

**NONLINEAR STABILITY OF RELATIVISTIC VORTEX
SHEETS IN THREE-DIMENSIONAL MINKOWSKI
SPACETIME**

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Abstract We are concerned with the nonlinear stability of vortex sheets for the relativistic Euler equations in three-dimensional Minkowski spacetime. This is a nonlinear hyperbolic problem with a characteristic free boundary. In this paper, we introduce a new symmetrization by choosing appropriate functions as primary unknowns. A necessary and sufficient condition for the weakly linear stability of relativistic vortex sheets is obtained by analyzing the roots of the Lopatinskiĭ determinant associated to the constant coefficient linearized problem. Under this stability condition, we show that the variable coefficient linearized problem obeys an energy estimate with a loss of derivatives. The construction of certain weight functions plays a crucial role in absorbing the error terms caused by microlocalization. Based on the weakly linear stability result, we establish the existence and nonlinear stability of relativistic vortex sheets under small initial perturbations by a Nash–Moser iteration scheme.

Keywords Relativistic fluid · compressible vortex sheet · three-dimensional Minkowski spacetime · weak stability · loss of derivatives · characteristic boundary · Nash–Moser iteration · existence · linear stability · nonlinear stability

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Contents

1	Introduction	2
2	Nonlinear Problems and the Main Theorem	5
3	Constant Coefficient Linearized Problem	12
4	Variable Coefficient Linearized Problem	29
5	Basic Energy Estimate for the Linearized Problem	33
6	Well-posedness for the Linearized Problem	72
7	Construction of the Approximate Solution	79
8	Nash–Moser Iteration	82
	Appendix A Symmetrization of the Relativistic Euler Equations	99

1 Introduction

We are concerned with the nonlinear stability of relativistic vortex sheets for the Euler equations describing the evolution of a relativistic compressible fluid. Relativistic vortex sheets arise as a very important feature in several models of phenomena occurring in astrophysics, plasma physics, and nuclear physics. Vortex sheets are interfaces between two incompressible or compressible flows across which there is a discontinuity in fluid velocity. In particular, across a vortex sheet, the tangential velocity field has a jump, while the normal component of the flow velocity is continuous. The discontinuity in the tangential velocity field creates a concentration of vorticity along the interface. Moreover, compressible vortex sheets are characteristic discontinuities to the Euler equations for compressible fluids and as such they are fundamental waves which play an important role in the study of general entropy solutions to multidimensional hyperbolic systems of conservation laws (*cf.* Chen–Feldman [7]).

It was observed in [21, 36], by the normal mode analysis, that rectilinear vortex sheets for non-relativistic isentropic compressible fluids in two space dimensions are linearly stable when the Mach number $M > \sqrt{2}$ and are violently unstable when $M < \sqrt{2}$, while planar vortex sheets are always violently unstable in three space dimensions. This kind of instabilities is the analogue of the Kelvin–Helmholtz instability for incompressible fluids. Artola–Majda [3] studied certain instabilities of two-dimensional supersonic vortex sheets by analyzing the interaction with highly oscillatory waves through geometric optics. A rigorous mathematical theory on nonlinear stability and local-in-time existence of two-dimensional non-relativistic supersonic vortex sheets was first established by Coulombel–Secchi [19, 20] based on their linear stability results in [17] and a Nash–Moser iteration scheme.

Motivated by the earlier results in [17, 19, 20], we aim to establish the nonlinear stability of relativistic vortex sheets in three-dimensional Minkowski

spacetime under the necessary condition for the linear stability on the piecewise constant background state. This problem is a nonlinear hyperbolic problem with a characteristic free boundary. The so-called Lopatinskiĭ condition holds only in a weak sense, which yields a loss of derivatives.

We first reformulate the relativistic Euler equations into a symmetrizable hyperbolic system by choosing appropriate functions as primary unknowns. Our symmetrization is purely algebraic and different from those obtained by Makino–Ukai in [33] and Trakhinin [48]. As in Francheteau–Métivier [22], we straighten the unknown front by lifting functions Φ^\pm that satisfy the *eikonal* equations (2.23a) on the whole domain. Consequently, the original problem can be transformed into a nonlinear problem in a half-space for which the boundary matrix has constant rank on the whole half-space. This constant rank property is essential to derive energy estimates for the variable coefficient linearized problem by developing the Kreiss’ symmetrizers technique from [15, 17, 27].

Then we consider the constant coefficient linearized problem around the piecewise constant background state. By computing the roots of the associated Lopatinskiĭ determinant, we deduce the necessary stability condition (*cf.* (2.25)):

$$M > M_c := \frac{\sqrt{2}}{\sqrt{1 + \epsilon^2 \bar{c}^2}},$$

where ϵ^{-1} is the speed of light and \bar{c} is the sound speed of the background state. In the non-relativistic limit $\epsilon \rightarrow 0$, this stability condition is reduced to $M > \sqrt{2}$, the well-known fact studied in [17, 36]. The critical Mach number M_c of the relativistic stability condition is always strictly smaller than $\sqrt{2}$, which means that the relativistic vortex sheets are stable in a larger physical regime. Moreover, when the sound speed \bar{c} is arbitrarily close to the light speed ϵ^{-1} , the critical Mach number M_c approaches 1 so that the stability holds precisely for supersonic relativistic flows. The symbol associated to the unknown front is elliptic, which enables us to eliminate the front and to consider a standard boundary value problem. We prove that the constant coefficient linearized problem obeys an *a priori* energy estimate, which exhibits a loss of derivatives with respect to the source terms, owing to the failure of the uniform Kreiss–Lopatinskiĭ condition. Since the boundary is characteristic, there exists a loss of control on the trace of the solution.

After that, we study the effective linear problem, which is deduced from the linearized problem around a perturbation of the background state by using the “good unknown” of Alinhac [1] and neglecting some zero-th order terms. The dropped terms will be considered as the error terms at each Nash–Moser iteration step in the subsequent nonlinear analysis. We first prove for small perturbations that the solution satisfies the same *a priori* estimate as the constant coefficient case. The energy estimate is deduced by the technique applied earlier to weakly stable shock waves in [15] and isentropic compressible vortex sheets in [17]. It consists of the parilinearization of the linearized problem, analysis of the Lopatinskiĭ determinant, microlocalization, and construction of the Kreiss’ symmetrizers. In particular, we introduce certain weight func-

tions, vanishing *only* on the bicharacteristic curves starting from the critical set, to absorb the error terms caused by microlocalization. Based on this basic energy estimate, we establish a well-posedness result for the effective linear problem in the usual Sobolev space H^s with s large enough. This is achieved by means of a duality argument and higher order energy estimates. Although our problem is a hyperbolic problem with characteristic boundary that yields a natural loss of normal derivatives, we manage to compensate this loss by estimating missing normal derivatives through the equations of the linearized vorticity. With the well-posedness and tame estimate for the effective linear problem in hand, we prove the local existence theorem for relativistic vortex sheets (see Theorem 2.1) by a Nash–Moser iteration scheme. We emphasize that our choice of new primary unknowns is essential for three main reasons: The system becomes symmetrizable hyperbolic; it has an appropriate form for the analysis of the Lopatinskiĭ determinant; and, most of all, it is suitable for getting a vorticity-type equation.

Characteristic discontinuities, especially vortex sheets, arise in a broad range of physical problems in fluid mechanics, oceanography, aerodynamics, plasma physics, astrophysics, and elastodynamics. The linear results in [17] have been generalized to cover the two-dimensional nonisentropic flows [38], the three-dimensional compressible steady flows [50, 52], and the two-dimensional two-phase flows [42]. It is worth mentioning that a key ingredient in all of these proofs is the constant rank property of the boundary matrix. Recently, the methodology in [17] has been developed to deal with several constant coefficient linearized problems arising in two-dimensional compressible magnetohydrodynamics (MHD) and elastic flows; *cf.* [5, 10, 49]. See also the very recent preprint [11] for the linear stability of elastic vortex sheets in the variable coefficient case. For three-dimensional MHD, Chen–Wang [8, 9] and Trakhinin [47] proved independently the nonlinear stability of compressible current-vortex sheets, which indicates that non-paralleled magnetic fields stabilize the motion of three-dimensional compressible vortex sheets. Moreover, the modified Nash–Moser iteration scheme developed in [19, 24] has been successfully applied to the compressible liquids in vacuum [48], the plasma-vacuum interface problem [44], MHD contact discontinuities [39], and vortex sheets for three-dimensional steady flow [51] and two-dimensional two-phase flow [25].

Let us also mention some earlier works on the relativistic fluids. The global existence of discontinuous solutions to the relativistic Euler equations in one space dimension was first investigated by Smoller–Temple [45]. Also, Makino–Ukai [33] showed the existence of local smooth solutions in three space dimensions when the initial data is away from the vacuum. The stability of relativistic compressible flows with vacuum was addressed in [26, 48]. Moreover, the blow-up in finite time of smooth solutions for the relativistic Euler equations was shown in Pan–Smoller [40]. Also see Christodoulou [12, 13] for the formation and development of shocks in the multidimensional relativistic compressible fluids.

The plan of this paper is as follows: In § 2, after introducing the free boundary problem for relativistic vortex sheets, we reformulate the relativistic Euler equations and reduce our nonlinear problem to that in a fixed domain. Then we state the main result in this paper and introduce the weighted spaces and norms. Section 3 is mainly devoted to proving Theorem 3.1, *i.e.* an energy estimate for the constant coefficient linearized problem. More precisely, after some reductions, we compute the roots of the associated Lopatinskiĭ determinant and deduce the criterion for weakly linear stability in § 3.2. Then we adopt the argument developed recently by Chen–Hu–Wang [10] to prove the energy estimate for the constant coefficient case. In § 4, we introduce the effective linear problem and its reformulation. Section 5 is devoted to the proof of Theorem 5.1, the energy estimate for the effective linear problem. After deriving a weighted energy estimate with certain weights vanishing *only* on the bicharacteristic curves starting from the critical set, we can absorb the error terms caused by microlocalization and complete the proof of Theorem 5.1. In § 6, we prove a well-posedness result of the effective linear problem in the usual Sobolev space H^s with s large enough. In § 7, we obtain the smooth “approximate solution” by imposing necessary compatible conditions on the initial data. Then the original problem (2.20) and (2.23) is reduced into a nonlinear problem with zero initial data. In § 8, by using a modification of the Nash–Moser iteration scheme, we show the existence of solutions to the reduced problem and conclude the proof of our main result, Theorem 2.1. Appendix A concerns the motivation of introducing new primary unknowns and the derivation of the new symmetrization.

2 Nonlinear Problems and the Main Theorem

In this section, we first introduce the free boundary problem for relativistic vortex sheets, then reformulate the relativistic Euler equations and reduce our nonlinear problem to that in a fixed domain, and finally state the main theorem of this paper and introduce the weighted spaces and norms.

2.1 Relativistic Vortex Sheets

We consider the equations of relativistic perfect fluid dynamics in the three-dimensional Minkowski spacetime \mathbb{R}^{2+1} , that is, the relativistic Euler equations (see Lichnerowicz [30]):

$$\partial_\alpha T^{\alpha\beta} = 0, \quad (2.1)$$

where T denotes the energy-momentum stress tensor with components

$$T^{\alpha\beta} = (p + \rho\epsilon^{-2})u^\alpha u^\beta + pg^{\alpha\beta}.$$

Here p is the pressure, ρ is the energy-mass density, ϵ^{-1} is the speed of light, $g^{\alpha\beta} = \text{diag}(-1, 1, 1)$ is the flat Minkowski metric, and $\mathbf{u} = (u^0, u^1, u^2)^\top$ is the flow velocity satisfying

$$g^{\alpha\beta} u^\alpha u^\beta = -1. \quad (2.2)$$

The notation, ∂_α , denotes the differentiation with respect to variable x_α , and the Greek indices “ α ” and “ β ” run from 0 to 2. Throughout this paper, we use the Einstein summation convention whereby a repeated index in a term implies the summation over all the values of that index.

We introduce the coordinate velocity $v = (v_1, v_2)^\top := (u^1, u^2)^\top / (\epsilon u^0)$. By virtue of (2.2), the physical constraint is:

$$|v| < \epsilon^{-1}. \quad (2.3)$$

We also introduce the spacetime coordinates (t, x) with $t := \epsilon x_0$ and $x := (x_1, x_2)$. Then system (2.1) can be equivalently rewritten as

$$\partial_t ((\rho + \epsilon^2 p) \Gamma^2 - \epsilon^2 p) + \nabla_x \cdot ((\rho + \epsilon^2 p) \Gamma^2 v) = 0, \quad (2.4a)$$

$$\partial_t ((\rho + \epsilon^2 p) \Gamma^2 v) + \nabla_x \cdot ((\rho + \epsilon^2 p) \Gamma^2 v \otimes v) + \nabla_x p = 0, \quad (2.4b)$$

where $\partial_t = \frac{\partial}{\partial t}$, $\nabla_x = (\partial_1, \partial_2)^\top$ with $\partial_j = \frac{\partial}{\partial x_j}$, matrix $v \otimes v$ has (i, j) -entry $v_i v_j$, and

$$\Gamma = \Gamma(v) := \frac{1}{\sqrt{1 - \epsilon^2 |v|^2}} \quad (2.5)$$

is the Lorentz factor. The fluid is assumed to be barotropic, which means that pressure p is given by an explicit function of ρ . We also assume that $p = p(\rho)$ is a C^∞ function defined on (ρ_*, ρ^*) and satisfies

$$0 < p'(\rho) < \epsilon^{-2} \quad \text{for all } \rho \in (\rho_*, \rho^*), \quad (2.6)$$

where ρ_* and ρ^* are some constants such that $0 \leq \rho_* < \rho^* \leq \infty$. Consequently, density ρ is a strictly increasing function of p defined on $(p(\rho_*), p(\rho^*))$, and system (2.4) is closed with three unknowns (ρ, v_1, v_2) . Barotropic fluids arise in many physical situations such as very cold matter, nuclear matter and ultrarelativistic fluids (*cf.* [2, Chapter II], [14, Chapter IX], and [45, § 1]).

Let $(\rho, v)(t, x)$ be smooth functions on either side of a smooth hypersurface $\Sigma(t) := \{x_2 = \varphi(t, x_1)\}$. Then (ρ, v) is a weak solution of (2.4) if and only if (ρ, v) is a classical solution of (2.4) on each side of $\Sigma(t)$ and satisfies the Rankine–Hugoniot conditions at every point of $\Sigma(t)$:

$$\begin{cases} \partial_t \varphi [(\rho + \epsilon^2 p) \Gamma^2 - \epsilon^2 p] - [(\rho + \epsilon^2 p) \Gamma^2 v \cdot \nu] = 0, \\ \partial_t \varphi [(\rho + \epsilon^2 p) \Gamma^2 v] - [(\rho + \epsilon^2 p) \Gamma^2 (v \cdot \nu) v] - [p] \nu = 0, \end{cases} \quad (2.7)$$

where $\nu := (-\partial_1 \varphi, 1)$ is a spatial normal vector to $\Sigma(t)$. As usual, for any function g , we denote by g^\pm the value of g in $\{\pm(x_2 - \varphi(t, x_1)) > 0\}$, and $[g] := g^+|_{\Sigma(t)} - g^-|_{\Sigma(t)}$ the jump across $\Sigma(t)$.

In this paper, we are interested in weak solutions (ρ, v) of (2.4) such that the tangential velocity (with respect to $\Sigma(t)$) is the only jump experienced by the solution (ρ, v) across $\Sigma(t)$. Then the Rankine–Hugoniot conditions (2.7) are reduced to

$$\partial_t \varphi = v^+ \cdot \nu = v^- \cdot \nu, \quad p^+ = p^- \quad \text{on } \Sigma(t). \quad (2.8)$$

A piecewise smooth weak solution (ρ, v) of (2.4) with discontinuities across $\Sigma(t)$ is called a *relativistic vortex sheet* if its trace on $\Sigma(t)$ satisfies (2.8).

We note that system (2.4) admits trivial vortex-sheet solutions that consist of two constant states separated by a rectilinear front:

$$(\rho, v)(t, x_1, x_2) = \begin{cases} (\bar{\rho}, \bar{v}, 0) & \text{if } x_2 > 0, \\ (\bar{\rho}, -\bar{v}, 0) & \text{if } x_2 < 0, \end{cases} \quad (2.9)$$

where $\bar{\rho}$ and \bar{v} are suitable positive constants. Every rectilinear relativistic vortex sheet is of this form by changing the observer if necessary. In view of (2.3) and (2.6), we may assume without loss of generality that

$$\bar{\rho} \in (\rho_*, \rho^*), \quad \bar{v} \in (0, \epsilon^{-1}). \quad (2.10)$$

The aim of this paper is to study the local-in-time existence and nonlinear stability of relativistic vortex sheets with initial data close to the piecewise constant state (2.9).

