

- Tangent bundle $TM \cong \mathbb{S} \otimes \mathbb{S}'$, with metric $g_{ab} = \epsilon_{AB} \epsilon_{A'B'}$
- Self-dual 2-forms $\Lambda_+^2 \cong \mathbb{S} \odot \mathbb{S} = \text{span}(\varphi_{AB}^i)$, with L-C connection

$$\nabla_a \varphi_{BC}^i = \Gamma_a^i{}_j \varphi_{BC}^j, \quad i, j = 1, 2, 3$$

- If $|\varphi|^2 = 1$, then $J^a{}_b = \sqrt{2} \varphi_B^A \delta_{B'}^{A'}$ is almost-complex structure
- (M, g) is Kähler iff $\nabla_a \varphi_{BC} = 0$
- Weyl tensor

$$W_{abcd} = \underbrace{\Psi_{ABCD} \epsilon_{A'B'} \epsilon_{C'D'}}_{W_{abcd}^+} + \underbrace{\tilde{\Psi}_{A'B'C'D'} \epsilon_{AB} \epsilon_{CD}}_{W_{abcd}^-}$$

Notation

Signature $(+ + + +)$

- Spin group $SU(2) \times SU(2)$, symplectic spin bundles (\mathbb{S}, ϵ) , (\mathbb{S}', ϵ')

$$\phi^A \in \mathbb{S}, \quad \psi^{A'} \in \mathbb{S}', \quad \phi_B = \epsilon_{AB} \phi^A, \quad \psi_{B'} = \epsilon_{A'B'} \psi^{A'}$$

- Tangent bundle $TM \cong \mathbb{S} \otimes \mathbb{S}'$, with metric $g_{ab} = \epsilon_{AB} \epsilon_{A'B'}$
- Self-dual 2-forms $\Lambda_+^2 \cong \mathbb{S} \odot \mathbb{S} = \text{span}(\varphi_{AB}^i)$, with L-C connection

$$\nabla_a \varphi_{BC}^i = \Gamma_a^i{}_j \varphi_{BC}^j, \quad i, j = 1, 2, 3$$

- If $|\varphi|^2 = 1$, then $J^a{}_b = \sqrt{2} \varphi_B^A \delta_B^{A'}$ is almost-complex structure
- (M, g) is Kähler iff $\nabla_a \varphi_{BC} = 0$
- Weyl tensor

$$W_{abcd} = \underbrace{\Psi_{ABCD} \epsilon_{A'B'} \epsilon_{C'D'}}_{W_{abcd}^+} + \underbrace{\tilde{\Psi}_{A'B'C'D'} \epsilon_{AB} \epsilon_{CD}}_{W_{abcd}^-}$$

Remark: A global spin structure is not assumed.

Eigenspinors

Consider the map (cf. [Penrose & Rindler Vol.2](#))

$$\Psi_{AB}{}^{CD} : \mathbb{S}^2 \rightarrow \mathbb{S}^2$$

Eigenspinors

Consider the map (cf. [Penrose & Rindler Vol.2](#))

$$\Psi_{AB}{}^{CD} : \mathbb{S}^2 \rightarrow \mathbb{S}^2$$

- Self-adjoint \Rightarrow orthonormal basis of real eigenspinors

Eigenspinors

Consider the map (cf. [Penrose & Rindler Vol.2](#))

$$\Psi_{AB}{}^{CD} : \mathbb{S}^2 \rightarrow \mathbb{S}^2$$

- Self-adjoint \Rightarrow orthonormal basis of real eigenspinors
- Trace-free $\Rightarrow \lambda_1 + \lambda_2 + \lambda_3 = 0$
- Can choose $\lambda_1 \leq 0$ and $\lambda_3 \geq 0$

