

Surface Energies Arising in Microscopic Modeling of Martensitic Transformations

Angkana Rüland

(joint work with G. Kitavtsev and S. Luckhaus)

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1 Introduction

2 The Model Hamiltonian

3 Results

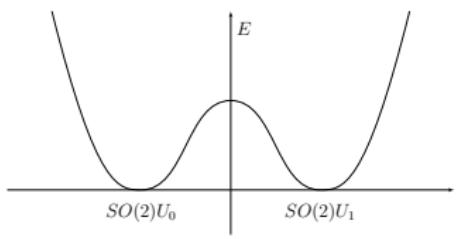
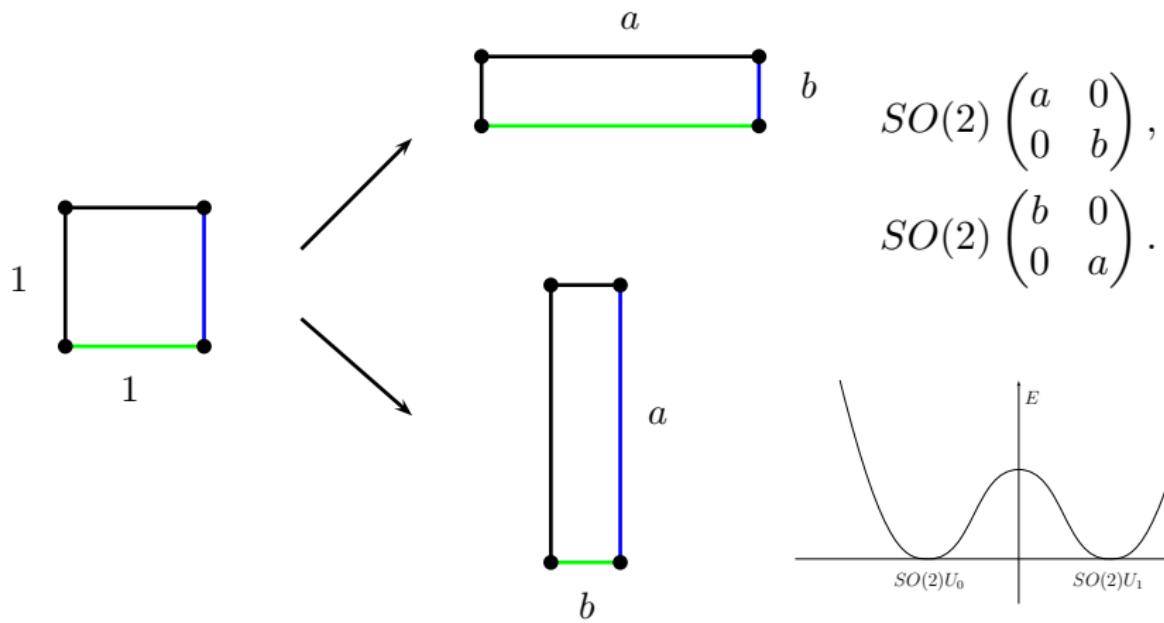
4 Ideas of Proof

5 Conclusion and Outlook

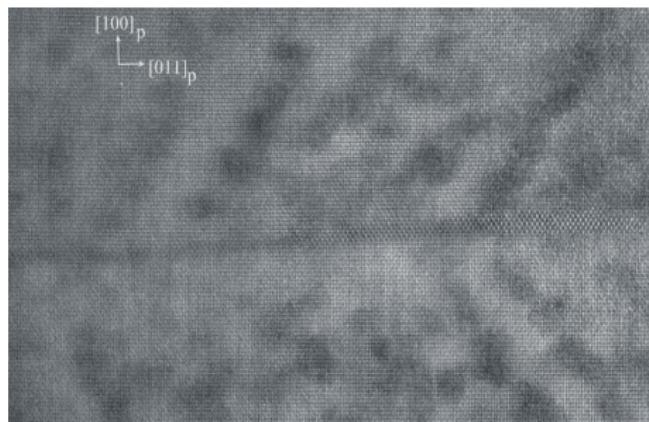
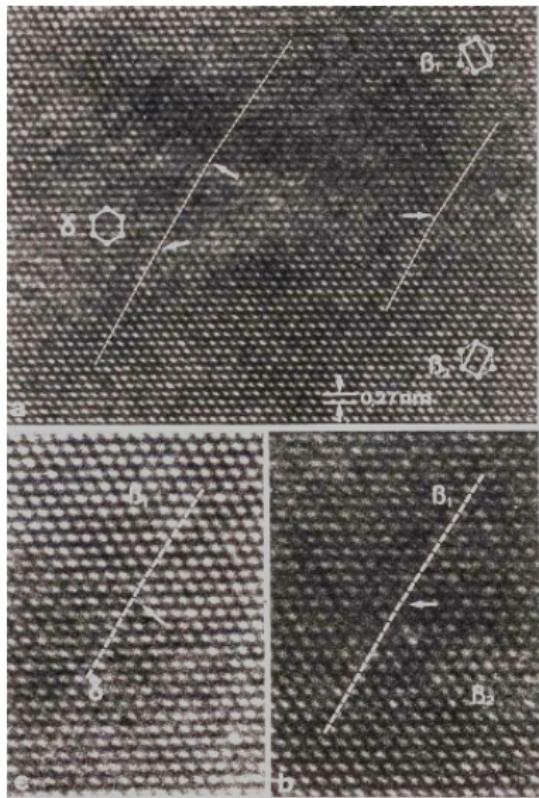
Questions and Objectives