2.2 Reformulation and the Main Theorem

Let us first reformulate the relativistic Euler equations (2.4) by choosing appropriate functions as primary unknowns. To this end, we define the particle number density $N = N(\rho)$, the sound speed $c = c(\rho)$, and $h = h(\rho)$ by

$$N(\rho) := \exp\left(\int_{\bar{\rho}}^{\rho} \frac{ds}{s + \epsilon^2 p(s)}\right), \quad c(\rho) := \sqrt{p'(\rho)}, \quad h(\rho) := \frac{\rho + \epsilon^2 p(\rho)}{N(\rho)}. \quad (2.11)$$

We also introduce

$$w := \Gamma v = \frac{v}{\sqrt{1 - \epsilon^2 |v|^2}}, \quad (2.12)$$

so that

$$\Gamma = \sqrt{1 + \epsilon^2 |w|^2}, \quad v = \frac{w}{\sqrt{1 + \epsilon^2 |w|^2}}. \quad (2.13)$$

Then we discover that smooth solutions to system (2.4) satisfy

$$\Gamma(\partial_t + v \cdot \nabla_x)(hw) + N^{-1} \nabla_x p = 0. \quad (2.14)$$

Let us take $U := (p, hw_1, hw_2)^\top$ as primary unknowns and define the following matrices:

$$A_0(U) := \begin{pmatrix} \Gamma(1 - \epsilon^4 c^2 |v|^2) & 2\epsilon^2 N c^2 v_1 & 2\epsilon^2 N c^2 v_2 \\ 0 & \Gamma(1 - \epsilon^2 v_1^2) & -\epsilon^2 \Gamma v_1 v_2 \\ 0 & -\epsilon^2 \Gamma v_1 v_2 & \Gamma(1 - \epsilon^2 v_2^2) \end{pmatrix}, \quad (2.15)$$

$$A_1(U) := \begin{pmatrix} \Gamma v_1(1 - \epsilon^4 c^2 |v|^2) & N c^2(1 + \epsilon^2 v_1^2) & \epsilon^2 v_1 v_2 N c^2 \\ N^{-1}(1 - \epsilon^2 v_1^2) & \Gamma v_1(1 - \epsilon^2 v_1^2) & -\epsilon^2 \Gamma v_1^2 v_2 \\ -\epsilon^2 v_1 v_2 N^{-1} & -\epsilon^2 \Gamma v_1^2 v_2 & \Gamma v_1(1 - \epsilon^2 v_2^2) \end{pmatrix}, \quad (2.16)$$

$$A_2(U) := \begin{pmatrix} \Gamma v_2(1 - \epsilon^4 c^2 |v|^2) & \epsilon^2 v_1 v_2 N c^2 & N c^2(1 + \epsilon^2 v_2^2) \\ -\epsilon^2 v_1 v_2 N^{-1} & \Gamma v_2(1 - \epsilon^2 v_1^2) & -\epsilon^2 \Gamma v_1 v_2^2 \\ N^{-1}(1 - \epsilon^2 v_2^2) & -\epsilon^2 \Gamma v_1 v_2^2 & \Gamma v_2(1 - \epsilon^2 v_2^2) \end{pmatrix}. \quad (2.17)$$

When the solution is in C^1 , system (2.4) equivalently reads

$$A_0(U)\partial_t U + A_1(U)\partial_1 U + A_2(U)\partial_2 U = 0. \quad (2.18)$$

We postpone proving the equivalence of systems (2.4) and (2.18) to Appendix A. The choice of the new unknowns U has several advantages. First, system (2.18) is symmetrizable hyperbolic in region $\{\rho_* < \rho < \rho^*, |v| < \epsilon^{-1}\}$ (see Appendix A for the precise expression of the Friedrichs symmetrizer). Second, we will see in the sequel that the form of (2.18) is appropriate for computing the roots of the Lopatinskiĭ determinant. Third, equations (2.14) will enable us to obtain the linearized vorticity equation through which the loss of derivatives can be compensated in the higher-order energy estimates.

Note that the first two identities in (2.8) are the *eikonal* equations:

$$\partial_t \varphi + \lambda_2(U^+, \partial_1 \varphi) = 0, \quad \partial_t \varphi + \lambda_2(U^-, \partial_1 \varphi) = 0,$$

where $\lambda_2(U, \xi) := v_1 \xi - v_2$, and $\xi \in \mathbb{R}$ is the second characteristic field of (2.18) with the corresponding eigenvector:

$$r_2(U, \xi) := (0, 1 - \epsilon^2 v_2^2 + \epsilon^2 v_1 v_2 \xi, (1 - \epsilon^2 v_1^2)\xi + \epsilon^2 v_1 v_2)^\top.$$

It follows from (2.11) and (2.13) that $\nabla_U \lambda_2(U, \xi) \cdot r_2(U, \xi) \equiv 0$, *i.e.* the characteristic field λ_2 is linearly degenerate in the sense of Lax [28]. As a consequence, a relativistic vortex sheet is a characteristic discontinuity.

Function φ describing the discontinuity front is a part of the unknowns, and thus the relativistic vortex sheet problem is a free boundary problem. To reformulate this problem in a fixed domain, we replace unknowns U , which are smooth on either side of $\Sigma(t)$, by

$$U_{\sharp}^\pm(t, x) := U(t, x_1, \Phi^\pm(t, x)), \quad (2.19)$$

where Φ^\pm are smooth functions satisfying the constraints:

$$\Phi^\pm(t, x_1, 0) = \varphi(t, x_1), \quad \pm \partial_2 \Phi^\pm(t, x) \geq \kappa > 0 \quad \text{if } x_2 \geq 0.$$

Then the existence of relativistic vortex sheets amounts to constructing solutions U_{\sharp}^{\pm} , which are smooth in the fixed domain $\{x_2 > 0\}$, to the following initial-boundary value problem:

$$\mathbb{L}(U^{\pm}, \Phi^{\pm}) = 0 \quad \text{if } x_2 > 0, \quad (2.20a)$$

$$\mathbb{B}(U^+, U^-, \varphi) = 0 \quad \text{if } x_2 = 0, \quad (2.20b)$$

$$(U^{\pm}, \varphi)|_{t=0} = (U_0^{\pm}, \varphi_0), \quad (2.20c)$$

where index “ \sharp ” has been dropped for notational simplicity. According to transformation (2.19), operators \mathbb{L} and \mathbb{B} take the forms:

$$\begin{cases} \mathbb{L}(U, \Phi) = L(U, \Phi)U \\ \text{with } L(U, \Phi) := A_0(U)\partial_t + A_1(U)\partial_1 + \tilde{A}_2(U, \Phi)\partial_2, \end{cases} \quad (2.21)$$

$$\mathbb{B}(U^+, U^-, \varphi) := \begin{pmatrix} [v_1]\partial_1\varphi - [v_2] \\ \partial_t\varphi + v_1^+\partial_1\varphi - v_2^+ \\ p^+ - p^- \end{pmatrix}, \quad (2.22)$$

where $A_j(U)$, $j = 0, 1, 2$, are defined by (2.15), (2.16), (2.17), respectively, and

$$\tilde{A}_2(U, \Phi) := \frac{1}{\partial_2\Phi} (A_2(U) - \partial_t\Phi A_0(U) - \partial_1\Phi A_1(U)).$$

As in Francheteau–Métivier [22], we choose the change of variables Φ^{\pm} such that

$$\partial_t\Phi^{\pm} + v_1^{\pm}\partial_1\Phi^{\pm} - v_2^{\pm} = 0 \quad \text{if } x_2 \geq 0, \quad (2.23a)$$

$$\pm \partial_2\Phi^{\pm} \geq \kappa > 0 \quad \text{if } x_2 \geq 0, \quad (2.23b)$$

$$\Phi^+ = \Phi^- = \varphi \quad \text{if } x_2 = 0. \quad (2.23c)$$

Not only does this choice simplify much the expression of system (2.20a), but it also implies that the boundary matrix for problem (2.20):

$$\text{diag}(-\tilde{A}_2(U^+, \Phi^+), -\tilde{A}_2(U^-, \Phi^-)),$$

has constant rank on the whole closed half-space $\{x_2 \geq 0\}$. This will play a crucial role in deriving the energy estimates for the variable coefficient linearized problem by developing further the Kreiss’ symmetrizers argument from [15, 17, 27].

In the new variables, the rectilinear vortex sheet (2.9) corresponds to the following stationary solution of (2.20a)–(2.20b) and (2.23):

$$\bar{U}^{\pm} := (\bar{p}, \pm \bar{h}\bar{w}, 0)^{\top}, \quad \bar{\varphi} := 0, \quad \bar{\Phi}^{\pm} := \pm x_2, \quad (2.24)$$

where $\bar{p} := p(\bar{\rho})$, $\bar{h} := h(\bar{\rho})$, and $\bar{w} := \bar{\Gamma}\bar{v}$ with $\bar{\Gamma}^{-1} := \sqrt{1 - \epsilon^2\bar{v}^2}$.

Imposing the smooth initial data (U_0^{\pm}, φ_0) close to (2.24), we aim to show the existence of solutions to the nonlinear problem (2.20) and (2.23) under the necessary condition for the linear stability on the background state (2.24). The main result is stated as follows:

Theorem 2.1 *Let $T > 0$ be any fixed constant and $\mu \in \mathbb{N}$ with $\mu \geq 13$. Assume that the background state (2.24) satisfies the physical constraints (2.10) and the necessary stability condition:*

$$M := \frac{|\bar{v}|}{\bar{c}} > \frac{\sqrt{2}}{\sqrt{1 + \epsilon^2 \bar{c}^2}} \quad \text{with} \quad \bar{c} := c(\bar{\rho}). \quad (2.25)$$

Assume further that the initial data U_0^\pm and φ_0 satisfy the compatibility conditions up to order μ (see § 7), and that $(U_0^\pm - \bar{U}^\pm, \varphi_0) \in H^{\mu+1/2}(\mathbb{R}_+^2) \times H^{\mu+1}(\mathbb{R})$ has a compact support. Then there exists a positive constant ε such that, if $\|U_0^\pm - \bar{U}^\pm\|_{H^{\mu+1/2}(\mathbb{R}_+^2)} + \|\varphi_0\|_{H^{\mu+1}(\mathbb{R})} \leq \varepsilon$, problem (2.20) and (2.23) has a solution $(U^\pm, \Phi^\pm, \varphi)$ on the time interval $[0, T]$ satisfying

$$(U^\pm - \bar{U}^\pm, \Phi^\pm - \bar{\Phi}^\pm) \in H^{\mu-7}((0, T) \times \mathbb{R}_+^2), \quad \varphi \in H^{\mu-6}((0, T) \times \mathbb{R}).$$

Remark 2.1 In the non-relativistic limit $\epsilon \rightarrow 0$, from (2.25), one obtains the classical stability condition $M > \sqrt{2}$ for compressible vortex sheets. The critical Mach number

$$M_c := \frac{\sqrt{2}}{\sqrt{1 + \epsilon^2 \bar{c}^2}}$$

of the relativistic stability condition is always strictly smaller than $\sqrt{2}$, which means that the relativistic vortex sheets are stable in a larger physical regime of the parameters. When \bar{c} is arbitrarily close to the light speed ϵ^{-1} , the critical Mach number M_c approaches 1 so that the stability holds precisely for supersonic relativistic flows.

2.3 Weighted Sobolev Spaces and Norms

We are going to introduce certain weighted Sobolev spaces in order to prove Theorem 2.1. Let Ω denote the half-space $\{(t, x_1, x_2) \in \mathbb{R}^3 : x_2 > 0\}$. Boundary $\partial\Omega$ is identified to \mathbb{R}^2 . For all $s \in \mathbb{R}$ and $\gamma \geq 1$, the usual Sobolev space $H^s(\mathbb{R}^2)$ is equipped with the following norm:

$$\|v\|_{s,\gamma}^2 := \frac{1}{(2\pi)^2} \int_{\mathbb{R}^2} \lambda^{2s,\gamma}(\xi) |\hat{v}(\xi)|^2 d\xi, \quad \lambda^{s,\gamma}(\xi) := (\gamma^2 + |\xi|^2)^{\frac{s}{2}},$$

where \hat{v} is the Fourier transform of v . We equip space $L^2(\mathbb{R}_+; H^s(\mathbb{R}^2))$ with the norm:

$$\| \| v \| \|_{s,\gamma}^2 := \int_{\mathbb{R}_+} \|v(\cdot, x_2)\|_{s,\gamma}^2 dx_2.$$

We will abbreviate the usual norms of $L^2(\mathbb{R}^2)$ and $L^2(\Omega)$ as

$$\| \cdot \| := \| \cdot \|_{0,\gamma} \quad \text{and} \quad \| \| \cdot \| \| := \| \| \cdot \| \|_{0,\gamma}.$$

The scalar products in $L^2(\mathbb{R}^2)$ and $L^2(\Omega)$ are denoted as follows:

$$\langle a, b \rangle := \int_{\mathbb{R}^2} a(x) \overline{b(y)} \, dy, \quad \langle\langle a, b \rangle\rangle := \int_{\Omega} a(y) \overline{b(y)} \, dy,$$

where $\overline{b(y)}$ is the complex conjugation of $b(y)$.

For $s \in \mathbb{R}$ and $\gamma \geq 1$, we introduce the weighted Sobolev space $H_\gamma^s(\mathbb{R}^2)$ as

$$H_\gamma^s(\mathbb{R}^2) := \{u \in \mathcal{D}'(\mathbb{R}^2) : e^{-\gamma t} u(t, x_1) \in H^s(\mathbb{R}^2)\},$$

and its norm $\|u\|_{H_\gamma^s(\mathbb{R}^2)} := \|e^{-\gamma t} u\|_{s, \gamma}$. We write $L_\gamma^2(\mathbb{R}^2) := H_\gamma^0(\mathbb{R}^2)$ and $\|u\|_{L_\gamma^2(\mathbb{R}^2)} := \|e^{-\gamma t} u\|$.

We define $L^2(\mathbb{R}_+; H_\gamma^s(\mathbb{R}^2))$, briefly denoted by $L^2(H_\gamma^s)$, as the space of distributions with finite $L^2(H_\gamma^s)$ -norm, where

$$\|u\|_{L^2(H_\gamma^s)}^2 := \int_{\mathbb{R}_+} \|u(\cdot, x_2)\|_{H_\gamma^s(\mathbb{R}^2)}^2 \, dx_2 = \|e^{-\gamma t} u\|_{s, \gamma}^2.$$

We set $L_\gamma^2(\Omega) := L^2(H_\gamma^0)$ and $\|u\|_{L_\gamma^2(\Omega)} := \|e^{-\gamma t} u\|$.

For all $k \in \mathbb{N}$ and $\gamma \geq 1$, we define the weighted Sobolev space $H_\gamma^k(\Omega)$ as

$$H_\gamma^k(\Omega) := \{u \in \mathcal{D}'(\Omega) : e^{-\gamma t} u \in H^k(\Omega)\}.$$

Throughout the paper, we introduce the notation: $A \lesssim B$ ($B \gtrsim A$) if $A \leq CB$ holds uniformly for some positive constant C that is *independent* of γ . The notation, $A \sim B$, means that both $A \lesssim B$ and $B \lesssim A$. Then, for $k \in \mathbb{N}$, one has

$$\|u\|_{k, \gamma} \sim \sum_{|\alpha| \leq k} \gamma^{k-|\alpha|} \|\partial^\alpha u\| \quad \text{for all } u \in H^k(\mathbb{R}^2). \quad (2.26)$$

For any real number T , we introduce $\omega_T := (-\infty, T) \times \mathbb{R}$ and $\Omega_T := \omega_T \times \mathbb{R}_+$. For all $k \in \mathbb{N}$ and $\gamma \geq 1$, we define the weighted space $H_\gamma^k(\Omega_T)$ as

$$H_\gamma^k(\Omega_T) := \{u \in \mathcal{D}'(\Omega_T) : e^{-\gamma t} u \in H^k(\Omega_T)\}.$$

In view of relation (2.26), we introduce the norm on $H_\gamma^k(\Omega_T)$ as

$$\|u\|_{H_\gamma^k(\Omega_T)} := \sum_{|\alpha| \leq k} \gamma^{k-|\alpha|} \|e^{-\gamma t} \partial^\alpha u\|_{L^2(\Omega_T)}. \quad (2.27)$$

The norm on $H_\gamma^k(\omega_T)$ is defined in the same way. For all $k \in \mathbb{N}$ and $\gamma \geq 1$, we define space $L^2(\mathbb{R}_+; H_\gamma^k(\omega_T))$, briefly denoted by $L^2(H_\gamma^k(\omega_T))$, as the space of distributions with finite $L^2(H_\gamma^k(\omega_T))$ -norm, where

$$\begin{aligned} \|u\|_{L^2(H_\gamma^k(\omega_T))}^2 &:= \int_{\mathbb{R}_+} \|u(\cdot, x_2)\|_{H_\gamma^k(\omega_T)}^2 \, dx_2 \\ &= \sum_{\alpha_0 + \alpha_1 \leq k} \gamma^{k-\alpha_0-\alpha_1} \|e^{-\gamma t} \partial_t^{\alpha_0} \partial_1^{\alpha_1} u\|_{L^2(\Omega_T)}. \end{aligned}$$

This is an anisotropic Sobolev space for measuring only the tangential regularity (with respect to boundary $\partial\Omega$). We write $L_\gamma^2(\Omega_T) := L^2(H_\gamma^0(\omega_T))$ and $\|u\|_{L_\gamma^2(\Omega_T)} := \|u\|_{L^2(H_\gamma^0(\omega_T))}$.

3 Constant Coefficient Linearized Problem

In order to deduce the necessary condition for the linear stability of the background state (2.24), in this section, we consider the following linearized problem of (2.20) and (2.23) around (2.24):

$$\mathbb{L}'_{\pm} V^{\pm} := \frac{d}{d\theta} \mathbb{L}(U_{\theta}^{\pm}, \Phi_{\theta}^{\pm}) \Big|_{\theta=0} = f^{\pm} \quad \text{if } x_2 > 0, \quad (3.1a)$$

$$\mathbb{B}'(V^+, V^-, \psi) := \frac{d}{d\theta} \mathbb{B}(U_{\theta}^+, U_{\theta}^-, \varphi_{\theta}) \Big|_{\theta=0} = g \quad \text{if } x_2 = 0, \quad (3.1b)$$

where $U_{\theta}^{\pm} := \bar{U}^{\pm} + \theta V^{\pm}$, $\Phi_{\theta}^{\pm} := \bar{\Phi}^{\pm} + \theta \Psi^{\pm}$, and φ_{θ} (resp. ψ) denotes the common trace of Φ_{θ}^{\pm} (resp. Ψ^{\pm}) on the boundary $\{x_2 = 0\}$. The differential operators \mathbb{L}'_{+} and \mathbb{L}'_{-} are given by

$$\mathbb{L}'_{\pm} := A_0(\bar{U}^{\pm}) \partial_t + A_1(\bar{U}^{\pm}) \partial_1 \pm A_2(\bar{U}^{\pm}) \partial_2,$$

which are both constant coefficient differential operators. It follows from (2.6) and (2.11) that

$$N(\bar{\rho}) = 1, \quad \bar{c} = c(\bar{\rho}) \in (0, \epsilon^{-1}). \quad (3.2)$$

To derive the boundary operator \mathbb{B}' , we infer from (2.13) that

$$\Gamma(U) = \frac{\sqrt{h(U_1)^2 + \epsilon^2 U_2^2 + \epsilon^2 U_3^2}}{h(U_1)}, \quad v_j(U) = \frac{U_{j+1}}{\Gamma(U)h(U_1)}, \quad j = 1, 2. \quad (3.3)$$

Utilizing the identity: $h'(U_1) = \epsilon^2/N(U_1)$ yields

$$\frac{\partial v_j}{\partial U_1} = -\frac{\epsilon^2 v_j}{Nh\Gamma^2}, \quad \frac{\partial v_j}{\partial U_{j+1}} = \frac{1 - \epsilon^2 v_j^2}{h\Gamma}, \quad \frac{\partial v_2}{\partial U_2} = \frac{\partial v_1}{\partial U_3} = -\frac{\epsilon^2 v_1 v_2}{h\Gamma}, \quad (3.4)$$

for $j = 1, 2$. The second component of $\mathbb{B}(U_{\theta}^+, U_{\theta}^-, \varphi_{\theta})$ is $\theta(\partial_t \psi + v_1(U_{\theta}^+) \partial_1 \psi) - v_2(U_{\theta}^+)$. Then we use (3.4) to obtain

$$\frac{d}{d\theta} (\mathbb{B}(U_{\theta}^+, U_{\theta}^-, \varphi_{\theta}))_2 \Big|_{\theta=0} = \partial_t \psi + \bar{v}^+ \partial_1 \psi - (\bar{h}\bar{\Gamma})^{-1} V_3^+.$$

After similar argument to the other components of $\mathbb{B}(U_{\theta}^+, U_{\theta}^-, \varphi_{\theta})$, we have

$$\mathbb{B}'(V^+, V^-, \psi) = \mathbb{B}'(V^{\text{nc}}, \psi) := \bar{b} \nabla \psi + \hat{B} V^{\text{nc}},$$

where $\nabla \psi := (\partial_t \psi, \partial_1 \psi)^{\text{T}}$, $V^{\text{nc}} := (V_1^+, V_3^+, V_1^-, V_3^-)^{\text{T}}$ denotes the ‘‘noncharacteristic part’’ of $V := (V^+, V^-)^{\text{T}}$, and coefficients \bar{b} and \hat{B} are given by

$$\bar{b} := \begin{pmatrix} 0 & 2\bar{v} \\ 1 & \bar{v} \\ 0 & 0 \end{pmatrix}, \quad \hat{B} := \begin{pmatrix} 0 & -(\bar{h}\bar{\Gamma})^{-1} & 0 & (\bar{h}\bar{\Gamma})^{-1} \\ 0 & -(\bar{h}\bar{\Gamma})^{-1} & 0 & 0 \\ 1 & 0 & -1 & 0 \end{pmatrix}. \quad (3.5)$$

We are now ready to state the main result for the constant coefficient case.