Eigenspinors

Consider the map (cf. [Penrose & Rindler Vol.2](#))

$$\Psi_{AB}{}^{CD} : \mathbb{S}^2 \rightarrow \mathbb{S}^2$$

- Self-adjoint \Rightarrow orthonormal basis of real eigenspinors
- Trace-free $\Rightarrow \lambda_1 + \lambda_2 + \lambda_3 = 0$
- Can choose $\lambda_1 \leq 0$ and $\lambda_3 \geq 0$
- If $\hat{g}_{ab} = \Omega^2 g_{ab}$, then

$$\hat{\Psi}_{AB}{}^{CD} = \Omega^{-2} \Psi_{AB}{}^{CD}, \quad \hat{\lambda}_i = \Omega^{-2} \lambda_i$$

Eigenspinors

Consider the map (cf. [Penrose & Rindler Vol.2](#))

$$\Psi_{AB}{}^{CD} : \mathbb{S}^2 \rightarrow \mathbb{S}^2$$

- Self-adjoint \Rightarrow orthonormal basis of real eigenspinors
- Trace-free $\Rightarrow \lambda_1 + \lambda_2 + \lambda_3 = 0$
- Can choose $\lambda_1 \leq 0$ and $\lambda_3 \geq 0$
- If $\hat{g}_{ab} = \Omega^2 g_{ab}$, then

$$\hat{\Psi}_{AB}{}^{CD} = \Omega^{-2} \Psi_{AB}{}^{CD}, \quad \hat{\lambda}_i = \Omega^{-2} \lambda_i$$

We will work with the following objects (cf. [LeBrun \(2019\)](#))

Eigenspinors

Consider the map (cf. [Penrose & Rindler Vol.2](#))

$$\Psi_{AB}{}^{CD} : \mathbb{S}^2 \rightarrow \mathbb{S}^2$$

- Self-adjoint \Rightarrow orthonormal basis of real eigenspinors
- Trace-free $\Rightarrow \lambda_1 + \lambda_2 + \lambda_3 = 0$
- Can choose $\lambda_1 \leq 0$ and $\lambda_3 \geq 0$
- If $\hat{g}_{ab} = \Omega^2 g_{ab}$, then

$$\hat{\Psi}_{AB}{}^{CD} = \Omega^{-2} \Psi_{AB}{}^{CD}, \quad \hat{\lambda}_i = \Omega^{-2} \lambda_i$$

We will work with the following objects (cf. [LeBrun \(2019\)](#))

- g_{ab} with L-C connection ∇_a and trace-free Ricci tensor E_{ab}

Eigenspinors

Consider the map (cf. [Penrose & Rindler Vol.2](#))

$$\Psi_{AB}{}^{CD} : \mathbb{S}^2 \rightarrow \mathbb{S}^2$$

- Self-adjoint \Rightarrow orthonormal basis of real eigenspinors
- Trace-free $\Rightarrow \lambda_1 + \lambda_2 + \lambda_3 = 0$
- Can choose $\lambda_1 \leq 0$ and $\lambda_3 \geq 0$
- If $\hat{g}_{ab} = \Omega^2 g_{ab}$, then

$$\hat{\Psi}_{AB}{}^{CD} = \Omega^{-2} \Psi_{AB}{}^{CD}, \quad \hat{\lambda}_i = \Omega^{-2} \lambda_i$$

We will work with the following objects (cf. [LeBrun \(2019\)](#))

- g_{ab} with L-C connection ∇_a and trace-free Ricci tensor E_{ab}
- $\hat{g}_{ab} = \Omega^2 g_{ab}$ with L-C connection $\hat{\nabla}_a$ and $\Omega = \hat{\lambda}_3$