- ▶ Justification of continuum models as limits of discrete models (closer to first principles)? How are continuum models related to the atomistic Hamiltonians governing the behavior of the atoms in a crystal?
- ▶ Explanations of surface energies? Is it possible to extract surface energy contributions from a discrete Hamiltonian?
- ▶ Discrete elastic energies as regularizations of continuum energies? Comparability to singular perturbation problems?

The Square-to-Rectangular Phase Transition and Shape Memory Alloys

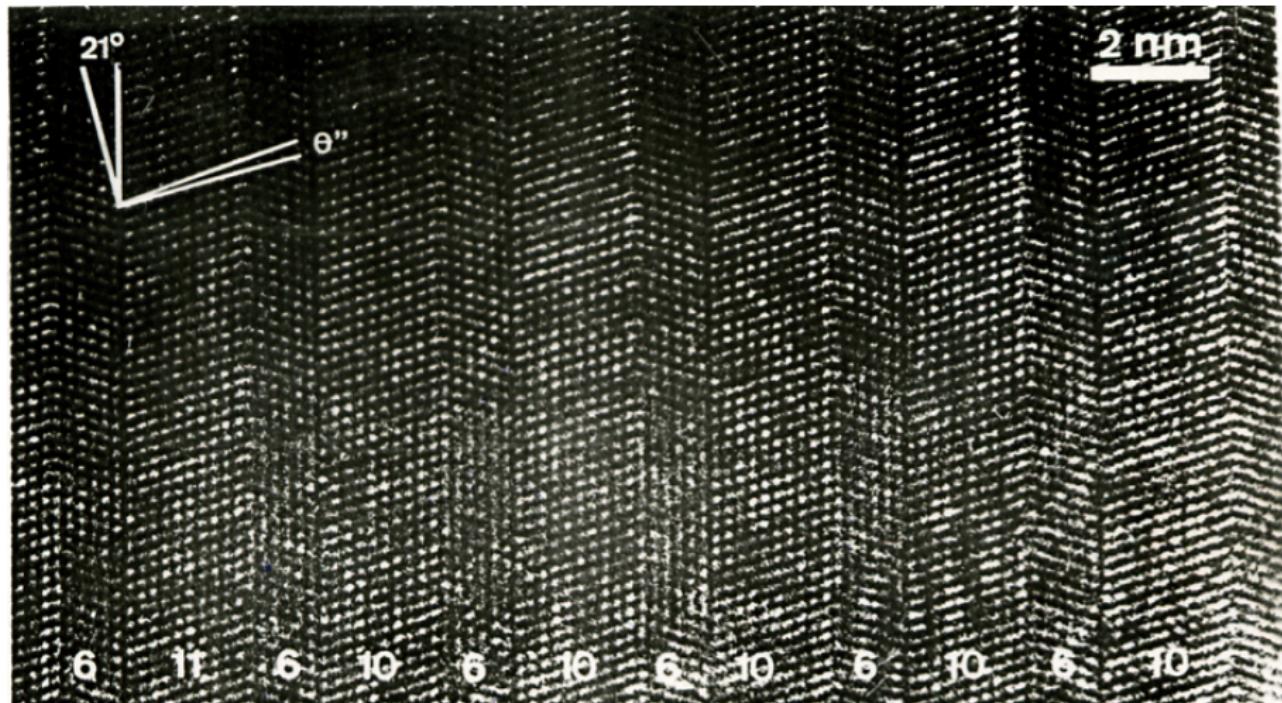


Experimental Observations: Diffuse Interfaces



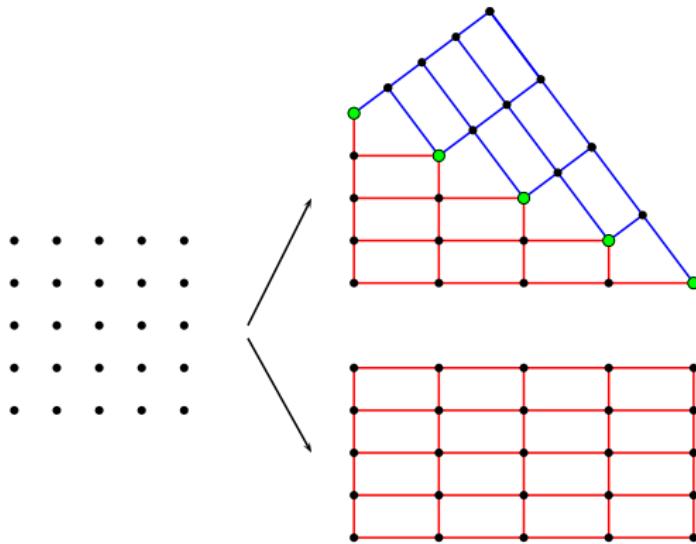
above: perovskite (Salje)
left: $Pb_3V_2O_8$ (Manolikas, van Tendeloo, Amelinckx)

Experimental Observations: Sharp Interfaces



NiMn (Baele, van Tendeloo, Amelinckx)

Deformations: Examples

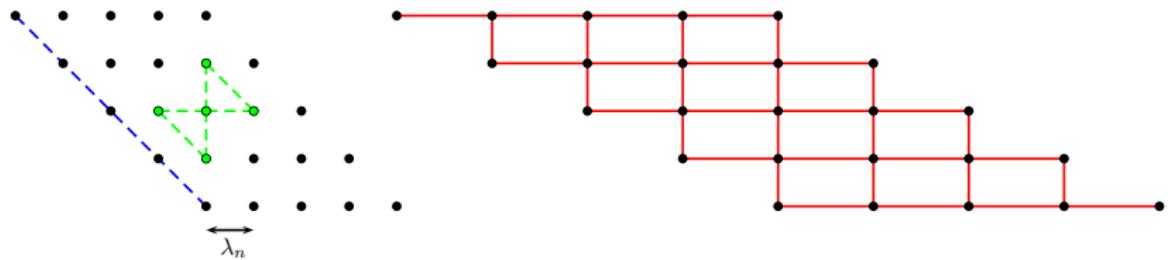