Theorem 3.1 *Assume that the stationary solution defined by (2.24) satisfies (2.10) and (2.25). Then, for all $\gamma \geq 1$ and for all $(V, \psi) \in H_\gamma^2(\Omega) \times H_\gamma^2(\mathbb{R}^2)$, the following estimate holds:*

$$\begin{aligned} & \gamma \|V\|_{L_\gamma^2(\Omega)}^2 + \|V^{\text{nc}}|_{x_2=0}\|_{L_\gamma^2(\mathbb{R}^2)}^2 + \|\psi\|_{H_\gamma^1(\mathbb{R}^2)}^2 \\ & \lesssim \gamma^{-3} \|\mathbb{L}'_\pm V^\pm\|_{L^2(H_\gamma^1)}^2 + \gamma^{-2} \|\mathbb{B}'(V^{\text{nc}}|_{x_2=0}, \psi)\|_{H_\gamma^1(\mathbb{R}^2)}^2. \end{aligned} \quad (3.6)$$

Remark 3.1 In the case of $M < M_c := \frac{\sqrt{2}}{\sqrt{1+\epsilon^2\bar{c}^2}}$, the relativistic vortex sheet (2.24) is violently unstable, *i.e.* the Lopatinskiĭ determinant admits the roots in the interior of frequency space. On the other hand, when $M \geq M_c$, all the roots of the Lopatinskiĭ determinant are localized on the boundary of frequency space. In particular, if $M = M_c$, the only root of the Lopatinskiĭ determinant is a triple one, which leads to the following weaker estimate than (3.6):

$$\begin{aligned} & \gamma \|V\|_{L_\gamma^2(\Omega)}^2 + \|V^{\text{nc}}|_{x_2=0}\|_{L_\gamma^2(\mathbb{R}^2)}^2 + \|\psi\|_{H_\gamma^1(\mathbb{R}^2)}^2 \\ & \lesssim \gamma^{-7} \|\mathbb{L}'_\pm V^\pm\|_{L^2(H_\gamma^3)}^2 + \gamma^{-6} \|\mathbb{B}'(V^{\text{nc}}|_{x_2=0}, \psi)\|_{H_\gamma^3(\mathbb{R}^2)}^2, \end{aligned} \quad (3.7)$$

for all $\gamma \geq 1$ and $(V, \psi) \in H_\gamma^4(\Omega) \times H_\gamma^4(\mathbb{R}^2)$. See Remarks 3.2–3.4 for more details. This latter case corresponds to a transition between a weakly stable zone and a violently unstable zone (*cf.* Coulombel–Secchi [18] for the non-relativistic case).

The rest of this section is devoted mainly to proving Theorem 3.1.

3.1 Some Reductions

Before proving Theorem 3.1, we first make some reductions of problem (3.1).

3.1.1 Reformulation of Theorem 3.1

We first transform our problem (3.1) into the one with diagonal boundary matrix. For this purpose, we calculate the eigenvalues and corresponding eigenvectors of $A_2(\bar{U}^\pm)$. The eigenvalues of

$$A_2(\bar{U}^+) = A_2(\bar{U}^-) = \begin{pmatrix} 0 & 0 & \bar{c}^2 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}$$

are $\lambda_1 = 0$, $\lambda_2 = -\bar{c}$, and $\lambda_3 = \bar{c}$, with the corresponding right eigenvectors:

$$r_1 = (0, 1, 0)^\top, \quad r_2 = (1, 0, -\frac{1}{\bar{c}})^\top, \quad r_3 = (1, 0, \frac{1}{\bar{c}})^\top.$$

Set $\bar{R} := (r_1 \ r_2 \ r_3)$. Then $\bar{R}^{-1}A_2(\bar{U}^\pm)\bar{R} = \text{diag}(0, -\bar{c}, \bar{c})$. We thus perform the linear transformation $W^\pm := \bar{R}^{-1}V^\pm$ with

$$W_1^\pm := V_2^\pm, \quad W_2^\pm := \frac{1}{2}(V_1^\pm - \bar{c}V_3^\pm), \quad W_3^\pm := \frac{1}{2}(V_1^\pm + \bar{c}V_3^\pm).$$

Let us multiply (3.1a) by $\bar{S}\bar{R}^{-1}$ with $\bar{S} := \text{diag}(1, 2/\bar{c}^2, 2/\bar{c}^2)$. Then problem (3.1) becomes equivalent to

$$\begin{cases} \mathcal{L}W := \mathcal{A}_0\partial_t W + \mathcal{A}_1\partial_1 W + \mathcal{A}_2\partial_2 W = f & \text{if } x_2 > 0, \\ \mathcal{B}(W^{\text{nc}}, \psi) := \bar{b}\nabla\psi + \bar{B}W^{\text{nc}} = g & \text{if } x_2 = 0, \end{cases} \quad (3.8)$$

with new f and g , where $W^{\text{nc}} := (W_2^+, W_3^+, W_2^-, W_3^-)^\top$ denotes the ‘‘non-characteristic part’’ of $W := (W^+, W^-)^\top$. The coefficient matrices $\mathcal{A}_j = \text{diag}(\mathcal{A}_j^+, \mathcal{A}_j^-)$, $j = 0, 1, 2$, are block diagonal with

$$\begin{aligned} \mathcal{A}_0^\pm &:= \bar{S}\bar{R}^{-1}A_0(\bar{U}^\pm)\bar{R} \\ &= \begin{pmatrix} \bar{\Gamma}(1 - \epsilon^2\bar{v}^2) & 0 & 0 \\ \pm 2\epsilon^2\bar{v} & \frac{\bar{\Gamma}(2 - \epsilon^4\bar{c}^2\bar{v}^2)}{\bar{c}^2} & -\epsilon^4\bar{\Gamma}\bar{v}^2 \\ \pm 2\epsilon^2\bar{v} & -\epsilon^4\bar{\Gamma}\bar{v}^2 & \frac{\bar{\Gamma}(2 - \epsilon^4\bar{c}^2\bar{v}^2)}{\bar{c}^2} \end{pmatrix}, \end{aligned} \quad (3.9)$$

$$\begin{aligned} \mathcal{A}_1^\pm &:= \bar{S}\bar{R}^{-1}A_1(\bar{U}^\pm)\bar{R} \\ &= \begin{pmatrix} \pm\bar{\Gamma}(1 - \epsilon^2\bar{v}^2)\bar{v} & 1 - \epsilon^2\bar{v}^2 & 1 - \epsilon^2\bar{v}^2 \\ 1 + \epsilon^2\bar{v}^2 & \pm\frac{\bar{\Gamma}\bar{v}(2 - \epsilon^4\bar{c}^2\bar{v}^2)}{\bar{c}^2} & \mp\epsilon^4\bar{\Gamma}\bar{v}^3 \\ 1 + \epsilon^2\bar{v}^2 & \mp\epsilon^4\bar{\Gamma}\bar{v}^3 & \pm\frac{\bar{\Gamma}\bar{v}(2 - \epsilon^4\bar{c}^2\bar{v}^2)}{\bar{c}^2} \end{pmatrix}, \end{aligned} \quad (3.10)$$

and

$$\mathcal{A}_2^\pm := \pm\bar{S}\bar{R}^{-1}A_2(\bar{U}^\pm)\bar{R} = \pm\text{diag}(0, -\frac{2}{\bar{c}}, \frac{2}{\bar{c}}). \quad (3.11)$$

We notice that (2.18) is a symmetrizable hyperbolic system with the Friedrichs symmetrizer $S_2(U)$ defined in (A.10). Consequently, operator \mathcal{L} is symmetrizable hyperbolic with the Friedrichs symmetrizer S_3 defined by

$$S_3 := \text{diag}(\bar{R}^\top S_2(\bar{U}^+)\bar{R}\bar{S}^{-1}, \bar{R}^\top S_2(\bar{U}^-)\bar{R}\bar{S}^{-1}). \quad (3.12)$$

Regarding the boundary coefficients, \bar{b} is given in (3.5), and \bar{B} is defined by

$$\bar{B} := \begin{pmatrix} (\bar{\Gamma}\bar{c}\bar{h})^{-1} & -(\bar{\Gamma}\bar{c}\bar{h})^{-1} & -(\bar{\Gamma}\bar{c}\bar{h})^{-1} & (\bar{\Gamma}\bar{c}\bar{h})^{-1} \\ (\bar{\Gamma}\bar{c}\bar{h})^{-1} & -(\bar{\Gamma}\bar{c}\bar{h})^{-1} & 0 & 0 \\ 1 & 1 & -1 & -1 \end{pmatrix}. \quad (3.13)$$

For $\gamma \geq 1$, we define

$$\mathcal{L}^\gamma := \mathcal{L} + \gamma\mathcal{A}_0, \quad \mathcal{B}^\gamma(W^{\text{nc}}, \psi) := \bar{b} \begin{pmatrix} \gamma\psi + \partial_t\psi \\ \partial_1\psi \end{pmatrix} + \bar{B}W^{\text{nc}}.$$

It is easily shown that Theorem 3.1 admits the following equivalent proposition.

Proposition 3.1 *Assume that the stationary solution (2.24) satisfies (2.10) and (2.25). Then, for all $\gamma \geq 1$ and $(W, \psi) \in H^2(\Omega) \times H^2(\mathbb{R}^2)$, the following estimate holds:*

$$\gamma\|W\|^2 + \|W^{\text{nc}}|_{x_2=0}\|^2 + \|\psi\|_{1,\gamma}^2 \lesssim \gamma^{-3}\|\mathcal{L}^\gamma W\|_{1,\gamma}^2 + \gamma^{-2}\|\mathcal{B}^\gamma(W^{\text{nc}}|_{x_2=0}, \psi)\|_{1,\gamma}^2.$$

3.1.2 Partial homogenization

In order to prove Proposition 3.1, we show that it suffices to study the homogeneous case $\mathcal{L}^\gamma W \equiv 0$. Given $(W, \psi) \in H^2(\Omega) \times H^2(\mathbb{R}^2)$, we set

$$f := \mathcal{L}^\gamma W \in H^1(\Omega), \quad g := \mathcal{B}^\gamma(W^{\text{nc}}|_{x_2=0}, \psi) \in H^1(\mathbb{R}^2),$$

and consider the following auxiliary problem:

$$\begin{cases} \mathcal{L}^\gamma W_1 = f & \text{if } x_2 > 0, \\ B^{\text{aux}} W_1^{\text{nc}} = 0 & \text{if } x_2 = 0, \end{cases} \quad (3.14)$$

where

$$B^{\text{aux}} := \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}.$$

The boundary matrix for problem (3.14) (*i.e.* $-\mathcal{A}_2$) has two negative eigenvalues and is nonnegative on $\ker B^{\text{aux}} = \{W_3^+ = W_2^- = 0\}$. Thus, the boundary conditions in (3.14) are maximally dissipative. From Lax–Phillips [29], there exists a unique solution $W_1 \in L^2(\mathbb{R}_+; H^1(\mathbb{R}^2))$ to problem (3.14) such that the trace of W_1 on $\{x_2 = 0\}$ is in $H^1(\mathbb{R}^2)$, and

$$\gamma \| \| W_1 \| \|^2 \lesssim \gamma^{-1} \| \| f \| \|^2, \quad \| \| W_1^{\text{nc}} |_{x_2=0} \| \|^2_{1,\gamma} \lesssim \gamma^{-1} \| \| f \| \|^2_{1,\gamma}. \quad (3.15)$$

It is clear that $W_2 := W - W_1$ satisfies

$$\mathcal{L}^\gamma W_2 = 0 \quad \text{if } x_2 > 0, \quad (3.16a)$$

$$\mathcal{B}^\gamma(W_2^{\text{nc}}, \psi) = \tilde{g} \quad \text{if } x_2 = 0, \quad (3.16b)$$

where $\tilde{g} := g - \bar{B} W_1^{\text{nc}}$. By virtue of (3.15), we obtain

$$\| \| \tilde{g} \| \|^2_{1,\gamma} \lesssim \| \| g \| \|^2_{1,\gamma} + \gamma^{-1} \| \| f \| \|^2_{1,\gamma}. \quad (3.17)$$

Multiplying (3.16a) by the symmetrizer S_3 (*cf.* (3.12)), then taking the scalar product of the resulting equations with W_2 , and employing integration by parts yield

$$\gamma \| \| W_2 \| \|^2 \lesssim \| \| W_2^{\text{nc}} |_{x_2=0} \| \|^2. \quad (3.18)$$

The next lemma follows directly from (3.15) and (3.17)–(3.18).

Lemma 3.1 *If the solution of (3.16) satisfies the estimate:*

$$\| \| W_2^{\text{nc}} |_{x_2=0} \| \|^2 + \| \| \psi \| \|^2_{1,\gamma} \lesssim \gamma^{-2} \| \| \tilde{g} \| \|^2_{1,\gamma},$$

then Proposition 3.1 holds.

3.1.3 Eliminating the front

We perform the Fourier transform of problem (3.16) in (t, x_1) , with dual variables denoted by (δ, η) . Setting $\tau = \gamma + i\delta$, we have

$$(\tau \mathcal{A}_0 + i\eta \mathcal{A}_1) \widehat{W} + \mathcal{A}_2 \frac{d\widehat{W}}{dx_2} = 0 \quad \text{if } x_2 > 0, \quad (3.19a)$$

$$b(\tau, \eta) \widehat{\psi} + \overline{B} \widehat{W}^{\text{nc}} = \widehat{g} \quad \text{if } x_2 = 0, \quad (3.19b)$$

where we write g for \tilde{g} and W for W_2 for simplicity when no confusion arises. The coefficient:

$$b(\tau, \eta) := \bar{b}(\tau, i\eta)^\top = (2i\bar{v}\eta, \tau + i\bar{v}\eta, 0)^\top$$

is homogeneous of degree 1 in (τ, η) . In order to take this homogeneity into account, we define the hemisphere:

$$\Xi_1 := \{(\tau, \eta) \in \mathbb{C} \times \mathbb{R} : |\tau|^2 + \eta^2 = 1, \operatorname{Re} \tau \geq 0\},$$

and the set of ‘‘frequencies’’:

$$\Xi := \{(\tau, \eta) \in \mathbb{C} \times \mathbb{R} : \operatorname{Re} \tau \geq 0, (\tau, \eta) \neq (0, 0)\} = (0, \infty) \cdot \Xi_1.$$

Notice that symbol $b(\tau, \eta)$ is elliptic, *i.e.* it is always different from zero on Ξ_1 .