Eigenspinors

Consider the map (cf. [Penrose & Rindler Vol.2](#))

$$\Psi_{AB}{}^{CD} : \mathbb{S}^2 \rightarrow \mathbb{S}^2$$

- Self-adjoint \Rightarrow orthonormal basis of real eigenspinors
- Trace-free $\Rightarrow \lambda_1 + \lambda_2 + \lambda_3 = 0$
- Can choose $\lambda_1 \leq 0$ and $\lambda_3 \geq 0$
- If $\hat{g}_{ab} = \Omega^2 g_{ab}$, then

$$\hat{\Psi}_{AB}{}^{CD} = \Omega^{-2} \Psi_{AB}{}^{CD}, \quad \hat{\lambda}_i = \Omega^{-2} \lambda_i$$

We will work with the following objects (cf. [LeBrun \(2019\)](#))

- g_{ab} with L-C connection ∇_a and trace-free Ricci tensor E_{ab}
- $\hat{g}_{ab} = \Omega^2 g_{ab}$ with L-C connection $\hat{\nabla}_a$ and $\Omega = \hat{\lambda}_3$
- $\hat{\varphi}_{AB}$ with $\hat{\Psi}_{AB}{}^{CD} \hat{\varphi}_{CD} = \hat{\lambda}_3 \hat{\varphi}_{AB}$ and $|\hat{\varphi}|^2 = 1$

Key identity

Lemma

With the above definitions, let (M, \hat{g}_{ab}) be arbitrary ($\hat{\lambda}_3 \neq 0$), and set

$$\hat{S}_B^{A'} = \hat{\nabla}^{AA'} \hat{\varphi}_{AB}, \quad \hat{V}^a = \hat{\varphi}^{AB} \hat{S}_B^{A'}, \quad \hat{P}^{abcd} = \hat{\varphi}^{(AB} \hat{\varphi}^{CD)} \hat{\epsilon}^{A'B'} \hat{\epsilon}^{C'D'}$$

Key identity

Lemma

With the above definitions, let (M, \hat{g}_{ab}) be arbitrary ($\hat{\lambda}_3 \neq 0$), and set

$$\hat{S}_B^{A'} = \hat{\nabla}^{AA'} \hat{\varphi}_{AB}, \quad \hat{V}^a = \hat{\varphi}^{AB} \hat{S}_B^{A'}, \quad \hat{P}^{abcd} = \hat{\varphi}^{(AB} \hat{\varphi}^{CD)} \hat{\epsilon}^{A'B'} \hat{\epsilon}^{C'D'}$$

Then the following identity holds:

$$\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$$

Key identity

Lemma

With the above definitions, let (M, \hat{g}_{ab}) be arbitrary ($\hat{\lambda}_3 \neq 0$), and set

$$\hat{S}_B^{A'} = \hat{\nabla}^{AA'} \hat{\varphi}_{AB}, \quad \hat{V}^a = \hat{\varphi}^{AB} \hat{S}_B^{A'}, \quad \hat{P}^{abcd} = \hat{\varphi}^{(AB} \hat{\varphi}^{CD)} \hat{\epsilon}^{A'B'} \hat{\epsilon}^{C'D'}$$

Then the following identity holds:

$$\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$$

where

$$\mathcal{A} = |\hat{S}|^2 + \frac{1}{6} |\hat{\nabla} \hat{\varphi}|^2 + \frac{1}{3\hat{\lambda}_3} \left[(\hat{\lambda}_1 - \hat{\lambda}_2)^2 - 2\hat{\lambda}_1 |\hat{\Gamma}^3{}_1|^2 - 2\hat{\lambda}_2 |\hat{\Gamma}^3{}_2|^2 \right]$$

$$\mathcal{B} = -\frac{1}{3} \hat{P}^{abcd} \hat{\nabla}_a (\hat{\lambda}_3^{-1} \nabla_d E_{bc})$$

Example

- We have $\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$, and

Example

- We have $\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$, and

$$\mathcal{A} = \underbrace{|\hat{S}|^2 + \frac{1}{6}|\hat{\nabla}\hat{\varphi}|^2}_{\geq 0} + \underbrace{\frac{1}{3\hat{\lambda}_3}}_{\geq 0} \left[\underbrace{(\hat{\lambda}_1 - \hat{\lambda}_2)^2}_{\geq 0} \underbrace{-2\hat{\lambda}_1|\hat{\Gamma}^3_1|^2}_{\geq 0} \underbrace{-2\hat{\lambda}_2|\hat{\Gamma}^3_2|^2}_{??} \right]$$