Homogeneous transformations to martensite are characterized by

- ▶ horizontal/vertical distances between neighboring atoms are given by a or b ,
- ▶ neighboring horizontal/vertical inter-atomic distances are equal,
- ▶ angles of 90° .

interfaces between the martensitic variants \rightsquigarrow violation of these properties.

Set-up



$$\Omega_n := \left\{ z \mid z = s \begin{pmatrix} 1 \\ 0 \end{pmatrix} + t \frac{1}{\sqrt{2}} \begin{pmatrix} -1 \\ 1 \end{pmatrix}, s, t \in [-1, 1] \right\} \cap [\lambda_n \mathbb{Z}]^2,$$

$$u_n : \Omega_n \rightarrow \mathbb{R}^2, \quad u_n^{j, -n-j} = F \lambda_n \begin{pmatrix} j \\ -n-j \end{pmatrix},$$

$u_n \in \mathcal{A}_n := \{v : \Omega_n \rightarrow \mathbb{R}^2 \mid \det(v(x_2) - v(x_1), v(x_3) - v(x_1)) > 0$
 for all $x_1, x_2, x_3 \in \Omega_n$ such that $\text{diam}(x_1, x_2, x_3) = \sqrt{2}\lambda_n$
 and $\det(x_2 - x_1, x_3 - x_1) > 0\}$.