We set $k := \sqrt{|\tau|^2 + \eta^2}$, and define

$$Q(\tau, \eta) := \frac{1}{k} \begin{pmatrix} 0 & 0 & k \\ \tau + i\bar{v}\eta & -2i\bar{v}\eta & 0 \\ -2i\bar{v}\eta & \bar{\tau} - i\bar{v}\eta & 0 \end{pmatrix} \quad \text{for } (\tau, \eta) \in \Xi,$$

where $\bar{\tau}$ denotes the complex conjugation of τ , so that $Q \in C^\infty(\Xi, GL_3(\mathbb{C}))$ is homogeneous of degree 0 in (τ, η) and satisfies

$$Q(\tau, \eta) b(\tau, \eta) = (0, 0, \theta(\tau, \eta))^\top \quad \text{with } \theta(\tau, \eta) = k^{-1} |b(\tau, \eta)|^2.$$

Since $\bar{v} \neq 0$, and Ξ_1 is compact, we obtain that $\min_{(\tau, \eta) \in \Xi_1} |\theta(\tau, \eta)| > 0$. Multiplying (3.19b) by $Q(\tau, \eta)$ yields

$$\begin{pmatrix} 0 \\ 0 \\ \theta(\tau, \eta) \end{pmatrix} \widehat{\psi}(\delta, \eta) + \begin{pmatrix} \beta(\tau, \eta) \\ \ell(\tau, \eta) \end{pmatrix} \widehat{W}^{\text{nc}}(\delta, \eta, 0) = Q(\tau, \eta) \widehat{g}, \quad (3.20)$$

where β is the 2×4 matrix given by the first two rows of $Q(\tau, \eta) \bar{B}$, and ℓ is the last row of $Q(\tau, \eta) \bar{B}$. Both β and ℓ are C^∞ and homogeneous of degree 0 on Ξ . In view of (3.13), symbol β satisfies

$$\beta(\tau, \eta) = \begin{pmatrix} 1 & 1 & -1 & -1 \\ \frac{\tau - i\bar{v}\eta}{\bar{\Gamma}\bar{c}\bar{h}} & \frac{-\tau + i\bar{v}\eta}{\bar{\Gamma}\bar{c}\bar{h}} & \frac{-\tau - i\bar{v}\eta}{\bar{\Gamma}\bar{c}\bar{h}} & \frac{\tau + i\bar{v}\eta}{\bar{\Gamma}\bar{c}\bar{h}} \end{pmatrix} \quad \text{on } \Xi_1. \quad (3.21)$$

The last component in (3.20) reads

$$\theta(\tau, \eta)\widehat{\psi} + \ell(\tau, \eta)\widehat{W}^{\text{nc}}(\delta, \eta, 0) = Q_3(\tau, \eta)\widehat{g},$$

where $Q_3(\tau, \eta)$ is the last row of $Q(\tau, \eta)$. Hence, it is homogeneous of degree 0. Thanks to the homogeneity of θ and ℓ , we obtain

$$k^2|\widehat{\psi}|^2 \lesssim |\widehat{W}^{\text{nc}}|_{x_2=0}|^2 + |\widehat{g}|^2 \quad \text{in } \Xi,$$

from which we employ Plancherel's theorem to deduce

$$\|\psi\|_{1,\gamma}^2 \lesssim \|W^{\text{nc}}|_{x_2=0}\|^2 + \gamma^{-2}\|g\|_{1,\gamma}^2. \quad (3.22)$$

After eliminating the front function ψ , we have

$$(\tau\mathcal{A}_0 + i\eta\mathcal{A}_1)\widehat{W} + \mathcal{A}_2\frac{d\widehat{W}}{dx_2} = 0 \quad \text{if } x_2 > 0, \quad (3.23a)$$

$$\beta(\tau, \eta)\widehat{W}^{\text{nc}} = \widehat{G} \quad \text{if } x_2 = 0, \quad (3.23b)$$

where \widehat{G} consists of the first two rows of $Q(\tau, \eta)\widehat{g}$. From (3.9)–(3.10), we have

$$\begin{aligned} & \tau\mathcal{A}_0^\pm + i\eta\mathcal{A}_1^\pm \\ &= \begin{pmatrix} a_\pm\bar{\Gamma}(1 - \epsilon^2\bar{v}^2) & i\eta(1 - \epsilon^2\bar{v}^2) & i\eta(1 - \epsilon^2\bar{v}^2) \\ i\eta(1 + \epsilon^2\bar{v}^2) \pm 2\epsilon^2\bar{v}\tau & \frac{\bar{\Gamma}(2 - \epsilon^4\bar{c}^2\bar{v}^2)}{\bar{c}^2}a_\pm & -\epsilon^4\bar{\Gamma}\bar{v}^2a_\pm \\ i\eta(1 + \epsilon^2\bar{v}^2) \pm 2\epsilon^2\bar{v}\tau & -\epsilon^4\bar{\Gamma}\bar{v}^2a_\pm & \frac{\bar{\Gamma}(2 - \epsilon^4\bar{c}^2\bar{v}^2)}{\bar{c}^2}a_\pm \end{pmatrix}, \end{aligned} \quad (3.24)$$

where $a_\pm := \tau \pm i\bar{v}\eta$. Recalling that formula (3.11) defines the boundary matrix \mathcal{A}_2^\pm , we write the first and fourth equations of (3.23a) as

$$a_\pm\bar{\Gamma}\widehat{W}_1^\pm + i\eta\widehat{W}_2^\pm + i\eta\widehat{W}_3^\pm = 0. \quad (3.25)$$

Then we utilize (3.25) to express \widehat{W}_1^\pm in terms of \widehat{W}_2^\pm and \widehat{W}_3^\pm , and plug the resulting expressions into the other four equations of (3.23a). As a consequence, we obtain a system of ordinary differential equations for \widehat{W}^{nc} in the following form:

$$\begin{cases} \frac{d}{dx_2}\widehat{W}^{\text{nc}} = \mathcal{A}(\tau, \eta)\widehat{W}^{\text{nc}} & \text{if } x_2 > 0, \\ \beta(\tau, \eta)\widehat{W}^{\text{nc}} = \widehat{G} & \text{if } x_2 = 0. \end{cases} \quad (3.26)$$

Here matrix $\mathcal{A}(\tau, \eta)$ is given by

$$\mathcal{A}(\tau, \eta) := \begin{pmatrix} \mu_+ & -m_+ & 0 & 0 \\ m_+ & -\mu_+ & 0 & 0 \\ 0 & 0 & -\mu_- & m_- \\ 0 & 0 & -m_- & \mu_- \end{pmatrix}, \quad (3.27)$$

where

$$\mu_{\pm} := \frac{\bar{\Gamma}a_{\pm}}{\bar{c}} - m_{\pm}, \quad m_{\pm} := \frac{\epsilon^2 \bar{c} \bar{v}^2 \bar{\Gamma}^2 a_{\pm}^2 + i\eta \bar{c} (i\eta(1 + \epsilon^2 \bar{v}^2) \pm 2\epsilon^2 \bar{v} \tau)}{2\bar{\Gamma}a_{\pm}}.$$

Using the relation, $\bar{\Gamma}^{-2} = 1 - \epsilon^2 \bar{v}^2$, yields

$$\mu_{\pm} = \frac{\bar{\Gamma}(\tau \pm i\bar{v}\eta)}{\bar{c}} - \frac{\bar{c}\bar{\Gamma}(i\eta \pm \epsilon^2 \bar{v} \tau)^2}{2(\tau \pm i\bar{v}\eta)}, \quad m_{\pm} = \frac{\bar{c}\bar{\Gamma}(i\eta \pm \epsilon^2 \bar{v} \tau)^2}{2(\tau \pm i\bar{v}\eta)}. \quad (3.28)$$

The reader may recognize the form of the symbol in (3.27) given also by [17, Page 957, (4.12)]. The poles of symbol $\mathcal{A}(\tau, \eta)$ on Ξ_1 are exactly the points: $(\tau, \eta) \in \Xi_1$ with $\tau = \mp i\bar{v}\eta$, where the coefficient of \widehat{W}_1^+ or \widehat{W}_1^- in (3.25) vanishes.

By virtue of (3.22) and Lemma 3.1, we infer that, in order to prove Proposition 3.1, it suffices to study problem (3.26). More precisely, we have the following lemma.

Lemma 3.2 *If the solution of (3.26) satisfies the estimate:*

$$\|W^{\text{nc}}|_{x_2=0}\|^2 \lesssim \gamma^{-2} \|G\|_{1,\gamma}^2, \quad (3.29)$$

then Proposition 3.1 holds.

3.2 Lopatinskiĭ Condition

In this subsection, we show that the Kreiss–Lopatinskiĭ condition (or briefly the Lopatinskiĭ condition) holds only in the weak form under assumption (2.25) by computing the Lopatinskiĭ determinant associated to problem (3.26).

We first calculate the stable subspace of the coefficient matrix $\mathcal{A}(\tau, \eta)$, that is, the sum of eigenspaces of $\mathcal{A}(\tau, \eta)$ corresponding to the eigenvalues of negative real parts.

Lemma 3.3 *The following properties hold:*

- (a) *If $(\tau, \eta) \in \Xi_1$ with $\text{Re } \tau > 0$, then the eigenvalues of $\mathcal{A}(\tau, \eta)$ are roots ω of*

$$\begin{aligned} \omega^2 = \mu_+^2 - m_+^2 &= \frac{\bar{\Gamma}^2}{\bar{c}^2} (\tau + i\eta\bar{v})^2 - \bar{\Gamma}^2 (i\eta + \epsilon^2 \bar{v} \tau)^2 \\ &= \bar{C}_0^2 (\bar{C}_1^2 (\tau + i\bar{C}_2 \eta)^2 + \eta^2), \end{aligned} \quad (3.30a)$$

$$\begin{aligned} \omega^2 = \mu_-^2 - m_-^2 &= \frac{\bar{\Gamma}^2}{\bar{c}^2} (\tau - i\eta\bar{v})^2 - \bar{\Gamma}^2 (i\eta - \epsilon^2 \bar{v} \tau)^2 \\ &= \bar{C}_0^2 (\bar{C}_1^2 (\tau - i\bar{C}_2 \eta)^2 + \eta^2), \end{aligned} \quad (3.30b)$$

where \bar{C}_j , $j = 0, 1, 2$, are positive constants defined by

$$\bar{C}_0 := \frac{\bar{\Gamma}(1 - \epsilon^2 \bar{v}^2)}{\sqrt{1 - \epsilon^4 \bar{c}^2 \bar{v}^2}}, \quad \bar{C}_1 := \frac{1 - \epsilon^4 \bar{c}^2 \bar{v}^2}{(1 - \epsilon^2 \bar{v}^2) \bar{c}}, \quad \bar{C}_2 := \frac{(1 - \epsilon^2 \bar{c}^2) \bar{v}}{1 - \epsilon^4 \bar{c}^2 \bar{v}^2}. \quad (3.31)$$

Moreover, (3.30a) (resp. (3.30b)) has a unique root ω_+ (resp. ω_-) of negative real part. The other root of (3.30a) (resp. (3.30b)) is $-\omega_+$ (resp. $-\omega_-$).

- (b) If $(\tau, \eta) \in \Xi_1$ with $\operatorname{Re} \tau > 0$, then the stable subspace $\mathcal{E}^-(\tau, \eta)$ of $\mathcal{A}(\tau, \eta)$ has dimension two and is spanned by

$$E_+(\tau, \eta) := ((\tau + i\bar{v}\eta)m_+, (\tau + i\bar{v}\eta)(\mu_+ - \omega_+), 0, 0)^\top, \quad (3.32a)$$

$$E_-(\tau, \eta) := (0, 0, (\tau - i\bar{v}\eta)(\mu_- - \omega_-), (\tau - i\bar{v}\eta)m_-)^\top. \quad (3.32b)$$

- (c) Both ω_+ and ω_- admit a continuous extension to any point $(\tau, \eta) \in \Xi_1$ with $\operatorname{Re} \tau = 0$. In particular, if $\tau = i\delta \in i\mathbb{R}$, then

$$\begin{aligned} & \omega_\pm(\tau, \eta) \\ &= \begin{cases} -\bar{C}_0 \sqrt{\eta^2 - \bar{C}_1^2(\delta \pm \bar{C}_2\eta)^2} & \text{if } \eta^2 \geq \bar{C}_1^2(\delta \pm \bar{C}_2\eta)^2, \\ -i \operatorname{sgn}(\delta \pm \bar{C}_2\eta) \bar{C}_0 \sqrt{\bar{C}_1^2(\delta \pm \bar{C}_2\eta)^2 - \eta^2} & \text{if } \eta^2 < \bar{C}_1^2(\delta \pm \bar{C}_2\eta)^2. \end{cases} \end{aligned} \quad (3.33)$$

- (d) Vectors $E_\pm(\tau, \eta)$ do not vanish at any point in Ξ_1 . Both $E_+(\tau, \eta)$ and $E_-(\tau, \eta)$ can be extended continuously to any point $(\tau, \eta) \in \Xi_1$ with $\operatorname{Re} \tau = 0$. These two vectors are linearly independent of the whole hemisphere Ξ_1 .
- (e) Matrix $\mathcal{A}(\tau, \eta)$ is diagonalizable as long as eigenvalues ω_\pm do not vanish, i.e. when $\tau \neq i(\mp \bar{C}_2 \pm \bar{C}_1^{-1})\eta$. Apart from these points, $\mathcal{A}(\tau, \eta)$ has a C^∞ basis of eigenvectors.

Proof The relations in (3.30) and assertions (b)–(c) and (e) can be deduced from straightforward calculations and the implicit functions theorem.

It follows from (2.10) and (3.2) that \bar{C}_j is positive. We now show that root ω of (3.30a) is not purely imaginary when $\operatorname{Re} \tau > 0$. If this were not true, there would exist $\sigma \in \mathbb{R}$ such that $i\sigma$ would be a root of (3.30a). Then we would have

$$\bar{C}_0 \bar{C}_1 (\tau + i\bar{C}_2\eta) = \pm i \sqrt{\sigma^2 + \bar{C}_0^2 \eta^2} \in i\mathbb{R},$$

which would imply $\operatorname{Re} \tau = 0$. This concludes assertion (a).

It remains to prove assertion (d). We see from (3.28) that, if $\tau + i\bar{v}\eta = 0$, then $(\tau + i\bar{v}\eta)m_+ = -\bar{c}\eta^2(1 - \epsilon^2\bar{v}^2)/(2\bar{\Gamma}) \neq 0$. Hence, when $(\tau + i\bar{v}\eta)m_+ = 0$,

$$\tau \neq -i\bar{v}\eta, \quad m_+ = 0, \quad \mu_+ = \bar{\Gamma}(\tau + i\bar{v}\eta)/\bar{c} \neq 0, \quad \mu_+^2 = \omega_+^2.$$

Using the relations: $\operatorname{Re} \mu_+ = \bar{\Gamma} \operatorname{Re} \tau / \bar{c} \geq 0$ and $\operatorname{Re} \omega_+ \leq 0$, we have

$$(\tau + i\bar{v}\eta)(\mu_+ - \omega_+) = 2(\tau + i\bar{v}\eta)\mu_+ \neq 0.$$

Therefore, $E_+(\tau, \eta)$ defined by (3.32a) does not vanish. We can also show in a similar way that $E_-(\tau, \eta)$ does not vanish. Assertion (d) then follows. \square

As in Majda–Osher [32], we define the Lopatinskiĭ determinant associated with problem (3.26) by

$$\Delta(\tau, \eta) := \det [\beta(\tau, \eta) (E_+(\tau, \eta) E_-(\tau, \eta))], \quad (3.34)$$

where β and E_{\pm} are given in (3.21) and (3.32), respectively. We say that the Lopatinskiĭ condition holds if $\Delta(\tau, \eta) \neq 0$ for all $(\tau, \eta) \in \Xi$ with $\operatorname{Re} \tau > 0$. Furthermore, if $\Delta(\tau, \eta) \neq 0$ for all $(\tau, \eta) \in \Xi$, we say that the *uniform* Lopatinskiĭ condition holds. To deduce the energy estimate, we need to study the zeros of $\Delta(\tau, \eta)$. For this, we have the following lemma.

Lemma 3.4 *Assume that (2.10) and (2.25) hold. Then, for any $(\tau, \eta) \in \Xi_1$,*

$$\Delta(\tau, \eta) = 0 \quad \text{if and only if} \quad \tau \in \{0, \pm iz_1 \eta\}, \quad (3.35)$$

where z_1 is some positive constant satisfying

$$0 < z_1 < \bar{C}_2 - \bar{C}_1^{-1} < \bar{C}_2 < \bar{v} < \bar{C}_2 + \bar{C}_1^{-1}. \quad (3.36)$$

Moreover, each of these roots is simple in the sense that, if $q \in \{0, -z_1, z_1\}$, then there exists a neighborhood \mathcal{V} of $(iq\eta, \eta)$ in Ξ_1 and a C^∞ -function h_q defined on \mathcal{V} such that

$$\Delta(\tau, \eta) = (\tau - iq\eta)h_q(\tau, \eta), \quad h_q(\tau, \eta) \neq 0 \quad \text{for all } (\tau, \eta) \in \mathcal{V}. \quad (3.37)$$

Proof We divide the proof into seven steps.

1. According to (3.21) and (3.32), we have

$$\begin{aligned} & \beta(\tau, \eta) (E_+(\tau, \eta) E_-(\tau, \eta)) \\ &= \begin{pmatrix} (\tau + i\bar{v}\eta)(m_+ + \mu_+ - \omega_+) & -(\tau - i\bar{v}\eta)(m_- + \mu_- - \omega_-) \\ \frac{\tau - i\bar{v}\eta}{\Gamma\bar{c}h}(\tau + i\bar{v}\eta)(m_+ - \mu_+ + \omega_+) & \frac{\tau + i\bar{v}\eta}{\Gamma\bar{c}h}(\tau - i\bar{v}\eta)(m_- - \mu_- + \omega_-) \end{pmatrix}. \end{aligned} \quad (3.38)$$

By using (3.28) and Lemma 3.3 (a), we have

$$m_{\pm} + \mu_{\pm} = \bar{\Gamma} \frac{\tau \pm i\bar{v}\eta}{\bar{c}}, \quad m_{\pm} - \mu_{\pm} = -\frac{\bar{c}\omega_{\pm}^2}{\bar{\Gamma}(\tau \pm i\bar{v}\eta)}. \quad (3.39)$$

It then follows that

$$\begin{aligned} \Delta(\tau, \eta) &= \frac{1}{\bar{c}^2 h} \left\{ \tau + i\bar{v}\eta - \frac{\bar{c}\omega_+}{\bar{\Gamma}} \right\} \left\{ \tau - i\bar{v}\eta - \frac{\bar{c}\omega_-}{\bar{\Gamma}} \right\} \\ &\quad \times \left\{ \omega_- (\tau + i\bar{v}\eta)^2 + \omega_+ (\tau - i\bar{v}\eta)^2 \right\}. \end{aligned} \quad (3.40)$$

We will check the zeros of each factors in this expression separately.

2. We show in this step that both $\bar{\Gamma}(\tau + i\bar{v}\eta) - \bar{c}\omega_+$ and $\bar{\Gamma}(\tau - i\bar{v}\eta) - \bar{c}\omega_-$ do not vanish at any point $(\tau, \eta) \in \Xi_1$. By contradiction, we assume without loss of generality that there exists a point $(\tau_0, \eta_0) \in \Xi_1$ such that

$$\bar{\Gamma}(\tau_0 + i\bar{v}\eta_0) = \bar{c}\omega_+(\tau_0, \eta_0). \quad (3.41)$$

From (3.30a), we have

$$\bar{\Gamma}^2(\tau_0 + i\bar{v}\eta_0)^2 - \bar{c}^2\omega_+(\tau_0, \eta_0)^2 = \bar{\Gamma}^2\bar{c}^2(i\eta_0 + \epsilon^2\bar{v}\tau_0)^2 = 0,$$

which implies $\tau_0 = i\delta_0 \in i\mathbb{R}$ with $\eta_0 = -\epsilon^2\bar{v}\delta_0$. Since $(\tau_0, \eta_0) \in \Xi_1$, we see that both η_0 and δ_0 are nonzero real numbers. If $\eta_0^2 \geq \bar{C}_1^2(\delta_0 + \bar{C}_2\eta_0)^2$, then $\omega_+(\tau_0, \eta_0) \in \mathbb{R}$ due to (3.33). Then $\bar{c}\omega_+(\tau_0, \eta_0) \neq \bar{\Gamma}(\tau_0 + i\bar{v}\eta_0)$, since

$$\bar{\Gamma}(\tau_0 + i\bar{v}\eta_0) = i\bar{\Gamma}(1 - \epsilon^2\bar{v}^2)\delta_0 \in i\mathbb{R} \setminus \{0\}.$$

According to (3.41), $\eta_0^2 < \bar{C}_1^2(\delta_0 + \bar{C}_2\eta_0)^2$ so that $\bar{C}_1^2(1 - \epsilon^2\bar{C}_2\bar{v})^2 > \epsilon^4\bar{v}^2$. It then follows from (3.33) that

$$\begin{aligned} \omega_+(\tau_0, \eta_0) &= -i \operatorname{sgn}(\delta_0) \operatorname{sgn}(1 - \epsilon^2\bar{C}_2\bar{v})\bar{C}_0\sqrt{(\bar{C}_1^2(1 - \epsilon^2\bar{C}_2\bar{v})^2 - \epsilon^4\bar{v}^2)\delta_0^2} \\ &= -i\delta_0\bar{C}_0\sqrt{\bar{C}_1^2(1 - \epsilon^2\bar{C}_2\bar{v})^2 - \epsilon^4\bar{v}^2}, \end{aligned}$$

where we have used that $1 - \epsilon^2\bar{C}_2\bar{v} = (1 - \epsilon^2\bar{v}^2)/(1 - \epsilon^4\bar{c}^2\bar{v}^2) > 0$ from (3.2). Consequently, we have

$$\begin{aligned} &\bar{\Gamma}(\tau_0 + i\bar{v}\eta_0) - \bar{c}\omega_+(\tau_0, \eta_0) \\ &= i\delta_0 \left\{ \bar{\Gamma}(1 - \epsilon^2\bar{v}^2) + \bar{c}\bar{C}_0\sqrt{\bar{C}_1^2(1 - \epsilon^2\bar{C}_2\bar{v})^2 - \epsilon^4\bar{v}^2} \right\} \neq 0. \end{aligned}$$

This contradicts (3.41).