Example

- We have $\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$, and

$$\mathcal{A} = \underbrace{|\hat{S}|^2 + \frac{1}{6}|\hat{\nabla}\hat{\varphi}|^2}_{\geq 0} + \underbrace{\frac{1}{3\hat{\lambda}_3}}_{\geq 0} \left[\underbrace{(\hat{\lambda}_1 - \hat{\lambda}_2)^2}_{\geq 0} \underbrace{-2\hat{\lambda}_1|\hat{\Gamma}^3{}_1|^2}_{\geq 0} \underbrace{-2\hat{\lambda}_2|\hat{\Gamma}^3{}_2|^2}_{??} \right]$$

$$\mathcal{B} = -\frac{1}{3}\hat{P}^{abcd}\hat{\nabla}_a(\hat{\lambda}_3^{-1}\nabla_d E_{bc})$$

Example

- We have $\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$, and

$$\mathcal{A} = \underbrace{|\hat{S}|^2 + \frac{1}{6}|\hat{\nabla}\hat{\varphi}|^2}_{\geq 0} + \underbrace{\frac{1}{3\hat{\lambda}_3}}_{\geq 0} \left[\underbrace{(\hat{\lambda}_1 - \hat{\lambda}_2)^2}_{\geq 0} \underbrace{-2\hat{\lambda}_1|\hat{\Gamma}^3_1|^2}_{\geq 0} \underbrace{-2\hat{\lambda}_2|\hat{\Gamma}^3_2|^2}_{??} \right]$$

$$\mathcal{B} = -\frac{1}{3}\hat{P}^{abcd}\hat{\nabla}_a(\hat{\lambda}_3^{-1}\nabla_d E_{bc})$$

- Integrate over M , use Stokes, and add assumptions,

Example

- We have $\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$, and

$$\mathcal{A} = \underbrace{|\hat{S}|^2 + \frac{1}{6}|\hat{\nabla}\hat{\varphi}|^2}_{\geq 0} + \underbrace{\frac{1}{3\hat{\lambda}_3}}_{\geq 0} \left[\underbrace{(\hat{\lambda}_1 - \hat{\lambda}_2)^2}_{\geq 0} \underbrace{-2\hat{\lambda}_1|\hat{\Gamma}^3_1|^2}_{\geq 0} \underbrace{-2\hat{\lambda}_2|\hat{\Gamma}^3_2|^2}_{??} \right]$$

$$\mathcal{B} = -\frac{1}{3}\hat{P}^{abcd}\hat{\nabla}_a(\hat{\lambda}_3^{-1}\nabla_d E_{bc})$$

- Integrate over M , use Stokes, and add assumptions,

	g Einstein	\Rightarrow	$\mathcal{B} = 0$
for example:	$\hat{\lambda}_2 < 0$	\Rightarrow	$\mathcal{A} \geq 0$
	Boundary term = 0	\Rightarrow	$\mathcal{A} = 0$

Example

- We have $\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$, and

$$\mathcal{A} = \underbrace{|\hat{S}|^2}_{\geq 0} + \underbrace{\frac{1}{6}|\hat{\nabla}\hat{\varphi}|^2}_{\geq 0} + \underbrace{\frac{1}{3\hat{\lambda}_3}[(\hat{\lambda}_1 - \hat{\lambda}_2)^2]}_{\geq 0} \underbrace{- 2\hat{\lambda}_1|\hat{\Gamma}^3_1|^2}_{\geq 0} \underbrace{- 2\hat{\lambda}_2|\hat{\Gamma}^3_2|^2}_{??}$$

$$\mathcal{B} = -\frac{1}{3}\hat{P}^{abcd}\hat{\nabla}_a(\hat{\lambda}_3^{-1}\nabla_d E_{bc})$$