Construction of a Model Hamiltonian

$$\begin{aligned}
 H_n(u) &:= \sum_{i,j=-n}^n \lambda_n^2 h\left(\frac{u^{ij} - u^{i\pm 1j}}{\lambda_n}, \frac{u^{ij} - u^{ij\pm 1}}{\lambda_n}\right) \\
 &= \sum_{i,j=-n}^n \lambda_n^2 \left[\left(\left(\frac{u^{ij\pm 1} - u^{ij}}{\lambda_n} \right)^2 - a^2 \right)^2 + \left(\left(\frac{u^{i\pm 1j} - u^{ij}}{\lambda_n} \right)^2 - b^2 \right)^2 \right. \\
 &\quad \left. + \left(\left(\frac{u^{ij\pm 1} - u^{ij}}{\lambda_n} \right) \cdot \left(\frac{u^{i\pm 1j} - u^{ij}}{\lambda_n} \right) \right)^2 \right] \times \\
 &\quad \times \left[\left(\left(\frac{u^{ij\pm 1} - u^{ij}}{\lambda_n} \right)^2 - b^2 \right)^2 + \left(\left(\frac{u^{i\pm 1j} - u^{ij}}{\lambda_n} \right)^2 - a^2 \right)^2 \right. \\
 &\quad \left. + \left(\left(\frac{u^{ij\pm 1} - u^{ij}}{\lambda_n} \right) \cdot \left(\frac{u^{i\pm 1j} - u^{ij}}{\lambda_n} \right) \right)^2 \right].
 \end{aligned}$$

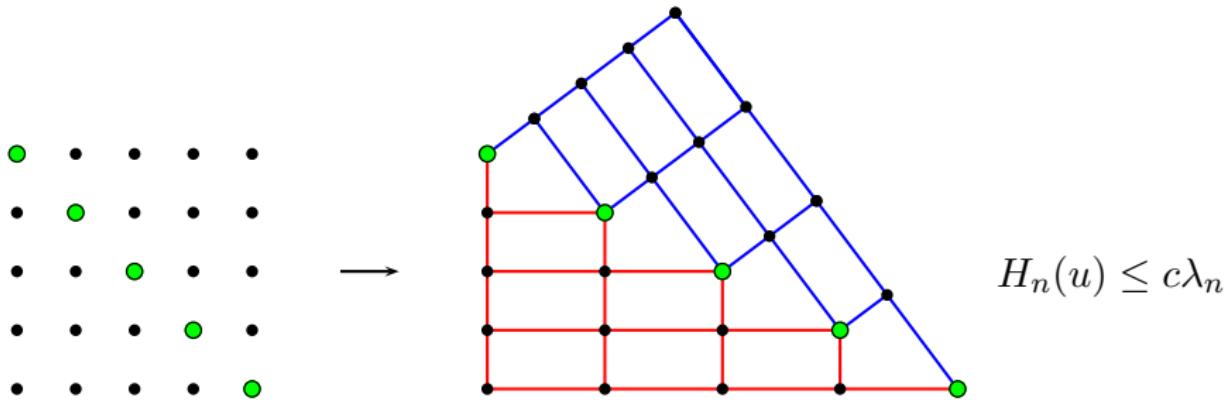
Construction of a Model Hamiltonian

$$H_n(u) := \sum_{i,j=-n}^n \lambda_n^2 h\left(\frac{u^{ij}-u^{i\pm 1j}}{\lambda_n}, \frac{u^{ij}-u^{ij\pm 1}}{\lambda_n}\right)$$

Advantages and Disadvantages

- ▶ based on geometric quantities,
- ▶ has $SO(2)U_0 \cup SO(2)U_1$ as wells,
- ▶ controls distance from wells,
- ▶ controls discrete second derivatives.
- ▶ ad hoc, no “first principles” justification for explicit form,
- ▶ uses underlying reference configuration,
- ▶ no defects allowed.

Martensitic Twins



$$U_0 - QU_1 = \sqrt{2} \frac{a^2 - b^2}{a^2 + b^2} \begin{pmatrix} a \\ -b \end{pmatrix} \otimes \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}, \quad Q \in SO(2),$$

$$U_0 - \tilde{Q}U_1 = \sqrt{2} \frac{a^2 - b^2}{a^2 + b^2} \begin{pmatrix} a \\ b \end{pmatrix} \otimes \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}, \quad \tilde{Q} \in SO(2).$$