3. From the above analysis, we know that $\Delta(\tau, \eta) = 0$ if and only if factor $\omega_-(\tau + i\bar{v}\eta)^2 + \omega_+(\tau - i\bar{v}\eta)^2$ vanishes. We first prove that this factor does not vanish when $\eta = 0$.

If $\eta = 0$, then we see from (3.30) that $\omega_{\pm}^2 = \bar{c}^{-2}\bar{\Gamma}^2\tau^2(1 - \epsilon^4\bar{c}^2\bar{v}^2)$. Using (3.2) and noting $\operatorname{Re} \tau \geq 0$, we find that $\omega_{\pm} = -\bar{c}^{-1}\bar{\Gamma}\tau\sqrt{1 - \epsilon^4\bar{c}^2\bar{v}^2}$, which yields

$$\omega_-(\tau + i\bar{v}\eta)^2 + \omega_+(\tau - i\bar{v}\eta)^2 = -2\bar{c}^{-1}\bar{\Gamma}\tau^3(1 - \epsilon^4\bar{c}^2\bar{v}^2)^{1/2} \neq 0.$$

We thus assume that $\eta \neq 0$. Introducing $z := \tau/(i\eta)$, we find from (3.30) that

$$\frac{\bar{c}^2\omega_-^2(\tau + i\bar{v}\eta)^4}{\bar{\Gamma}^2(i\eta)^6} = (z + \bar{v})^4 \left\{ (z - \bar{v})^2 - \bar{c}^2(1 - \epsilon^2\bar{v}z)^2 \right\} =: P_1(z), \quad (3.42)$$

$$\frac{\bar{c}^2\omega_+^2(\tau - i\bar{v}\eta)^4}{\bar{\Gamma}^2(i\eta)^6} = (z - \bar{v})^4 \left\{ (z + \bar{v})^2 - \bar{c}^2(1 + \epsilon^2\bar{v}z)^2 \right\} =: P_2(z). \quad (3.43)$$

Define

$$P(z) := P_1(z) - P_2(z). \quad (3.44)$$

Then $\Delta(\tau, \eta) = 0$ holds only if $\omega_-^2(\tau + i\bar{v}\eta)^4 = \omega_+^2(\tau - i\bar{v}\eta)^4$, which is equivalent to $P(z) = 0$. A straightforward calculation yields

$$P(z) = -4z\bar{v}P_0(z), \quad P_0(z) := E_1z^4 + E_2z^2 + E_3,$$

where $E_1 = 2\epsilon^4\bar{c}^2\bar{v}^2 - \epsilon^2\bar{c}^2 - 1$, $E_2 = 2\epsilon^4\bar{c}^2\bar{v}^4 - 6\epsilon^2\bar{c}^2\bar{v}^2 + 2\bar{v}^2 + 2\bar{c}^2$, and

$$E_3 = 2\bar{c}^2\bar{v}^2 - \epsilon^2\bar{c}^2\bar{v}^4 - \bar{v}^4. \quad (3.45)$$

It is trivial that $z = 0$ is one zero of $P(z)$. Function $P_0(z)$ is a polynomial one of z^2 with the following zeros:

$$-\frac{E_2 \pm \sqrt{E_2^2 - 4E_1E_3}}{2E_1}. \quad (3.46)$$

By virtue of (2.10) and (3.2), we have

$$\begin{cases} E_1 = -(1 - \epsilon^4\bar{c}^2\bar{v}^2) - \epsilon^2\bar{c}^2(1 - \epsilon^2\bar{v}^2) < 0, \\ E_2 = 2\bar{v}^2(1 - \epsilon^2\bar{c}^2) + 2\bar{c}^2(1 - \epsilon^2\bar{v}^2)^2 > 0, \\ E_2^2 - 4E_1E_3 \\ = 4\bar{c}^2(\epsilon\bar{v} - 1)^2(\epsilon\bar{v} + 1)^2(\epsilon^4\bar{c}^4\bar{v}^4 - 2\epsilon^2\bar{c}^2\bar{v}^2 + 4\bar{v}^2 + \bar{c}^2) > 0, \end{cases} \quad (3.47)$$

which yields that the zeros in (3.46) are real and distinct. If (2.25) holds, then $E_3 < 0$, which immediately implies that the zeros in (3.46) are also positive. Let us denote these zeros by z_1^2 and z_2^2 with $0 < z_1 < z_2$ so that

$$z_1^2 = \frac{E_2 - \sqrt{E_2^2 - 4E_1E_3}}{-2E_1}, \quad z_2^2 = \frac{E_2 + \sqrt{E_2^2 - 4E_1E_3}}{-2E_1}. \quad (3.48)$$

Consequently, the Lopatinskiĭ determinant vanishes only if $z \in \{0, \pm z_1, \pm z_2\}$.

4. In this step, we show that the Lopatinskiĭ determinant vanishes when $z = 0$, *i.e.* when $\tau = 0$. We note that $\bar{c} < \bar{v}$ by combining (2.25) and (3.2). Then

$$\bar{C}_2 - \bar{C}_1^{-1} = \frac{(\bar{v} - \bar{c})(1 + \epsilon^2\bar{c}\bar{v})}{1 - \epsilon^4\bar{c}^2\bar{v}^2} > 0. \quad (3.49)$$

It then follows directly from (3.33) that

$$\omega_{\pm}(0, \eta) = -i \operatorname{sgn}(\pm\bar{C}_2\eta)\bar{C}_0\sqrt{\bar{C}_1^2\bar{C}_2^2\eta^2 - \eta^2} = \mp i\eta\bar{C}_0\sqrt{\bar{C}_1^2\bar{C}_2^2 - 1}.$$

Then we infer

$$\{\omega_{-}(\tau + i\bar{v}\eta)^2 + \omega_{+}(\tau - i\bar{v}\eta)^2\}|_{\tau=0} = -\bar{v}^2\eta^2(\omega_{+}(0, \eta) + \omega_{-}(0, \eta)) = 0,$$

and hence $\Delta(0, \eta) = 0$.

5. We prove that $\omega_{-}(\tau + i\bar{v}\eta)^2 + \omega_{+}(\tau - i\bar{v}\eta)^2 \neq 0$ when $z = \pm z_2$, *i.e.* when $\tau = \pm iz_2\eta$. To this end, we need to show that

$$z_2 + \bar{C}_2 > z_2 - \bar{C}_2 > \bar{C}_1^{-1}. \quad (3.50)$$

The first inequality is trivial, so it suffices to prove the second one. From (3.31) and (3.48), we have

$$\begin{aligned} z_2^2 - (\bar{C}_1^{-1} + \bar{C}_2)^2 &= -\frac{E_2 + \sqrt{E_2^2 - 4E_1E_3}}{2E_1} - \frac{((1 - \epsilon^2\bar{v}^2)\bar{c} + (1 - \epsilon^2\bar{c}^2)\bar{v})^2}{(1 - \epsilon^4\bar{c}^2\bar{v}^2)^2} \\ &= -\frac{(1 - \epsilon^4\bar{c}^2\bar{v}^2)^2(\sqrt{E_2^2 - 4E_1E_3} + E_2) + 2E_1((1 - \epsilon^2\bar{v}^2)\bar{c} + (1 - \epsilon^2\bar{c}^2)\bar{v})^2}{2E_1(1 - \epsilon^4\bar{c}^2\bar{v}^2)^2} \\ &= -\frac{2c(1 - \epsilon^2\bar{v}^2)(1 - \epsilon^2\bar{c}\bar{v})^2(R_2 - L_2)}{2E_1(1 - \epsilon^4\bar{c}^2\bar{v}^2)^2}, \end{aligned}$$

where $R_2 := (1 + \epsilon^2\bar{c}\bar{v})^2\sqrt{\epsilon^4\bar{c}^2\bar{v}^4 - 2\epsilon^2\bar{c}^2\bar{v}^2 + 4\bar{v}^2 + \bar{c}^2}$ and

$$L_2 := \epsilon^6\bar{c}^3\bar{v}^4 + 2\epsilon^4\bar{c}^2\bar{v}^3 - 2\epsilon^4\bar{c}^3\bar{v}^2 + 4\epsilon^2\bar{c}\bar{v}^2 + 2\bar{v} + \epsilon^2\bar{c}^3.$$

Then we obtain from (3.2) that

$$R_2^2 - L_2^2 = \bar{c}^2(\epsilon\bar{c} - 1)(\epsilon\bar{c} + 1)(\epsilon\bar{v} - 1)^2(\epsilon\bar{v} + 1)^2(2\epsilon^4\bar{c}^2\bar{v}^2 - \epsilon^2\bar{c}^2 - 1) > 0,$$

which, combined with (3.2) and (3.47), implies that $z_2^2 > (\bar{C}_1^{-1} + \bar{C}_2)^2$. Then (3.50) follows.

In view of (3.50), we see from (3.33) that, for $\tau = iz_2\eta$,

$$\begin{aligned} \omega_{\pm}(\tau, \eta) &= -i \operatorname{sgn}(z_2\eta \pm \bar{C}_2\eta)\bar{C}_0\sqrt{\bar{C}_1^2(z_2 \pm \bar{C}_2)^2\eta^2 - \eta^2} \\ &= -i\eta\bar{C}_0\sqrt{\bar{C}_1^2(z_2 \pm \bar{C}_2)^2 - 1}. \end{aligned}$$

Therefore, we obtain that

$$\begin{aligned} \omega_{-}(\tau + i\bar{v}\eta)^2 + \omega_{+}(\tau - i\bar{v}\eta)^2 &= -\eta^2(\omega_{+}(z_2 - \bar{v})^2 + \omega_{-}(z_2 + \bar{v})^2) \\ &= i\eta^3\left\{\bar{C}_0\sqrt{\bar{C}_1^2(z_2 + \bar{C}_2)^2 - 1}(z_2 - \bar{v})^2 + \bar{C}_0\sqrt{\bar{C}_1^2(z_2 - \bar{C}_2)^2 - 1}(z_2 + \bar{v})^2\right\}, \end{aligned}$$

which is away from zero. Applying a similar argument and using (3.50) imply that the Lopatinskiĭ determinant Δ does not vanish either for the case: $z = -z_2$.

6. Let us now show that $\omega_{-}(\tau + i\bar{v}\eta)^2 + \omega_{+}(\tau - i\bar{v}\eta)^2 = 0$ if $z = \pm z_1$, *i.e.* if $\tau = \pm iz_1\eta$. For this purpose, we first prove

$$z_1 + \bar{C}_2 > \bar{C}_1^{-1}, \quad z_1 - \bar{C}_2 < -\bar{C}_1^{-1}. \quad (3.51)$$

The first inequality in (3.51) follows from (3.49). For the second in (3.51), it suffices to derive that $z_1^2 < (\bar{C}_2 - \bar{C}_1^{-1})^2$. From (3.31) and (3.48), we have

$$\begin{aligned} (\bar{C}_2 - \bar{C}_1^{-1})^2 - z_1^2 &= \frac{((1 - \epsilon^2\bar{c}^2)\bar{v} - (1 - \epsilon^2\bar{v}^2)\bar{c})^2}{(1 - \epsilon^4\bar{c}^2\bar{v}^2)^2} + \frac{E_2 - \sqrt{E_2^2 - 4E_1E_3}}{2E_1} \\ &= \frac{(1 - \epsilon^4\bar{c}^2\bar{v}^2)^2(E_2 - \sqrt{E_2^2 - 4E_1E_3}) + 2E_1((1 - \epsilon^2\bar{c}^2)\bar{v} - (1 - \epsilon^2\bar{v}^2)\bar{c})^2}{2E_1(1 - \epsilon^4\bar{c}^2\bar{v}^2)^2} \\ &= \frac{2c(1 - \epsilon^2\bar{v}^2)(1 + \epsilon^2\bar{c}\bar{v})^2(R_4 + L_4)}{-2E_1(1 - \epsilon^4\bar{c}^2\bar{v}^2)^2}, \end{aligned}$$

where $R_4 := (1 - \epsilon^2 \bar{c}\bar{v})^2 \sqrt{\epsilon^4 \bar{c}^2 \bar{v}^4 - 2\epsilon^2 \bar{c}^2 \bar{v}^2 + 4\bar{v}^2 + \bar{c}^2}$ and

$$L_4 := \epsilon^6 \bar{c}^3 \bar{v}^4 - 2\epsilon^4 \bar{c}^2 \bar{v}^3 - 2\epsilon^4 \bar{c}^3 \bar{v}^2 + 4\epsilon^2 \bar{c}\bar{v}^2 - 2\bar{v} + \epsilon^2 \bar{c}^3.$$

We compute that $R_4^2 - L_4^2 = R_2^2 - L_2^2 > 0$. Hence, we deduce the second inequality in (3.51).

By virtue of (3.33) and (3.51), we derive that, for $\tau = iz_1\eta$,

$$\omega_+(\tau, \eta) = -i\eta \bar{C}_0 \sqrt{\bar{C}_1^2 (z_1 + \bar{C}_2)^2 - 1}, \quad \omega_-(\tau, \eta) = i\eta \bar{C}_0 \sqrt{\bar{C}_1^2 (z_1 - \bar{C}_2)^2 - 1}.$$

Since $z = z_1$ solves $P(z) = 0$, if $\tau = iz_1\eta$, it follows from the definition of $P(z)$ that $\omega_-^2(\tau + i\bar{v}\eta)^4 = \omega_+^2(\tau - i\bar{v}\eta)^4$. Hence, the Lopatinskiĭ determinant vanishes for $z = z_1$ (*i.e.* $\tau = iz_1\eta$). The same argument can be applied to show that the Lopatinskiĭ determinant $\Delta(\tau, \eta)$ also vanishes for $z = -z_1$, *i.e.* for $\tau = -iz_1\eta$.

7. We obtain from (3.2) by a direct computation that $\bar{C}_2 < \bar{v} < \bar{C}_2 + \bar{C}_1^{-1}$, which, combined with (3.51), yields (3.36).

It remains to show that the roots of the Lopatinskiĭ determinant are simple. By introducing $\Omega_\pm := \omega_\pm/(i\eta)$, we find that, for $\eta \neq 0$,

$$\frac{\omega_-(\tau + i\bar{v}\eta)^2}{(i\eta)^3} = \Omega_-(z + \bar{v})^2 =: Q_1(z), \quad \frac{\omega_+(\tau - i\bar{v}\eta)^2}{(i\eta)^3} = \Omega_+(z - \bar{v})^2 =: Q_2(z).$$

It follows from (3.36) and Lemma 3.3 that $\omega_\pm(\tau, \eta) \neq 0$ and $\eta \neq 0$ when (τ, η) are near any root of the Lopatinskiĭ determinant. Hence, Ω_\pm are analytic functions of z only and satisfy

$$\Omega_\pm^2 = \bar{c}^{-2} \bar{I}^2 ((z \pm \bar{v})^2 - \bar{c}^2 (1 \pm \epsilon^2 \bar{v}z)^2).$$

Since $\Delta(\tau, \eta) = 0$ if and only if $\omega_-(\tau + i\bar{v}\eta)^2 + \omega_+(\tau - i\bar{v}\eta)^2 = 0$, it suffices to prove

$$\left. \frac{d(Q_1 + Q_2)}{dz} \right|_{z=q} \neq 0 \quad \text{for all } q \in \{0, -z_1, z_1\}.$$

Using (3.42)–(3.44) and the fact that $Q_1(q) = -Q_2(q) \neq 0$ for $q \in \{0, -z_1, z_1\}$, we derive that, for $q \in \{0, -z_1, z_1\}$,

$$\begin{aligned} \left. \frac{d(Q_1 + Q_2)}{dz} \right|_{z=q} &= \frac{1}{2Q_1(q)} \left. \frac{d(Q_1^2 - Q_2^2)}{dz} \right|_{z=q} = \frac{\bar{I}^2}{2\bar{c}^2 Q_1(q)} \left. \frac{dP}{dz} \right|_{z=q} \\ &= \frac{\bar{I}^2}{\bar{c}^2 Q_1(q)} \{-2\bar{v}P_0(q) - 4\bar{v}q^2(2E_1q^2 + E_2)\} \neq 0. \end{aligned}$$

Using the factorization property of holomorphic functions, we obtain

$$Q_1(z) + Q_2(z) = (z - q)H_q(z) \quad \text{for all } q \in \{0, -z_1, z_1\},$$

where H_q is holomorphic near q and $H_q(q) \neq 0$. This yields that the Lopatinskiĭ determinant Δ has a factorization:

$$\Delta(\tau, \eta) = (\tau - iq\eta)h_q(\tau, \eta) \quad \text{for all } q \in \{0, -z_1, z_1\},$$

where $h_q(\tau, \eta)$ is C^∞ and does not vanish near $(iq\eta, \eta) \in \Xi_1$. The proof is completed. \square

Remark 3.2 If $M = M_c$, then both E_3 and z_1 defined by (3.45) and (3.48) vanish. Employing a similar argument, we can show that the Lopatinskiĭ determinant $\Delta(\tau, \eta)$ has only one triple root $\tau = 0$. On the other hand, if $M < M_c$, then $E_3 > 0$. In the latter case, the Lopatinskiĭ determinant $\Delta(\tau, \eta)$ vanishes if and only if $\tau/(i\eta) \in \{0, \pm z_1\}$ with nonreal number z_1 given by (3.48). Therefore, the relativistic vortex sheet (2.24) is violently unstable, which means that the Lopatinskiĭ condition does not hold.