- Integrate over M , use Stokes, and add assumptions,

$$\begin{array}{lll} g \text{ Einstein} & \Rightarrow & \mathcal{B} = 0 \\ \text{for example: } \hat{\lambda}_2 < 0 & \Rightarrow & \mathcal{A} \geq 0 \\ \text{Boundary term} = 0 & \Rightarrow & \mathcal{A} = 0 \end{array}$$

- Then $\hat{\nabla}\hat{\varphi} = 0$, so \hat{g} is Kähler, i.e. g is conformally Kähler

Example

- We have $\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$, and

$$\mathcal{A} = \underbrace{|\hat{S}|^2}_{\geq 0} + \underbrace{\frac{1}{6}|\hat{\nabla}\hat{\varphi}|^2}_{\geq 0} + \underbrace{\frac{1}{3\hat{\lambda}_3}[(\hat{\lambda}_1 - \hat{\lambda}_2)^2]}_{\geq 0} \underbrace{- 2\hat{\lambda}_1|\hat{\Gamma}^3_1|^2}_{\geq 0} \underbrace{- 2\hat{\lambda}_2|\hat{\Gamma}^3_2|^2}_{??}$$

$$\mathcal{B} = -\frac{1}{3}\hat{P}^{abcd}\hat{\nabla}_a(\hat{\lambda}_3^{-1}\nabla_d E_{bc})$$

- Integrate over M , use Stokes, and add assumptions,

$$\begin{array}{lll} g \text{ Einstein} & \Rightarrow & \mathcal{B} = 0 \\ \text{for example: } \hat{\lambda}_2 < 0 & \Rightarrow & \mathcal{A} \geq 0 \\ \text{Boundary term} = 0 & \Rightarrow & \mathcal{A} = 0 \end{array}$$

- Then $\hat{\nabla}\hat{\varphi} = 0$, so \hat{g} is Kähler, i.e. g is conformally Kähler

Remark

This result was obtained by [Wu \(2019\)](#) and [LeBrun \(2019\)](#) in the case that (M, g) is compact, Einstein and $\det W^+ > 0$

ALF curves

Recall key identity $\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$.

Lemma

Let $g(s)$ be a curve of ALF metrics on M , with $g_0 := g(0)$ Hermitian instanton and $\delta g := \frac{dg}{ds}|_{s=0}$ infinitesimal deformation. Then

ALF curves

Recall key identity $\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$.

Lemma

Let $g(s)$ be a curve of ALF metrics on M , with $g_0 := g(0)$ Hermitian instanton and $\delta g := \frac{dg}{ds}|_{s=0}$ infinitesimal deformation. Then

$$\int_M (\hat{\nabla}_a \hat{V}^a)(s) = \mathcal{O}(s^3)$$

$$\int_M \mathcal{A}(s) = \mathcal{O}(s^2)$$

$$\int_M \mathcal{B}(s) = \mathcal{O}(s^3)$$

ALF curves

Recall key identity $\hat{\nabla}_a \hat{V}^a = \mathcal{A} + \mathcal{B}$.

Lemma

Let $g(s)$ be a curve of ALF metrics on M , with $g_0 := g(0)$ Hermitian instanton and $\delta g := \frac{dg}{ds}|_{s=0}$ infinitesimal deformation. Then

$$\int_M (\hat{\nabla}_a \hat{V}^a)(s) = \mathcal{O}(s^3)$$

$$\int_M \mathcal{A}(s) = \mathcal{O}(s^2)$$

$$\int_M \mathcal{B}(s) = \mathcal{O}(s^3)$$

and the almost-complex structure $J^a_b = \sqrt{2} \hat{\varphi}^A_B \delta^{A'}_{B'}$ satisfies