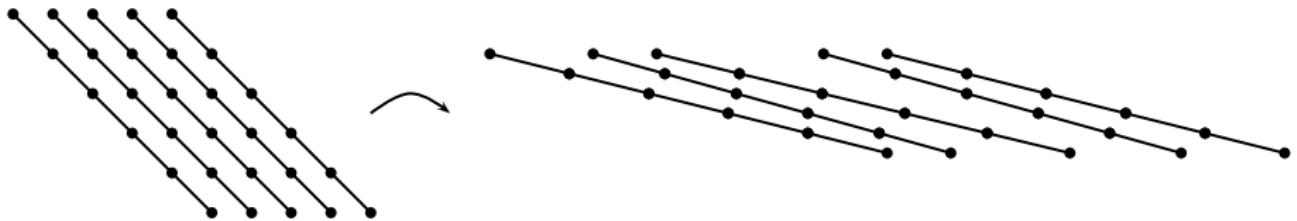
The Chain Hamiltonian

Additional “chain” assumption:

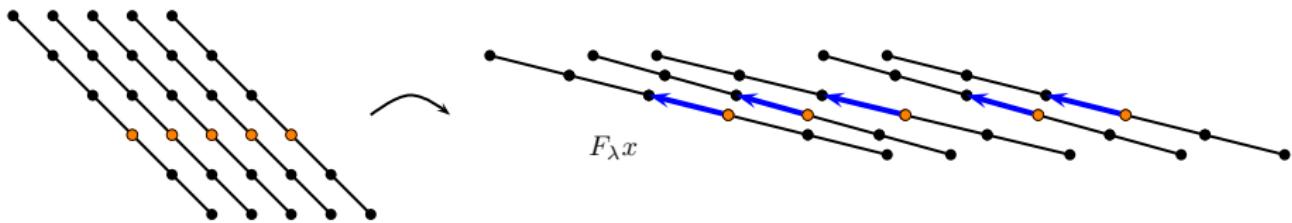
$$u_n^{i+1j} - u_n^{ij+1} = -\lambda_n \tau^{\textcolor{red}{i+j+1}}, \quad \tau^i \in SO(2) \begin{pmatrix} -a \\ b \end{pmatrix},$$

↔ corresponding adaptations for Hamiltonian

$$H_n(u_n) = \lambda_n^2 \sum_{i,j=-n}^n h \left(\frac{u_n^{\textcolor{blue}{i}} - u_n^{\textcolor{blue}{i}\pm 1}}{\lambda_n}, \tau_n^i, \tau_n^{i\pm 1}, j \right).$$



Set-up



$$H_n(u_n) := \sum_{i,j} \lambda_n^2 h(u^{i\pm 1} - u^i, \tau_n^i, \tau_n^{i\pm 1}, j) \leq C \lambda_n,$$

$$u_n^{j,-n-j} = \lambda_n F_\mu \binom{j}{-n-j}, \quad u_n \in \mathcal{A}_{n,\tau}^{F_\mu},$$

$$F_\mu = \mu Q U_0 + (1-\mu) U_1, \quad Q \in SO(2), \mu \in [0,1].$$

Rigidity

Proposition

Let $F_\mu := \mu U_0 + (1 - \mu) Q U_1$ with $\mu \in [0, 1]$. Let $\{u_n\}_{n \in \mathbb{N}} \in \mathcal{A}_{n,\tau}^{F_\mu}$ s.t.