3.3 Proof of Theorem 3.1

The following lemma relies heavily on the fact that each root of the Lopatinskiĭ determinant is simple (see Lemma 3.4).

Lemma 3.5 *For every point $(\tau_0, \eta_0) \in \Xi_1$, there exists a neighborhood \mathcal{V} of (τ_0, η_0) in Ξ_1 and a positive constant c depending on (τ_0, η_0) such that*

$$|\beta(\tau, \eta)(E_+(\tau, \eta) E_-(\tau, \eta))Z| \geq c\gamma|Z| \quad \text{for all } (\tau, \eta) \in \mathcal{V}, Z \in \mathbb{C}^2. \quad (3.52)$$

Proof The proof is divided into two steps.

1. Let $(\tau_0, \eta_0) \in \Xi_1$ with $\Delta(\tau_0, \eta_0) \neq 0$. Since the Lopatinskiĭ determinant $\Delta(\tau, \eta)$ is continuous in (τ, η) , then there exists a neighborhood \mathcal{V} of (τ_0, η_0) in Ξ_1 such that $\Delta(\tau, \eta) \neq 0$ for all $(\tau, \eta) \in \mathcal{V}$. It follows from definition (3.34) that $\beta(\tau, \eta)(E_+ E_-)$ is invertible in \mathcal{V} . We combine this with the fact that $\gamma \leq 1$ to obtain (3.52).

2. Let $(\tau_0, \eta_0) \in \Xi_1$ such that $\Delta(\tau_0, \eta_0) = 0$. We see from Lemma 3.4 that $\tau_0 = iq\eta_0$ for some $q \in \{0, -z_1, z_1\}$. Let us write (3.38) as

$$\beta(E_+ E_-) = \begin{pmatrix} \zeta_1 & \zeta_2 \\ \zeta_3 & \zeta_4 \end{pmatrix},$$

where the upper left corner ζ_1 is given by

$$\zeta_1 = (\tau - i\bar{v}\eta)(m_+ + \mu_+ - \omega_+) = \frac{\tau + i\bar{v}\eta}{\bar{c}} (\bar{I}(\tau + i\bar{v}\eta) - \bar{c}\omega_+).$$

From (3.36) and the proof of Lemma 3.4 (especially, Step 2), we know that $\tau \neq -i\bar{v}\eta$ and $\bar{I}(\tau + i\bar{v}\eta) \neq \bar{c}\omega_+$ when (τ, η) is close to (τ_0, η_0) . Hence, there exists a neighborhood \mathcal{V} of (τ_0, η_0) in Ξ_1 such that $\zeta_1(\tau, \eta) \neq 0$ for all $(\tau, \eta) \in \mathcal{V}$. Using the identity (*cf.* [15, Page 439]):

$$\begin{pmatrix} 1/\zeta_1 & 0 \\ -\zeta_3/(\zeta_1\zeta_5) & 1/\zeta_5 \end{pmatrix} \beta(E_+ E_-) \begin{pmatrix} 1 & -\zeta_2 \\ 0 & \zeta_1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & (\zeta_1\zeta_4 - \zeta_2\zeta_3)/\zeta_5 \end{pmatrix} \quad (3.53)$$

with $\zeta_5 = 1$, and noting $\Delta(\tau, \eta) = \det[\beta(E_+ E_-)] = \zeta_1\zeta_4 - \zeta_2\zeta_3$, we have

$$|\beta(\tau, \eta)(E_+(\tau, \eta) E_-(\tau, \eta))Z| \geq c \min(1, |\Delta(\tau, \eta)|)|Z| \quad (3.54)$$

for all $(\tau, \eta) \in \mathcal{V}$, $Z \in \mathbb{C}^2$. It thus remains to show that $|\Delta(\tau, \eta)| \geq c\gamma$ for all $(\tau, \eta) \in \mathcal{V}$. Employ Lemma 3.4 and shrink \mathcal{V} if necessary to find that factorization (3.37) holds. Thus, we have

$$\partial_\gamma \Delta(\tau, \eta) = h_q(\tau, \eta) + (\tau - iq\eta)\partial_\gamma h_q(\tau, \eta) \quad \text{for all } (\tau, \eta) \in \mathcal{V}. \quad (3.55)$$

Let $(i\delta, \eta) \in \mathcal{V}$ so that $(i\delta, \eta) \in \mathcal{V}$ is close to $(iq\eta_0, \eta_0)$. It follows from (3.36) that $\bar{C}_1^2(\delta + \bar{C}_2\eta)^2 > \eta^2$, which, combined with (3.33), implies

$$\omega_\pm(i\delta, \eta) \in i\mathbb{R} \setminus \{0\}. \quad (3.56)$$

Then we obtain from (3.37) and expression (3.40) that

$$h_q(i\delta, \eta) \neq 0, \quad i(\delta - q\eta)h_q(i\delta, \eta) = \Delta(i\delta, \eta) \in i\mathbb{R},$$

from which we have

$$h_q(i\delta, \eta) \in \mathbb{R} \setminus \{0\}. \quad (3.57)$$

Since $\tau_0 \neq i(\pm\bar{C}_2 \pm \bar{C}_1^{-1})\eta_0$, eigenvalues ω_\pm depend analytically on (τ, η) in a neighborhood of (τ_0, η_0) by the implicit function theorem. We then use (3.30) to obtain that, for (τ, η) near (τ_0, η_0) ,

$$\omega_+(\tau, \eta)\partial_\gamma \omega_+(\tau, \eta) = \bar{C}_0^2 \bar{C}_1^2 (\tau + i\bar{C}_2\eta). \quad (3.58)$$

From (3.56) and (3.58), we infer that the derivative, $\partial_\gamma \omega_+(i\delta, \eta)$, is real by shrinking \mathcal{V} if necessary. Employ (3.40) to derive $\partial_\gamma \Delta(i\delta, \eta) \in \mathbb{R}$. We then deduce from (3.55) and (3.57) that

$$\partial_\gamma h_q(i\delta, \eta) \in i\mathbb{R}. \quad (3.59)$$

Using (3.37) and the Taylor formula for h_q , we find that, for $(\tau, \eta) \in \mathcal{V}$,

$$\begin{aligned} \Delta(\tau, \eta) &= (\gamma + i(\delta - q\eta)) (h_q(i\delta, \eta) + \gamma\partial_\gamma h_q(i\delta, \eta) + O(\gamma^2)) \\ &= i(\delta - q\eta)h_q(i\delta, \eta) + \{h_q(i\delta, \eta) + i\partial_\gamma h_q(i\delta, \eta)(\delta - q\eta)\}\gamma + O(\gamma^2) \quad (\gamma \rightarrow 0), \end{aligned}$$

where we have used the Landau symbol $f = O(g)$ ($x \rightarrow x_0$), which means that there exists a constant C such that $|f(x)| \leq C|g(x)|$ for all x sufficiently close to x_0 . We can conclude from (3.57) and (3.59) that

$$\operatorname{Re} \Delta(\tau, \eta) = \{h_q(i\delta, \eta) + i\partial_\gamma h_q(i\delta, \eta)(\delta - q\eta)\}\gamma + O(\gamma^2) \quad (\gamma \rightarrow 0).$$

Shrinking \mathcal{V} if necessary, we derive from (3.57) that

$$|\Delta(\tau, \eta)| \geq |\operatorname{Re} \Delta(\tau, \eta)| \geq c\gamma \quad \text{for all } (\tau, \eta) \in \mathcal{V}.$$

Plug this into (3.54) to complete the proof of this lemma. \square

Remark 3.3 In the case of $\mathbf{M} = \mathbf{M}_c$, we know from Remark 3.2 that the Lopatinskiĭ determinant $\Delta(\tau, \eta)$ has only one triple root $\tau = 0$. In a similar way, we can find neighborhoods \mathcal{V}_\pm of $(0, \pm 1)$ in Ξ_1 and a positive constant c such that

$$|\beta(\tau, \eta)(E_+(\tau, \eta) E_-(\tau, \eta))Z| \geq c\gamma^3|Z| \quad \text{for all } (\tau, \eta) \in \mathcal{V}_\pm, Z \in \mathbb{C}^2. \quad (3.60)$$

We now adopt the argument developed recently by Chen–Hu–Wang [10] to avoid constructing the Kreiss’ symmetrizers in the derivation of energy estimates for the constant coefficient case. To this end, we need the following lemma.

Lemma 3.6 *For each point $(\tau_0, \eta_0) \in \Xi_1$, there exist a neighborhood \mathcal{V} of (τ_0, η_0) in Ξ_1 and a continuous invertible matrix $T(\tau, \eta)$ defined on \mathcal{V} such that*

$$T^{-1}\mathcal{A}T(\tau, \eta) = \begin{pmatrix} \omega_+ & z_+ & 0 & 0 \\ 0 & -\omega_+ & 0 & 0 \\ 0 & 0 & \omega_- & z_- \\ 0 & 0 & 0 & -\omega_- \end{pmatrix} \quad (3.61)$$

for all $(\tau, \eta) \in \mathcal{V} \setminus \{\tau = \pm i\bar{v}\eta\}$, where $z_{\pm} = z_{\pm}(\tau, \eta)$ are complex-valued functions defined on $\mathcal{V} \setminus \{\tau = \pm i\bar{v}\eta\}$. Moreover, the first and third columns of $T(\tau, \eta)$ are $E_+(\tau, \eta)$ and $E_-(\tau, \eta)$, respectively.

Proof We set $a_{\pm}(\tau, \eta) := \tau \pm i\bar{v}\eta$ and define the following vectors on a neighborhood \mathcal{V} of (τ_0, η_0) :

$$Y_+(\tau, \eta) := \begin{cases} (0, 1, 0, 0)^{\top} & \text{if } a_+m_+(\tau_0, \eta_0) \neq 0, \\ (1, 0, 0, 0)^{\top} & \text{if } a_+(\mu_+ - \omega_+)(\tau_0, \eta_0) \neq 0, \end{cases}$$

$$Y_-(\tau, \eta) := \begin{cases} (0, 0, 1, 0)^{\top} & \text{if } a_-m_-(\tau_0, \eta_0) \neq 0, \\ (0, 0, 0, 1)^{\top} & \text{if } a_-(\mu_- - \omega_-)(\tau_0, \eta_0) \neq 0. \end{cases}$$

Recall that $E_{\pm}(\tau, \eta)$ defined by (3.32) are continuous and never vanish on Ξ_1 . Hence, one can define the following continuous and invertible matrix on \mathcal{V} :

$$T(\tau, \eta) := (E_+(\tau, \eta) \ Y_+(\tau, \eta) \ E_-(\tau, \eta) \ Y_-(\tau, \eta)).$$

When $\tau \neq \pm i\bar{v}\eta$, by a direct computation and using (3.30), we obtain (3.61) with

$$z_{\pm}(\tau, \eta) := \begin{cases} -\frac{1}{a_{\pm}(\tau, \eta)} & \text{if } a_{\pm}m_{\pm}(\tau_0, \eta_0) \neq 0, \\ -\frac{m_{\pm}}{a_{\pm}(\mu_{\pm} - \omega_{\pm})(\tau, \eta)} & \text{if } a_{\pm}(\mu_{\pm} - \omega_{\pm})(\tau_0, \eta_0) \neq 0, \end{cases}$$

which are well-defined apart from the poles of \mathcal{A} , *i.e.* from $\tau = \pm i\bar{v}\eta$. \square

Proof of Theorem 3.1 According to Lemma 3.2, it suffices to show estimate (3.29) in order to prove Theorem 3.1. Using Lemmas 3.5–3.6, for each point $(\tau_0, \eta_0) \in \Xi_1$, there exists a neighborhood \mathcal{V} of (τ_0, η_0) in Ξ_1 and a continuous invertible matrix $T(\tau, \eta)$ defined on \mathcal{V} such that (3.52) and (3.61)

hold. Thanks to the compactness of hemisphere Ξ_1 , there exists a finite covering $\{\mathcal{V}_1, \dots, \mathcal{V}_J\}$ of Ξ_1 by such neighborhoods with corresponding matrices $\{T_1(\tau, \eta), \dots, T_J(\tau, \eta)\}$, and a smooth partition of unity $\{\chi_j(\tau, \eta)\}_{j=1}^J$ such that

$$\chi_j \in C_c^\infty(\mathcal{V}_j), \quad \sum_{j=1}^J \chi_j^2 = 1 \quad \text{on } \Xi_1.$$

We now derive an energy estimate in $\Pi_j := \{(\tau, \eta) \in \Xi : s \cdot (\tau, \eta) \in \mathcal{V}_j \text{ for some } s > 0\}$ and then patch them together to obtain (3.29). We first extend χ_j and T_j to the conic zone Π_j as homogeneous mappings of degree 0 with respect to (τ, η) . Note that both $T_j(\tau, \eta)$ and its inverse are bounded on Π_j , and identity (3.61) holds for all $(\tau, \eta) \in \Pi_j$ with $\tau \neq \pm i\bar{v}\eta$. Define

$$W(\tau, \eta, x_2) := \chi_j T_j(\tau, \eta)^{-1} \widehat{W}^{\text{nc}}(\tau, \eta, x_2) \quad \text{for all } (\tau, \eta) \in \Pi_j.$$

Assume that $(\tau, \eta) \in \Pi_j$ with $\text{Re } \tau > 0$. In light of (3.26), we obtain that W satisfies

$$\frac{dW}{dx_2} = T_j(\tau, \eta)^{-1} \mathcal{A} T_j(\tau, \eta) W.$$

Since (3.61) holds when $(\tau, \eta) \in \Pi_j$ with $\text{Re } \tau > 0$, the equations for W_2 and W_4 read

$$\frac{dW_2}{dx_2} = -\omega_+ W_2, \quad \frac{dW_4}{dx_2} = -\omega_- W_4. \quad (3.62)$$

Recall from Lemma 3.3 (a) that $\text{Re } \omega_\pm(\tau, \eta) < 0$ whenever $\text{Re } \tau > 0$. Integration by parts for (3.62) yields

$$\|W_2(\tau, \eta, \cdot)\|_{L^2(\mathbb{R}_+)} = \|W_4(\tau, \eta, \cdot)\|_{L^2(\mathbb{R}_+)} = 0,$$

from which we immediately deduce

$$W_2(\tau, \eta, x_2) = W_4(\tau, \eta, x_2) = 0 \quad (3.63)$$

for all $x_2 \in \mathbb{R}_+$ and $(\tau, \eta) \in \Pi_j$ with $\text{Re } \tau > 0$, where we have used the continuity of W_2 and W_4 . Using the boundary equations in (3.26) yields

$$\chi_j \widehat{G} = \beta(\tau, \eta) T_j(\tau, \eta) W(\tau, \eta, 0) = \beta(\tau, \eta) (E_+ \ E_-) \begin{pmatrix} W_1(\tau, \eta, 0) \\ W_3(\tau, \eta, 0) \end{pmatrix} \quad (3.64)$$

for all $(\tau, \eta) \in \Pi_j$ with $\text{Re } \tau > 0$. By the homogeneity of T_j and β , we obtain from (3.52) that

$$(|\tau| + |\eta|) |\beta(\tau, \eta) (E_+(\tau, \eta) \ E_-(\tau, \eta)) Z| \geq c_j \gamma |Z| \quad \text{for all } (\tau, \eta) \in \Pi_j, Z \in \mathbb{C}^2.$$

Combine this with (3.64) to deduce

$$|(W_1(\tau, \eta, 0), W_3(\tau, \eta, 0))| \leq \frac{|\tau| + |\eta|}{c_j \gamma} |\chi_j \widehat{G}(\tau, \eta)| \quad (3.65)$$

for all $(\tau, \eta) \in \Pi_j$ with $\operatorname{Re} \tau > 0$. Combining (3.63) and (3.65) yields

$$|\mathbb{W}(\tau, \eta, 0)| \leq \frac{|\tau| + |\eta|}{c_j \gamma} |\chi_j \widehat{G}(\tau, \eta)| \quad \text{for all } (\tau, \eta) \in \Pi_j \text{ with } \operatorname{Re} \tau > 0.$$

We then obtain from the definition of \mathbb{W} and boundedness of $T_j(\tau, \eta)$ that

$$|\chi_j \widehat{W}^{\text{nc}}(\tau, \eta, 0)| \leq \frac{|\tau| + |\eta|}{c_j \gamma} |\chi_j \widehat{G}(\tau, \eta)|$$

for all $(\tau, \eta) \in \Pi_j$ with $\gamma = \operatorname{Re} \tau > 0$ and new positive constants c_j . Adding the above estimates for all $j \in \{1, \dots, J\}$ and integrating the resulting estimate over \mathbb{R}^2 with respect to (δ, η) , we can derive the desired estimate (3.29) from the Plancherel theorem. This completes the proof of Theorem 3.1. \square

Remark 3.4 In the case of $\mathbb{M} = \mathbb{M}_c$, we can derive the energy estimate (3.7) by using (3.60) and employing a completely similar argument as above.

4 Variable Coefficient Linearized Problem

In this section, we derive the linearized problem of (2.20) around a basic state $(\bar{U}^\pm, \bar{\Phi}^\pm)$ that is a small perturbation of $(\bar{U}^\pm, \bar{\Phi}^\pm)$ given in (2.24). More precisely, we assume that the perturbations: $\mathring{V}^\pm := \mathring{U}^\pm - \bar{U}^\pm$ and $\mathring{\Psi}^\pm := \mathring{\Phi}^\pm - \bar{\Phi}^\pm$ satisfy

$$\operatorname{supp} (\mathring{V}^\pm, \mathring{\Psi}^\pm) \subset \{-T \leq t \leq 2T, x_2 \geq 0, |x| \leq R\}, \quad (4.1)$$

$$\mathring{V}^\pm \in W^{2,\infty}(\Omega), \quad \mathring{\Psi}^\pm \in W^{3,\infty}(\Omega), \quad \|\mathring{V}^\pm\|_{W^{2,\infty}(\Omega)} + \|\mathring{\Psi}^\pm\|_{W^{3,\infty}(\Omega)} \leq K, \quad (4.2)$$

where T , R , and K are positive constants. Moreover, we assume that $(\mathring{U}^\pm, \mathring{\Phi}^\pm)$ satisfies constraints (2.23) and the Rankine–Hugoniot conditions (2.20b):

$$\partial_t \mathring{\Phi}^\pm + v_1^\pm \partial_1 \mathring{\Phi}^\pm - v_2^\pm = 0 \quad \text{if } x_2 \geq 0, \quad (4.3a)$$

$$\pm \partial_2 \mathring{\Phi}^\pm \geq \kappa_0 > 0 \quad \text{if } x_2 \geq 0, \quad (4.3b)$$

$$\mathring{\Phi}^+ = \mathring{\Phi}^- = \mathring{\varphi} \quad \text{if } x_2 = 0, \quad (4.3c)$$

$$\mathbb{B}(\mathring{U}^+, \mathring{U}^-, \mathring{\varphi}) = 0 \quad \text{if } x_2 = 0, \quad (4.3d)$$

where κ_0 is a positive constant. We will use $\mathring{V} := (\mathring{V}^+, \mathring{V}^-)^\top$ and $\mathring{\Psi} := (\mathring{\Psi}^+, \mathring{\Psi}^-)^\top$ to avoid overloaded expressions.