$$(\hat{\nabla} J)(s) = \mathcal{O}(s^2)$$

Sketch of proof

- Set $\hat{J}_{ab} = \sqrt{2} \hat{\varphi}_{AB} \hat{\epsilon}_{A'B'}$, then

$$\int_M \hat{\nabla}_a \hat{V}^a = \left\{ \oint_{\partial M} \hat{J} \wedge \hat{*} d\delta \hat{J} \right\} s + \left\{ \oint_{\partial M} \hat{J} \wedge \hat{*} d\delta^2 \hat{J} \right\} \frac{s^2}{2!} + \mathcal{O}(s^3)$$

Sketch of proof

- Set $\hat{J}_{ab} = \sqrt{2} \hat{\varphi}_{AB} \hat{\epsilon}_{A'B'}$, then

$$\int_M \hat{\nabla}_a \hat{V}^a = \left\{ \oint_{\partial M} \hat{J} \wedge \hat{*} d\delta \hat{J} \right\} s + \left\{ \oint_{\partial M} \hat{J} \wedge \hat{*} d\delta^2 \hat{J} \right\} \frac{s^2}{2!} + \mathcal{O}(s^3)$$

- Using that g_0 is ALF Hermitian instanton,

$$\mathcal{A} = \left\{ \frac{|\delta(\hat{\nabla} \hat{J})|^2}{24} + |\delta \hat{S}|^2 + \frac{(\delta \hat{\lambda}_1 - \delta \hat{\lambda}_2)^2}{3 \hat{\lambda}_3} + \frac{(|\delta \hat{\Gamma}_1^3|^2 + |\delta \hat{\Gamma}_2^3|^2)}{3} \right\} \frac{s^2}{2!} + \mathcal{O}(s^3)$$

Sketch of proof

- Set $\hat{J}_{ab} = \sqrt{2} \hat{\varphi}_{AB} \hat{\epsilon}_{A'B'}$, then

$$\int_M \hat{\nabla}_a \hat{V}^a = \left\{ \oint_{\partial M} \hat{J} \wedge \hat{*} d\delta \hat{J} \right\} s + \left\{ \oint_{\partial M} \hat{J} \wedge \hat{*} d\delta^2 \hat{J} \right\} \frac{s^2}{2!} + \mathcal{O}(s^3)$$

- Using that g_0 is ALF Hermitian instanton,

$$\mathcal{A} = \left\{ \frac{|\delta(\hat{\nabla} \hat{J})|^2}{24} + |\delta \hat{S}|^2 + \frac{(\delta \hat{\lambda}_1 - \delta \hat{\lambda}_2)^2}{3 \hat{\lambda}_3} + \frac{(|\delta \hat{\Gamma}^3_1|^2 + |\delta \hat{\Gamma}^3_2|^2)}{3} \right\} \frac{s^2}{2!} + \mathcal{O}(s^3)$$

- Using that δg is infinitesimal deformation,

$$\int_M \mathcal{B} = \left\{ -\frac{1}{3} \oint_{\partial M} (\hat{\lambda}_3^{-1} \hat{P}^{abcd} \nabla_d \delta^2 E_{bc}) d\Sigma_a \right\} \frac{s^2}{2!} + \mathcal{O}(s^3)$$

Sketch of proof

- Set $\hat{J}_{ab} = \sqrt{2} \hat{\varphi}_{AB} \hat{e}_{A'B'}$, then

$$\int_M \hat{\nabla}_a \hat{V}^a = \left\{ \oint_{\partial M} \hat{J} \wedge \star d\hat{J} \right\} s + \left\{ \oint_{\partial M} \hat{J} \wedge \star d\delta^2 \hat{J} \right\} \frac{s^2}{2!} + \mathcal{O}(s^3)$$