$$\limsup_{n \rightarrow \infty} \lambda_n^{-1} H_n(u_n) < \infty.$$

Then there exists a number $K \in \mathbb{N}$ and a subsequence such that

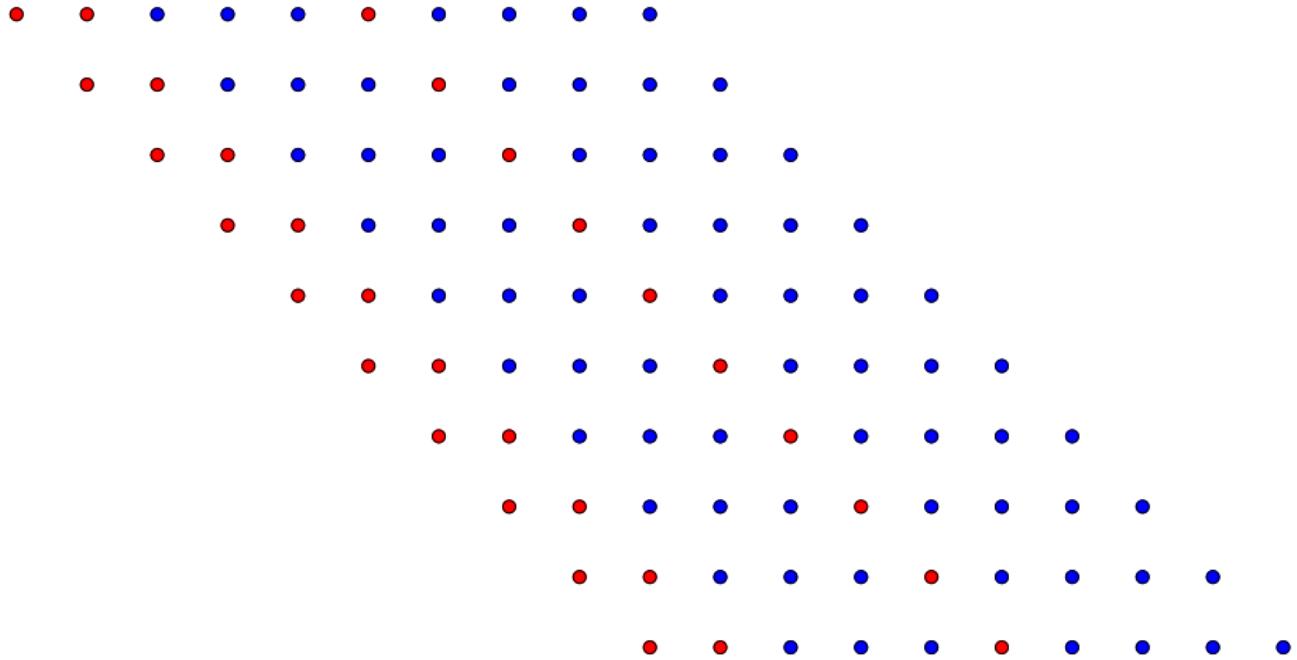
- ▶ $u_n \rightarrow u$ in $W^{1,4}(\Omega, \mathbb{R}^2)$,
- ▶ for each $s \in \{1, \dots, K-1\}$ there exists $m_s \in \{0, 1\}$, $x_s \in [-1, 1]$ such that

$$\nabla u(z) = Q^{m_s} U_{m_s},$$

for $z \in \Omega(x_s, x_{s+1})$ where $Q^0 = Id$, $Q^1 := Q$ and $x_K = 1$,

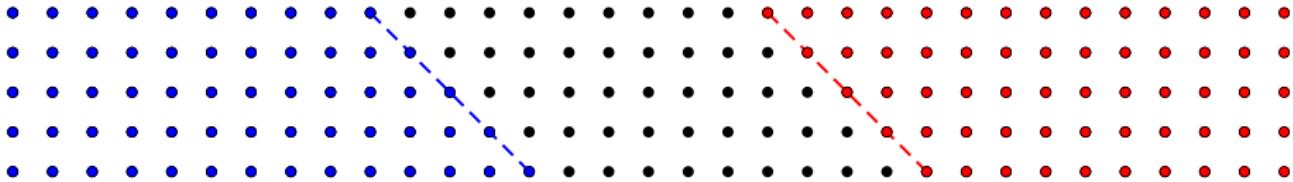
- ▶ $\bigcup_{s=1}^{K-1} [x_s, x_{s+1}] = [-1, 1]$.

Rigidity



Surface Energies

$$\begin{aligned}
 C(\textcolor{blue}{V}_2, \textcolor{red}{V}_3, r^*) := \liminf_{n \rightarrow \infty} \min_{\tau_i, u^i} & \Big\{ \sum_{i \in \mathbb{Z}} \frac{1}{n} \sum_{j=-n}^n h(u_n^i - u_n^{i \pm 1}, \tau_n^i, \tau_n^{i \pm 1}, j) : \\
 & u \in \mathcal{A}_{n, \tau}^r, \quad u^{i-jj} = \textcolor{blue}{V}_2 \binom{i-j}{j} + r_1, \quad i \leq -n, \quad |j| \leq n, \\
 & u^{i-jj} = \textcolor{red}{V}_3 \binom{i-j}{j} + r_2, \quad r^* = r_2 - r_1, \quad i \geq n, \quad |j| \leq n \Big\}.
 \end{aligned}$$



Surface Energies

Proposition

$H_n^1 := \lambda_n^{-1} H_n \xrightarrow{\Gamma} E_{surf}$ with respect to the L^∞ topology.