4.1 Linearized Problem

Let us consider the families, $U_\theta^\pm = \mathring{U}^\pm + \theta V^\pm$ and $\Phi_\theta^\pm = \mathring{\Phi}^\pm + \theta \Psi^\pm$, with a small parameter θ . The linearized operators are given by

$$\begin{cases} \mathbb{L}'(\mathring{U}^\pm, \mathring{\Phi}^\pm)(V^\pm, \Psi^\pm) := \frac{d}{d\theta} \mathbb{L}(U_\theta^\pm, \Phi_\theta^\pm) \Big|_{\theta=0}, \\ \mathbb{B}'(\mathring{U}^\pm, \mathring{\Phi}^\pm)(V, \psi) := \frac{d}{d\theta} \mathbb{B}(U_\theta^+, U_\theta^-, \varphi_\theta) \Big|_{\theta=0}, \end{cases}$$

where $V := (V^+, V^-)^\top$, and φ_θ (resp. ψ) denotes the common trace of Φ_θ^\pm (resp. Ψ^\pm) on boundary $\{x_2 = 0\}$. A standard computation yields the following expression for \mathbb{L}' :

$$\mathbb{L}'(U, \Phi)(V, \Psi) = L(U, \Phi)V + \mathcal{C}(U, \Phi)V - \frac{1}{\partial_2 \Phi} L(U, \Phi)\Psi \partial_2 U, \quad (4.4)$$

where $\mathcal{C}(U, \Phi)$ is the zero-th order operator defined by

$$\mathcal{C}(U, \Phi)V := (\partial_{U_i} A_0(U) \partial_t U + \partial_{U_i} A_1(U) \partial_1 U + \partial_{U_i} \tilde{A}_2(U, \Phi) \partial_2 U) V_i. \quad (4.5)$$

We notice that matrices $\mathcal{C}(\mathring{U}^\pm, \mathring{\Phi}^\pm)$ are C^∞ -functions of $(\mathring{V}^\pm, \nabla \mathring{V}^\pm, \nabla \mathring{\Psi}^\pm)$ vanishing at the origin.

We recall that the first component of $\mathbb{B}(U_\theta^+, U_\theta^-, \varphi_\theta)$ is $[v_1(U_\theta)] \partial_1 \varphi_\theta - [v_2(U_\theta)]$. Ignoring indices “+” and “-” for the moment, it follows from (3.4) and (4.3a) that

$$\begin{aligned} \frac{d}{d\theta} (\mathbb{B}(U_\theta^+, U_\theta^-, \varphi_\theta)) \Big|_{\theta=0} &= \mathring{v}_1 \partial_1 \psi + \partial_1 \mathring{\varphi} \nabla_U v_1(\mathring{U}) \cdot V - \nabla_U v_2(\mathring{U}) \cdot V \\ &= \mathring{v}_1 \partial_1 \psi - \frac{\epsilon^2 (\partial_1 \mathring{\varphi} \mathring{v}_1 - \mathring{v}_2)}{\mathring{N} \mathring{h} \mathring{\Gamma}^2} V_1 + \frac{\partial_1 \mathring{\varphi} (1 - \epsilon^2 \mathring{v}_1^2) + \epsilon^2 \mathring{v}_1 \mathring{v}_2}{\mathring{h} \mathring{\Gamma}} V_2 \\ &\quad - \frac{\partial_1 \mathring{\varphi} \epsilon^2 \mathring{v}_1 \mathring{v}_2 + (1 - \epsilon^2 \mathring{v}_2^2)}{\mathring{h} \mathring{\Gamma}} V_3 \\ &= \mathring{v}_1 \partial_1 \psi + \frac{\epsilon^2 \partial_t \mathring{\varphi}}{\mathring{N} \mathring{h} \mathring{\Gamma}^2} V_1 + \frac{\partial_1 \mathring{\varphi} + \epsilon^2 \mathring{v}_1 \partial_t \mathring{\varphi}}{\mathring{h} \mathring{\Gamma}} V_2 - \frac{1 - \epsilon^2 \mathring{v}_2 \partial_t \mathring{\varphi}}{\mathring{h} \mathring{\Gamma}} V_3 \quad \text{on } \{x_2 = 0\}, \end{aligned}$$

where $\mathring{v}_j := v_j(\mathring{U})$, $\mathring{N} := N(\mathring{U}_1)$, $\mathring{h} := h(\mathring{U}_1)$, and $\mathring{\Gamma} := \Gamma(\mathring{U})$. Performing a similar analysis to the other components of $\mathbb{B}(U_\theta^+, U_\theta^-, \varphi_\theta)$ implies

$$\mathbb{B}'(\mathring{U}^\pm, \mathring{\Phi}^\pm)(V, \psi) := \mathring{b} \nabla \psi + \mathring{B} V|_{x_2=0}, \quad (4.6)$$

where $\nabla \psi := (\partial_t \psi, \partial_1 \psi)^\top$. Coefficients \mathring{b} and \mathring{B} are defined by

$$\mathring{b}(t, x_1) := \begin{pmatrix} 0 & (\mathring{v}_1^+ - \mathring{v}_1^-)|_{x_2=0} \\ 1 & \mathring{v}_1^+|_{x_2=0} \\ 0 & 0 \end{pmatrix}, \quad (4.7)$$

$$\mathring{B}(t, x_1) := \begin{pmatrix} \frac{\epsilon^2 \partial_t \mathring{\varphi}}{\mathring{N}_+ \mathring{h}_+ \mathring{\Gamma}_+^2} & \frac{\mathring{\varrho}_+}{\mathring{h}_+ \mathring{\Gamma}_+} & \frac{-\mathring{\zeta}_+}{\mathring{h}_+ \mathring{\Gamma}_+} & \frac{-\epsilon^2 \partial_t \mathring{\varphi}}{\mathring{N}_- \mathring{h}_- \mathring{\Gamma}_-^2} & \frac{-\mathring{\varrho}_-}{\mathring{h}_- \mathring{\Gamma}_-} & \frac{\mathring{\zeta}_-}{\mathring{h}_- \mathring{\Gamma}_-} \\ \frac{\epsilon^2 \partial_t \mathring{\varphi}}{\mathring{N}_+ \mathring{h}_+ \mathring{\Gamma}_+^2} & \frac{\mathring{\varrho}_+}{\mathring{h}_+ \mathring{\Gamma}_+} & \frac{-\mathring{\zeta}_+}{\mathring{h}_+ \mathring{\Gamma}_+} & 0 & 0 & 0 \\ 1 & 0 & 0 & -1 & 0 & 0 \end{pmatrix} \Big|_{x_2=0}. \quad (4.8)$$

In expression (4.8), we have set $\hat{\varrho}_\pm := \varrho(\hat{U}^\pm, \hat{\Phi}^\pm)$ and $\hat{\zeta}_\pm := \zeta(\hat{U}^\pm, \hat{\Phi}^\pm)$, where

$$\varrho(U, \Phi) := \partial_1 \Phi + \epsilon^2 v_1 \partial_t \Phi, \quad \zeta(U, \Phi) := 1 - \epsilon^2 v_2 \partial_t \Phi. \quad (4.9)$$

In particular, if $\hat{\Psi}^\pm \equiv 0$, then $\hat{\varrho}_\pm \equiv 0$ and $\hat{\zeta}_\pm \equiv 1$. Moreover, \hat{b} is a C^∞ -function of $\hat{V}|_{x_2=0}$, and \hat{B} is a C^∞ -function of $(\hat{V}|_{x_2=0}, \nabla \hat{\varphi})$.

We simplify expression (4.4) as Alinhac [1] by employing the ‘‘good unknown’’:

$$\dot{V}^\pm := V^\pm - \frac{\partial_2 \hat{U}^\pm}{\partial_2 \hat{\Phi}^\pm} \Psi^\pm. \quad (4.10)$$

After some direct calculation, we find (*cf.* Métivier [34, Proposition 1.3.1]) that

$$\begin{aligned} \mathbb{L}'(\hat{U}^\pm, \hat{\Phi}^\pm)(V^\pm, \Psi^\pm) \\ = L(\hat{U}^\pm, \hat{\Phi}^\pm)\dot{V}^\pm + \mathcal{C}(\hat{U}^\pm, \hat{\Phi}^\pm)\dot{V}^\pm + \frac{\Psi^\pm}{\partial_2 \hat{\Phi}^\pm} \partial_2(L(\hat{U}^\pm, \hat{\Phi}^\pm)\hat{U}^\pm). \end{aligned} \quad (4.11)$$

In view of the nonlinear results obtained in [1, 19, 22], we neglect the zero-th order term in Ψ^\pm and consider the following *effective linear problem*:

$$\mathbb{L}'_e(\hat{U}^\pm, \hat{\Phi}^\pm)\dot{V}^\pm := L(\hat{U}^\pm, \hat{\Phi}^\pm)\dot{V}^\pm + \mathcal{C}(\hat{U}^\pm, \hat{\Phi}^\pm)\dot{V}^\pm = f^\pm \quad \text{if } x_2 > 0, \quad (4.12a)$$

$$\mathbb{B}'_e(\hat{U}^\pm, \hat{\Phi}^\pm)(\dot{V}, \psi) := \hat{b} \nabla \psi + b_\# \psi + \hat{B} \dot{V}|_{x_2=0} = g \quad \text{if } x_2 = 0, \quad (4.12b)$$

$$\Psi^+ = \Psi^- = \psi \quad \text{if } x_2 = 0, \quad (4.12c)$$

where $\mathcal{C}(\hat{U}^\pm, \hat{\Phi}^\pm)$, \hat{b} , and \hat{B} are defined by (4.5), (4.7), and (4.8) respectively, $\dot{V} := (\dot{V}^+, \dot{V}^-)^\top$, and

$$b_\#(t, x_1) := \hat{B}(t, x_1) \left(\frac{\partial_2 \hat{U}^+}{\partial_2 \hat{\Phi}^+}, \frac{\partial_2 \hat{U}^-}{\partial_2 \hat{\Phi}^-} \right)^\top |_{x_2=0}. \quad (4.13)$$

Note that $b_\#$ is a C^∞ -function of $(\hat{V}|_{x_2=0}, \partial_2 \hat{V}|_{x_2=0}, \nabla \hat{\varphi}, \partial_2 \hat{\Psi}|_{x_2=0})$ that vanishes at the origin. By virtue of (4.2), it follows that $\mathcal{C}(\hat{U}^\pm, \hat{\Phi}^\pm) \in W^{1,\infty}(\Omega)$, and the coefficients of operators $L(\hat{U}^\pm, \hat{\Phi}^\pm)$ are in $W^{2,\infty}(\Omega)$. We observe that the trace of vector $\hat{B} \dot{V}$ involved in boundary conditions (4.12b) depends solely on the traces of $\mathbb{P}^+(\hat{\varphi}) \dot{V}^+$ and $\mathbb{P}^-(\hat{\varphi}) \dot{V}^-$ on $\{x_2 = 0\}$, where $\mathbb{P}^\pm(\hat{\varphi})$ are defined as

$$\mathbb{P}^\pm(\hat{\varphi})V := (V_1, \hat{\zeta}_\pm|_{x_2=0} V_3 - \hat{\varrho}_\pm|_{x_2=0} V_2)^\top, \quad (4.14)$$

with $\hat{\varrho}_\pm$ and $\hat{\zeta}_\pm$ defined by (4.9). We will consider the dropped term in (4.11) as an error term at each Nash–Moser iteration step in the subsequent nonlinear analysis.

4.2 Reformulation

It is more convenient to transform the linearized problem (4.12) into a problem with a constant and diagonal boundary matrix. This is possible because the boundary matrix for (4.12) has constant rank on the whole closed half-space $\{x_2 \geq 0\}$.

Let us calculate the eigenvalues and the corresponding eigenvectors of the boundary matrix for (4.12). Using constraint (4.3a) reduces the coefficient matrices $\tilde{A}_2(\mathring{U}^\pm, \mathring{\Phi}^\pm)$ to:

$$\tilde{A}_2(\mathring{U}^\pm, \mathring{\Phi}^\pm) = \frac{1}{\partial_2 \mathring{\Phi}^\pm} \begin{pmatrix} 0 & -\mathring{N}^\pm \mathring{c}_\pm^2 \mathring{\rho}_\pm & \mathring{N}^\pm \mathring{c}_\pm^2 \mathring{\varsigma}_\pm \\ -\mathring{\rho}_\pm / \mathring{N}^\pm & 0 & 0 \\ \mathring{\varsigma}_\pm / \mathring{N}^\pm & 0 & 0 \end{pmatrix}. \quad (4.15)$$

After a direct calculation, we obtain that the eigenvalues are

$$\lambda_1 = 0, \quad \lambda_2 = -\frac{\mathring{c}_\pm \sqrt{\mathring{\rho}_\pm^2 + \mathring{\varsigma}_\pm^2}}{\partial_2 \mathring{\Phi}^\pm}, \quad \lambda_3 = \frac{\mathring{c}_\pm \sqrt{\mathring{\rho}_\pm^2 + \mathring{\varsigma}_\pm^2}}{\partial_2 \mathring{\Phi}^\pm},$$

with corresponding eigenvectors

$$r_1 = \begin{pmatrix} 0 \\ \mathring{\varsigma}_\pm \\ \mathring{\rho}_\pm \end{pmatrix}, \quad r_2 = \begin{pmatrix} \sqrt{\mathring{\rho}_\pm^2 + \mathring{\varsigma}_\pm^2} \\ \mathring{\rho}_\pm / (\mathring{N}^\pm \mathring{c}_\pm) \\ -\mathring{\varsigma}_\pm / (\mathring{N}^\pm \mathring{c}_\pm) \end{pmatrix}, \quad r_3 = \begin{pmatrix} \sqrt{\mathring{\rho}_\pm^2 + \mathring{\varsigma}_\pm^2} \\ -\mathring{\rho}_\pm / (\mathring{N}^\pm \mathring{c}_\pm) \\ \mathring{\varsigma}_\pm / (\mathring{N}^\pm \mathring{c}_\pm) \end{pmatrix}.$$

Define the matrices

$$R(\mathring{U}^\pm, \mathring{\Phi}^\pm) := \begin{pmatrix} 0 & \sqrt{\mathring{\rho}_\pm^2 + \mathring{\varsigma}_\pm^2} & \sqrt{\mathring{\rho}_\pm^2 + \mathring{\varsigma}_\pm^2} \\ \mathring{\varsigma}_\pm & \mathring{\rho}_\pm / (\mathring{N}^\pm \mathring{c}_\pm) & -\mathring{\rho}_\pm / (\mathring{N}^\pm \mathring{c}_\pm) \\ \mathring{\rho}_\pm & -\mathring{\varsigma}_\pm / (\mathring{N}^\pm \mathring{c}_\pm) & \mathring{\varsigma}_\pm / (\mathring{N}^\pm \mathring{c}_\pm) \end{pmatrix}, \quad (4.16)$$

and $\tilde{A}_0(\mathring{U}^\pm, \mathring{\Phi}^\pm) := \text{diag}(1, \lambda_2^{-1}, \lambda_3^{-1})$. Then it follows that

$$\tilde{A}_0 R^{-1} \tilde{A}_2 R(\mathring{U}^\pm, \mathring{\Phi}^\pm) = \mathbf{I}_2 := \text{diag}(0, 1, 1).$$

We thus perform the transformation:

$$W^\pm := R^{-1}(\mathring{U}^\pm, \mathring{\Phi}^\pm) \dot{V}^\pm. \quad (4.17)$$

Multiplying (4.12a) by matrices $\tilde{A}_0 R^{-1}(\mathring{U}^\pm, \mathring{\Phi}^\pm)$ yields the equivalent system of (4.12a):

$$\mathbf{A}_0^\pm \partial_t W^\pm + \mathbf{A}_1^\pm \partial_1 W^\pm + \mathbf{I}_2 \partial_2 W^\pm + \mathbf{C}^\pm W^\pm = F^\pm, \quad (4.18)$$

where $F^\pm := \tilde{A}_0 R^{-1}(\mathring{U}^\pm, \mathring{\Phi}^\pm) f^\pm$, and

$$\begin{aligned} \mathbf{A}_0^\pm &:= \tilde{A}_0 R^{-1} A_0 R(\mathring{U}^\pm, \mathring{\Phi}^\pm), & \mathbf{A}_1^\pm &:= \tilde{A}_0 R^{-1} A_1 R(\mathring{U}^\pm, \mathring{\Phi}^\pm), \\ \mathbf{C}^\pm &:= \tilde{A}_0 \left(R^{-1} A_0 \partial_t R + R^{-1} A_1 \partial_1 R + R^{-1} \tilde{A}_2 \partial_2 R + R^{-1} C R \right) (\mathring{U}^\pm, \mathring{\Phi}^\pm). \end{aligned} \quad (4.19)$$