- Using that g_0 is ALF Hermitian instanton,

$$\mathcal{A} = \left\{ \frac{|\delta(\hat{\nabla} \hat{J})|^2}{24} + |\delta \hat{S}|^2 + \frac{(\delta \hat{\lambda}_1 - \delta \hat{\lambda}_2)^2}{3 \hat{\lambda}_3} + \frac{(|\delta \hat{\Gamma}_1^3|^2 + |\delta \hat{\Gamma}_2^3|^2)}{3} \right\} \frac{s^2}{2!} + \mathcal{O}(s^3)$$

- Using that δg is infinitesimal deformation,

$$\int_M \mathcal{B} = \left\{ -\frac{1}{3} \oint_{\partial M} (\hat{\lambda}_3^{-1} \hat{P}^{abcd} \nabla_d \delta^2 E_{bc}) d\Sigma_a \right\} \frac{s^2}{2!} + \mathcal{O}(s^3)$$

- ALF fall-off at ∞ implies that boundary terms vanish $\Rightarrow \delta^2 \mathcal{A} = 0$
 $\Rightarrow \delta(\hat{\nabla} \hat{J}) = 0 \quad \Rightarrow \quad \hat{\nabla} J = \mathcal{O}(s^2)$

Integrability of $\mathcal{E}(M)$

We then have an ALF curve $(g(s), J(s))$ on M such that

Integrability of $\mathcal{E}(M)$

We then have an ALF curve $(g(s), J(s))$ on M such that

$$(g_0, J_0) = \text{ALF Hermitian instanton}$$

$$\text{Ric}[g(s)] = \mathcal{O}(s^2)$$

$$(\hat{\nabla} J)(s) = \mathcal{O}(s^2)$$

Integrability of $\mathcal{E}(M)$

We then have an ALF curve $(g(s), J(s))$ on M such that

$$(g_0, J_0) = \text{ALF Hermitian instanton}$$

$$\text{Ric}[g(s)] = \mathcal{O}(s^2)$$

$$(\hat{\nabla} J)(s) = \mathcal{O}(s^2)$$

Therefore:

Integrability of $\mathcal{E}(M)$

We then have an ALF curve $(g(s), J(s))$ on M such that

$$(g_0, J_0) = \text{ALF Hermitian instanton}$$

$$\text{Ric}[g(s)] = \mathcal{O}(s^2)$$

$$(\hat{\nabla} J)(s) = \mathcal{O}(s^2)$$

Therefore:

- The infinitesimal deformation δg solves the ALF linearised Hermitian-Einstein system, i.e. points in the Hermitian direction

Integrability of $\mathcal{E}(M)$

We then have an ALF curve $(g(s), J(s))$ on M such that

$$(g_0, J_0) = \text{ALF Hermitian instanton}$$

$$\text{Ric}[g(s)] = \mathcal{O}(s^2)$$

$$(\hat{\nabla} J)(s) = \mathcal{O}(s^2)$$

Therefore:

- The infinitesimal deformation δg solves the ALF linearised Hermitian-Einstein system, i.e. points in the Hermitian direction
- ALF Hermitian instantons are classified and topology is fixed, so δg must be deformation within same family \Rightarrow perturbation w.r.t moduli

Integrability of $\mathcal{E}(M)$

We then have an ALF curve $(g(s), J(s))$ on M such that

$$(g_0, J_0) = \text{ALF Hermitian instanton}$$

$$\text{Ric}[g(s)] = \mathcal{O}(s^2)$$

$$(\hat{\nabla} J)(s) = \mathcal{O}(s^2)$$

Therefore:

- The infinitesimal deformation δg solves the ALF linearised Hermitian-Einstein system, i.e. points in the Hermitian direction
- ALF Hermitian instantons are classified and topology is fixed, so δg must be deformation within same family \Rightarrow perturbation w.r.t moduli
- Hermitian moduli spaces are smooth $\Rightarrow \delta g$ integrates to a curve of Ricci-flat metrics in $\mathcal{E}(M)$.

Conclusions

- There is a counterexample to Euclidean BH Uniqueness; currently unknown if there are others.