Here,

$$E_{surf}(u) := \begin{cases} E^K(F_\mu, \nabla u(x_1-, 0), \dots, \nabla u(x_{(K-1)-}, 0), F_\mu), \\ \quad \text{if } u \in W_0^{1,\infty}(\Omega) + F_\mu x, \quad \nabla u \in \{U_0, QU_1\} \\ \quad \text{in } \Omega(x_j, x_{j+1}), u \text{ satisfies the b.c.,} \\ \infty, \quad \text{else,} \end{cases}$$

$$\begin{aligned} E^K(V_0, \dots, V_K) := \inf_r & \left\{ B^+(V_0, V_1, r_0) + \sum_{s=1}^{K-2} C(V_s, V_{s+1}, r_s) \right. \\ & \left. + B^-(V_{K-1}, V_K, r_{K-1}) \right\}. \end{aligned}$$

Comparison with Literature

Discrete:

- ▶ Blanc, Le Bris, Lions (2002): From molecular models to continuum mechanics.
- ▶ Braides, Cicalese (2007): Surface energies in nonconvex discrete systems.
- ▶ Luckhaus, Mugnai (2009): On a mesoscopic many-body Hamiltonian describing elastic shears and dislocations.

Continuous:

- ▶ Conti, Schweizer (2006): Rigidity and Gamma Convergence for Solid-Solid Phase Transitions with $\text{SO}(2)$ Invariance.

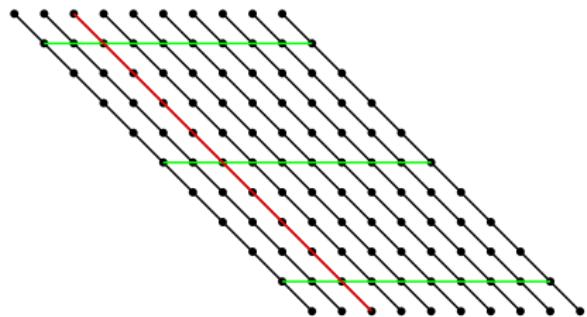
$$\sum_{i,j=-n}^n \frac{1}{n} h(u^{i\pm 1,j\pm 1} - u^{ij}) \leq C \quad \xrightarrow{\epsilon = \frac{1}{n}} \quad \int_{\Omega} \frac{1}{\epsilon} W(\nabla u) + \epsilon |\nabla^2 u|^2 dx \leq C$$

Compactness: Good Layers

There exist j_{-1}^n, j_0^n, j_1^n s.t.

- ▶ $j_{-1}^n \in [-n, -n + 2\delta n]$, $j_0^n \in [-\delta n, \delta n]$, $j_1^n \in [n - 2\delta n, n]$,
- ▶ $\lambda_n \sum_{i=-n}^n h\left(\frac{u_n^i - u_n^{i\pm 1}}{\lambda_n}, \tau_n^i, \tau_n^{i\pm 1}, j_l^n\right) \lesssim n^{-\alpha}$,
- ▶ $\#\left\{i \in [-n, n] : h\left(\frac{u_n^i - u_n^{i\pm 1}}{\lambda_n}, \tau_n^i, \tau_n^{i\pm 1}, j_l^n\right) \geq n^{-\alpha}\right\} \lesssim \delta^{-1} n^\alpha$,
- ▶ there exists a number $M_\delta > 0$, independent of n with

$$\#\left\{i \in [-n, n] : h\left(\frac{u_n^i - u_n^{i\pm 1}}{\lambda_n}, \tau_n^i, \tau_n^{i\pm 1}, j_l^n\right) \geq \tilde{c}\right\} \leq M_\delta.$$



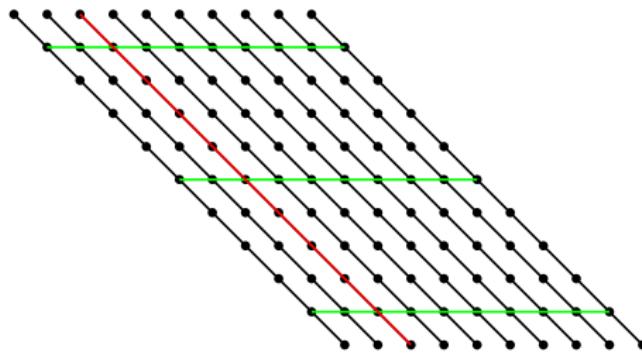
Chain structure:

- ▶ L^∞ bound:

$$|\nabla u_n^{ij}| \leq c < \infty$$

for all $i, j \in \{-n, \dots, n\}$.