Conclusions

- There is a counterexample to Euclidean BH Uniqueness; currently unknown if there are others.
- We studied infinitesimal deformations. Connected to integrability of moduli spaces —open for general Einstein manifolds.

Conclusions

- There is a counterexample to Euclidean BH Uniqueness; currently unknown if there are others.
- We studied infinitesimal deformations. Connected to integrability of moduli spaces —open for general Einstein manifolds.
- Using Witten-like identity, proved the deformations have complex structures, moduli space is integrable, instantons cannot be smoothly deformed away.

Conclusions

- There is a counterexample to Euclidean BH Uniqueness; currently unknown if there are others.
- We studied infinitesimal deformations. Connected to integrability of moduli spaces —open for general Einstein manifolds.
- Using Witten-like identity, proved the deformations have complex structures, moduli space is integrable, instantons cannot be smoothly deformed away.
- New connections: asymptotic flatness, complex structures, global regularity.

Conclusions

- There is a counterexample to Euclidean BH Uniqueness; currently unknown if there are others.
- We studied infinitesimal deformations. Connected to integrability of moduli spaces —open for general Einstein manifolds.
- Using Witten-like identity, proved the deformations have complex structures, moduli space is integrable, instantons cannot be smoothly deformed away.
- New connections: asymptotic flatness, complex structures, global regularity.

Some questions

- Other asymptotics? Beyond Ricci-flatness?

Conclusions

- There is a counterexample to Euclidean BH Uniqueness; currently unknown if there are others.
- We studied infinitesimal deformations. Connected to integrability of moduli spaces —open for general Einstein manifolds.
- Using Witten-like identity, proved the deformations have complex structures, moduli space is integrable, instantons cannot be smoothly deformed away.
- New connections: asymptotic flatness, complex structures, global regularity.

Some questions

- Other asymptotics? Beyond Ricci-flatness?
- Key identity \leadsto **deduce** \exists complex structures on instantons?

Conclusions

- There is a counterexample to Euclidean BH Uniqueness; currently unknown if there are others.
- We studied infinitesimal deformations. Connected to integrability of moduli spaces —open for general Einstein manifolds.
- Using Witten-like identity, proved the deformations have complex structures, moduli space is integrable, instantons cannot be smoothly deformed away.
- New connections: asymptotic flatness, complex structures, global regularity.

Some questions

- Other asymptotics? Beyond Ricci-flatness?
- Key identity \leadsto **deduce** \exists complex structures on instantons?
- Classification of toric instantons? (Twistor approach?)

Conclusions

- There is a counterexample to Euclidean BH Uniqueness; currently unknown if there are others.
- We studied infinitesimal deformations. Connected to integrability of moduli spaces —open for general Einstein manifolds.
- Using Witten-like identity, proved the deformations have complex structures, moduli space is integrable, instantons cannot be smoothly deformed away.
- New connections: asymptotic flatness, complex structures, global regularity.

Some questions

- Other asymptotics? Beyond Ricci-flatness?
- Key identity \leadsto **deduce** \exists complex structures on instantons?
- Classification of toric instantons? (Twistor approach?)
- Applications to **Lorentzian** Relativity? (e.g. BH stability?)

Conclusions

- There is a counterexample to Euclidean BH Uniqueness; currently unknown if there are others.
- We studied infinitesimal deformations. Connected to integrability of moduli spaces —open for general Einstein manifolds.
- Using Witten-like identity, proved the deformations have complex structures, moduli space is integrable, instantons cannot be smoothly deformed away.
- New connections: asymptotic flatness, complex structures, global regularity.

Some questions

- Other asymptotics? Beyond Ricci-flatness?
- Key identity \leadsto **deduce** \exists complex structures on instantons?
- Classification of toric instantons? (Twistor approach?)
- Applications to **Lorentzian** Relativity? (e.g. BH stability?)

Thank you!