Compactness: Simultaneously Good Points



Most points “simultaneously good” points:

$$h\left(\frac{u_n^i - u_n^{i\pm 1}}{\lambda_n}, \tau_n^i, \tau_n^{i\pm 1}, j_l^n\right) \leq n^{-\alpha}$$

for all $l \in \{-1, 0, 1\}$.

For each “simultaneously good” $i \in [-n, n]$ there exists $Q_{i,n} \in SO(2)$ such that either

$$\|\nabla u_n^{i-jj} - Q_{i,n} U_0\|_{C(\Omega_{i-jj})} \lesssim n^{-\alpha/4} \quad \text{for all } j \in [-n, n],$$

$$\text{or } \|\nabla u_n^{i-jj} - Q_{i,n} U_1\|_{C(\Omega_{i-jj})} \lesssim n^{-\alpha/4} \quad \text{for all } j \in [-n, n].$$

Compactness: Conclusion

Up to subsequences,

- ▶ there exist $K \in \mathbb{N}$, $x_1, \dots, x_K \in (-1, 1)$ independent of n ,
- ▶ and for any n there exist associated points $x_1^n, \dots, x_K^n \in (-1, 1)$ and $y_{s,1}^n, \dots, y_{s,K_s^n} \in (x_s^n, x_{s+1}^n)$

such that

- ▶ $x_s^n \rightarrow x_s$, $s \in \{1, \dots, K\}$,
- ▶ $u_n \rightharpoonup u$ in $W^{1,4}(\Omega)$,
- ▶ in the interval (x_s^n, x_{s+1}^n) the following dichotomy holds: For each i with $\lambda_n i \in (y_{s,l}^n, y_{s,l+1}^n) \subset (x_s^n, x_{s+1}^n)$ and $l \in \{1, \dots, K_s^n\}$, either

$$\text{dist}(\nabla u_n^{i-jj}, SO(2)U_0) \lesssim n^{-\alpha/4} \quad \text{or} \quad \text{dist}(\nabla u_n^{i-jj}, SO(2)U_1) \lesssim n^{-\alpha/4}$$

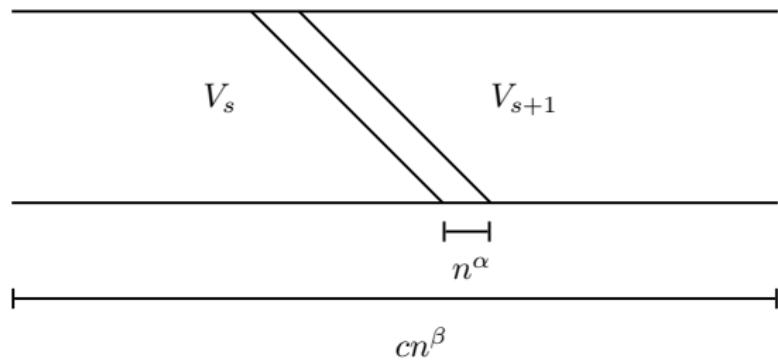
for all $j \in [-n, n]$.

Gamma-Limit

- ▶ **Idea:** Use infimizing sequences, modify boundary conditions.
- ▶ **Difficulties:** Ensure boundary conditions without violating admissibility (in particular non-interpenetration condition).

- ▶ **Techniques:**

- ▶ averaging,
- ▶ cutting.



The Full 2D Model and Further Questions

$$H_n(u) := \sum_{i,j=-n}^n \lambda_n^2 h\left(\frac{u^{ij}-u^{i\pm 1,j}}{\lambda_n}, \frac{u^{ij}-u^{i,j\pm 1}}{\lambda_n}\right)$$

Results for the full 2D model:

- ▶ Rigidity (one sided comparability to spin system).
- ▶ Sharp-interface limit:

$$H_n^1 := \lambda_n^{-1} H_n \xrightarrow{\Gamma} \tilde{E}_{surf}$$

with respect to the L^1 topology.

Further questions:

- ▶ More general m -well problem, e.g. three wells?
- ▶ Higher dimensional problem, e.g. 3D?
- ▶ Form of the energy densities?
- ▶ Minimizers of the energy densities? Relation to diffuse/sharp interfaces?