Matrices \mathbf{A}_0^\pm and \mathbf{A}_1^\pm belong to $W^{2,\infty}(\Omega)$, while matrices \mathbf{C}^\pm are in $W^{1,\infty}(\Omega)$. Moreover, \mathbf{A}_0^\pm and \mathbf{A}_1^\pm are C^∞ -functions of their arguments $(\mathring{V}^\pm, \nabla \mathring{\Psi}^\pm)$, and \mathbf{C}^\pm are C^∞ -functions of their arguments $(\mathring{V}^\pm, \nabla \mathring{V}^\pm, \nabla \mathring{\Psi}^\pm, \nabla^2 \mathring{\Psi}^\pm)$. Under transformation (4.17), the boundary conditions (4.12b)–(4.12c) become

$$\mathcal{B}^\gamma(W, \psi) := \mathring{b} \nabla \psi + b_\sharp \psi + \mathbf{B}W = g \quad \text{if } x_2 = 0, \quad (4.20a)$$

$$\Psi^+ = \Psi^- = \psi \quad \text{if } x_2 = 0, \quad (4.20b)$$

where \mathring{b} and b_\sharp are given by (4.7) and (4.13) respectively, $W := (W^+, W^-)^\top$, and

$$\begin{aligned} \mathbf{B}(t, x_1) &:= \mathring{B} \begin{pmatrix} R(\mathring{U}^+, \mathring{\Phi}^+) & 0 \\ 0 & R(\mathring{U}^-, \mathring{\Phi}^-) \end{pmatrix} \Big|_{x_2=0} \\ &= \begin{pmatrix} 0 & m_1^+ + m_2^+ & m_1^+ - m_2^+ & 0 & -m_1^- - m_2^- & -m_1^- + m_2^- \\ 0 & m_1^+ + m_2^+ & m_1^+ - m_2^+ & 0 & 0 & 0 \\ 0 & \sqrt{\mathring{\varrho}_+^2 + \mathring{\varsigma}_+^2} & \sqrt{\mathring{\varrho}_+^2 + \mathring{\varsigma}_+^2} & 0 & -\sqrt{\mathring{\varrho}_-^2 + \mathring{\varsigma}_-^2} & -\sqrt{\mathring{\varrho}_-^2 + \mathring{\varsigma}_-^2} \end{pmatrix} \Big|_{x_2=0}. \end{aligned} \quad (4.21)$$

In the last expression, for notational simplicity, we have introduced m_j^\pm as

$$m_1^\pm := \frac{\epsilon^2 \partial_t \mathring{\Phi}^\pm \sqrt{\mathring{\varrho}_\pm^2 + \mathring{\varsigma}_\pm^2}}{\mathring{I}_\pm^2 \mathring{h}_\pm \mathring{N}^\pm}, \quad m_2^\pm := \frac{\mathring{\varrho}_\pm^2 + \mathring{\varsigma}_\pm^2}{\mathring{I}_\pm^2 \mathring{c}_\pm \mathring{h}_\pm \mathring{N}^\pm}. \quad (4.22)$$

It is clear that matrix \mathbf{B} is a C^∞ -function of $(\mathring{V}|_{x_2=0}, \nabla \mathring{\varphi})$. According to (4.14) and (4.17), we have

$$\mathbb{P}^\pm(\mathring{\varphi}) \mathring{V}^\pm|_{x_2=0} = \begin{pmatrix} \sqrt{\mathring{\varrho}_\pm^2 + \mathring{\varsigma}_\pm^2} (W_2^\pm + W_3^\pm) \\ -\frac{\mathring{\varrho}_\pm^2 + \mathring{\varsigma}_\pm^2}{\mathring{N}^\pm \mathring{c}_\pm} (W_2^\pm - W_3^\pm) \end{pmatrix} \Big|_{x_2=0}. \quad (4.23)$$

We find that the trace of vector $\mathbf{B}W$ involved in boundary conditions (4.20) depends only on the traces of the noncharacteristic part of vector W , *i.e.* sub-vector $W^{\text{nc}} := (W_2^+, W_3^+, W_2^-, W_3^-)^\top$.

5 Basic Energy Estimate for the Linearized Problem

In this section, we are going to prove the following theorem, which provides the basic energy estimate for the effective linear problem (4.12).

Theorem 5.1 *Assume that the stationary solution (2.24) satisfies (2.10) and (2.25). Assume further that the basic state $(\mathring{U}^\pm, \mathring{\Phi}^\pm)$ satisfies (4.1)–(4.3). Then there exist constants $K_0 > 0$ and $\gamma_0 \geq 1$ such that, if $K \leq K_0$ and $\gamma \geq \gamma_0$, then, for all $(\mathring{V}, \psi) \in H_\gamma^2(\Omega) \times H_\gamma^2(\mathbb{R}^2)$, the following estimate holds:*

$$\begin{aligned} &\gamma \|\mathring{V}\|_{L_\gamma^2(\Omega)}^2 + \|\mathbb{P}^\pm(\mathring{\varphi}) \mathring{V}^\pm|_{x_2=0}\|_{L_\gamma^2(\mathbb{R}^2)}^2 + \|\psi\|_{H_\gamma^1(\mathbb{R}^2)}^2 \\ &\lesssim \gamma^{-3} \|\mathbb{L}'_e(\mathring{U}^\pm, \mathring{\Phi}^\pm) \mathring{V}^\pm|_{L^2(H_\gamma^1)}\|^2 + \gamma^{-2} \|\mathbb{B}'_e(\mathring{U}^\pm, \mathring{\Phi}^\pm)(\mathring{V}|_{x_2=0}, \psi)\|_{H_\gamma^1(\mathbb{R}^2)}^2. \end{aligned} \quad (5.1)$$

Remark 5.1 Since the Lopatinskiĭ determinant associated with problem (4.12) admits the roots on the boundary of frequency space, the energy estimate (5.1) has a loss of regularity of the solution with respect to the source terms. Furthermore, there is a loss of control on the traces of the solution in (5.1), which is mainly owing to the fact that (4.12) is a characteristic boundary problem.

We notice that systems (4.12a) are symmetrizable hyperbolic with the Friedrichs symmetrizers $S_2(\dot{U}^\pm)$ for operators $\mathbb{L}'_e(\dot{U}^\pm, \dot{\Phi}^\pm)$, where function $S_2(U)$ is defined in (A.10). By virtue of (4.3a), we compute

$$\begin{aligned} & S_2(\dot{U}^\pm) \tilde{A}_2(\dot{U}^\pm, \dot{\Phi}^\pm) \\ &= \frac{1}{\partial_2 \dot{\Phi}^\pm} S_2(\dot{U}^\pm) (A_2(\dot{U}^\pm) - \partial_t \dot{\Phi}^\pm A_0(\dot{U}^\pm) - \partial_1 \dot{\Phi}^\pm A_1(\dot{U}^\pm)) \\ &= \frac{1}{\partial_2 \dot{\Phi}^\pm} \begin{pmatrix} 0 & -\dot{N}^\pm \dot{c}_\pm^2 \dot{\varrho}_\pm & \dot{N}^\pm \dot{c}_\pm^2 \dot{\varsigma}_\pm \\ -\dot{N}^\pm \dot{c}_\pm^2 \dot{\varrho}_\pm & 0 & 0 \\ \dot{N}^\pm \dot{c}_\pm^2 \dot{\varsigma}_\pm & 0 & 0 \end{pmatrix}, \end{aligned}$$

where $\dot{\varrho}_\pm$ and $\dot{\varsigma}_\pm$ are defined in (4.9). Multiplying (4.12a) by the Friedrichs symmetrizers $S_2(\dot{U}^\pm)$ and employing integration by parts yield the following lemma:

Lemma 5.1 *There exists a constant $\gamma_0 \geq 1$ such that, for all $\gamma \geq \gamma_0$, the following estimate holds:*

$$\gamma \|\dot{V}^\pm\|_{L^2_\gamma(\Omega)}^2 \lesssim \gamma^{-1} \|\mathbb{L}'_e(\dot{U}^\pm, \dot{\Phi}^\pm) \dot{V}^\pm\|_{L^2_\gamma(\Omega)}^2 + \|\mathbb{P}^\pm(\dot{\varphi}) \dot{V}^\pm|_{x_2=0}\|_{L^2_\gamma(\mathbb{R}^2)}^2.$$

To prove Theorem 5.1, it remains to deduce the desired energy estimate for the discontinuity front ψ and the traces of $\mathbb{P}^\pm(\dot{\varphi}) \dot{V}^\pm$ on $\{x_2 = 0\}$ in terms of the source terms in the interior domain and on the boundary.

Introducing $\widetilde{W}^\pm := e^{-\gamma t} W^\pm$, system (4.18) equivalently reads

$$\begin{aligned} & \mathcal{L}_\pm^\gamma \widetilde{W}^\pm \\ &:= \gamma \mathbf{A}_0^\pm \widetilde{W}^\pm + \mathbf{A}_0^\pm \partial_t \widetilde{W}^\pm + \mathbf{A}_1^\pm \partial_1 \widetilde{W}^\pm + \mathbf{I}_2 \partial_2 \widetilde{W}^\pm + \mathbf{C}^\pm \widetilde{W}^\pm = e^{-\gamma t} F^\pm. \end{aligned} \quad (5.2)$$

We also introduce $\widetilde{W} := (\widetilde{W}^+, \widetilde{W}^-)^\top$, $\widetilde{\Psi}^\pm := e^{-\gamma t} \Psi^\pm$, and $\widetilde{\psi} := e^{-\gamma t} \psi$. Then the boundary conditions (4.20) are equivalent to

$$\mathcal{B}^\gamma(\widetilde{W}, \widetilde{\psi}) := \gamma \mathbf{b}_0 \widetilde{\psi} + \mathbf{b} \nabla \widetilde{\psi} + \mathbf{b}_\sharp \widetilde{\psi} + \mathbf{B} \widetilde{W} = e^{-\gamma t} g \quad \text{if } x_2 = 0, \quad (5.3a)$$

$$\widetilde{\Psi}^+ = \widetilde{\Psi}^- = \widetilde{\psi} \quad \text{if } x_2 = 0, \quad (5.3b)$$

where $\mathbf{b}_0 := (0, 1, 0)^\top$. In view of (4.23), we obtain the estimate:

$$\|\mathbb{P}^\pm(\dot{\varphi}) \dot{V}^\pm|_{x_2=0}\|_{L^2_\gamma(\mathbb{R}^2)} \lesssim \|W^{\text{nc}}|_{x_2=0}\|_{L^2_\gamma(\mathbb{R}^2)} \lesssim \|\widetilde{W}^{\text{nc}}|_{x_2=0}\|, \quad (5.4)$$

where $\widetilde{W}^{\text{nc}} := (\widetilde{W}_2^+, \widetilde{W}_3^+, \widetilde{W}_2^-, \widetilde{W}_3^-)^\top$. By virtue of (5.4) and Lemma 5.1, we obtain that Theorem 5.1 admits the following equivalent proposition.

Proposition 5.1 *Assume that the stationary solution (2.24) satisfies (2.10) and (2.25). Assume further that the basic state $(\tilde{U}^\pm, \tilde{\Phi}^\pm)$ satisfies (4.1)–(4.3). Then there exist some constants $K_0 > 0$ and $\gamma_0 \geq 1$ such that, if $K \leq K_0$ and $\gamma \geq \gamma_0$, then the following estimate holds: for all $(W, \psi) \in H^2(\Omega) \times H^2(\mathbb{R}^2)$,*

$$\|W^{\text{nc}}|_{x_2=0}\|^2 + \|\psi\|_{1,\gamma}^2 \lesssim \gamma^{-3} \|\mathcal{L}_\pm^\gamma W^\pm\|_{1,\gamma}^2 + \gamma^{-2} \|\mathcal{B}^\gamma(W^{\text{nc}}|_{x_2=0}, \psi)\|_{1,\gamma}^2, \quad (5.5)$$

where operators \mathcal{L}_\pm^γ and \mathcal{B}^γ are given by (5.2) and (5.3a), respectively.

In the rest of this section, we give the proof of Proposition 5.1.

5.1 Parilinearization

We now perform the parilinearization of the interior equations and the boundary conditions.

5.1.1 Some results on paradifferential calculus

For self-containedness, we list some definitions and results about paradifferential calculus with a parameter that will be used in this paper. See [4, Appendix C] and the references cited therein for the rigorous proofs.

Definition 5.1 For any $m \in \mathbb{R}$ and $k \in \mathbb{N}$, we define the following:

- (i) A function $a(x, \xi, \gamma) : \mathbb{R}^2 \times \mathbb{R}^2 \times [1, \infty) \rightarrow \mathbb{C}^{N \times N}$ is called a paradifferential symbol of degree m and regularity k if a is C^∞ in ξ and, for each $\alpha \in \mathbb{N}^2$, there exists a positive constant C_α such that

$$\|\partial_\xi^\alpha a(\cdot, \xi, \gamma)\|_{W^{k,\infty}(\mathbb{R}^2)} \leq C_\alpha \lambda^{m-|\alpha|,\gamma}(\xi) \quad \text{for all } (\xi, \gamma) \in \mathbb{R}^2 \times [1, \infty),$$

where $\lambda^{s,\gamma}(\xi) := (\gamma^2 + |\xi|^2)^{s/2}$ for $s \in \mathbb{R}$.

- (ii) Γ_k^m denotes the set of paradifferential symbols of degree m and regularity k . We denote by α_m a generic symbol in the class Γ_1^m .
- (iii) We say that a family of operators $\{P^\gamma\}_{\gamma \geq 1}$ is of order $\leq m$, if, for every $s \in \mathbb{R}$ and $\gamma \geq 1$, there exists a constant $C(s, m)$ independent of γ such that

$$\|P^\gamma u\|_{s,\gamma} \leq C(s, m) \|u\|_{s+m,\gamma} \quad \text{for all } u \in H^{s+m}.$$

We use \mathcal{R}_m to denote a generic family of operators of order $\leq m$.

- (iv) For $s \in \mathbb{R}$, operator $A^{s,\gamma}$ is defined in such a way that

$$A^{s,\gamma} u(x) := \frac{1}{(2\pi)^2} \int_{\mathbb{R}^2} e^{ix \cdot \xi} \lambda^{s,\gamma}(\xi) \hat{u}(\xi) d\xi$$

for all u in the Schwartz class \mathcal{S} .

- (v) To any symbol $a \in \Gamma_0^m$, we associate the family of paradifferential operators $\{T_a^\gamma\}_{\gamma \geq 1}$ defined in such a way that

$$T_a^\gamma u(x) := \frac{1}{(2\pi)^2} \int_{\mathbb{R}^2} \int_{\mathbb{R}^2} e^{ix \cdot \xi} K^\psi(x-y, \xi, \gamma) a(y, \xi, \gamma) \hat{u}(\xi) dy d\xi$$

for all $u \in \mathcal{S}$. In the last expression, $K^\psi(\cdot, \xi, \gamma)$ is the inverse Fourier transform of $\psi(\cdot, \xi, \gamma)$ with ψ given by

$$\psi(x, \xi, \gamma) := \sum_{q \in \mathbb{N}} \chi(2^{2^{-1}x}, 0) \phi(2^{-q}\xi, 2^{-q}\gamma),$$

where $\phi(\xi, \gamma) := \chi(2^{-1}\xi, 2^{-1}\gamma) - \chi(\xi, \gamma)$, and χ is a C^∞ -function on \mathbb{R}^3 such that

$$\chi(z) \geq \chi(z') \quad \text{if } |z| \leq |z'|, \quad \chi(z) = \begin{cases} 1 & \text{if } |z| \leq \frac{1}{2}, \\ 0 & \text{if } |z| \geq 1. \end{cases}$$

Lemma 5.2 *The following statements hold:*

- (i) *If $a \in W^{1,\infty}(\mathbb{R}^2)$, $u \in L^2(\mathbb{R}^2)$, and $\gamma \geq 1$, then*

$$\gamma \|au - T_a^\gamma u\| + \|a\partial_j u - T_{i\xi_j a}^\gamma u\| + \|au - T_a^\gamma u\|_{1,\gamma} \lesssim \|a\|_{W^{1,\infty}(\mathbb{R}^2)} \|u\|.$$

- (ii) *If $a \in W^{2,\infty}(\mathbb{R}^2)$, $u \in L^2(\mathbb{R}^2)$, and $\gamma \geq 1$, then*

$$\gamma \|au - T_a^\gamma u\|_{1,\gamma} + \|a\partial_j u - T_{i\xi_j a}^\gamma u\|_{1,\gamma} \lesssim \|a\|_{W^{2,\infty}(\mathbb{R}^2)} \|u\|.$$

- (iii) *If $a \in \Gamma_k^m$, then T_a^γ is of order $\leq m$. In particular, if $a \in L^\infty(\mathbb{R}^2)$ is independent of ξ , then*

$$\|T_a^\gamma u\|_{s,\gamma} \lesssim \|a\|_{L^\infty(\mathbb{R}^2)} \|u\|_{s,\gamma} \quad \text{for all } s \in \mathbb{R}, u \in H^s(\mathbb{R}^2).$$

- (iv) *If $a \in \Gamma_1^m$ and $b \in \Gamma_1^{m'}$, then product $ab \in \Gamma_1^{m+m'}$, family $\{T_a^\gamma T_b^\gamma - T_{ab}^\gamma\}_{\gamma \geq 1}$ is of order $\leq m+m'-1$, and family $\{(T_a^\gamma)^* - T_{a^*}^\gamma\}_{\gamma \geq 1}$ is of order $\leq m-1$.*

- (v) *If $a \in \Gamma_2^m$ and $b \in \Gamma_2^{m'}$, then $\{T_a^\gamma T_b^\gamma - T_{ab}^\gamma - T_{-i \sum_j \partial_{\xi_j} a \partial_{x_j} b}\}_{\gamma \geq 1}$ is of order $\leq m+m'-2$.*

- (vi) *Gårding's inequality: If $a \in \Gamma_1^{2m}$ is a square matrix symbol that satisfies*

$$\operatorname{Re} a(x, \xi, \gamma) \geq c(\gamma^2 + |\xi|^2)^m I \quad \text{for all } (x, \xi, \gamma) \in \mathbb{R}^4 \times [1, \infty)$$

for some constant c , then there exists $\gamma_0 \geq 1$ such that

$$\operatorname{Re} \langle T_a^\gamma u, u \rangle \geq \frac{c}{4} \|u\|_{m,\gamma}^2 \quad \text{for all } u \in H^m(\mathbb{R}^2) \text{ and } \gamma \geq \gamma_0.$$

(vii) *Microlocalized Gårding's inequality:* Let $a \in \Gamma_1^{2m}$ be a square matrix symbol and $\chi \in \Gamma_1^0$. If there exist a scalar real symbol $\tilde{\chi} \in \Gamma_1^0$ and a constant $c > 0$ such that $\tilde{\chi} \geq 0$, $\chi\tilde{\chi} \equiv \chi$, and

$$\tilde{\chi}^2(x, \xi, \gamma) \operatorname{Re} a(x, \xi, \gamma) \geq c\tilde{\chi}^2(x, \xi, \gamma)(\gamma^2 + |\xi|^2)^m I$$

for all $(x, \xi, \gamma) \in \mathbb{R}^4 \times [1, \infty)$, then there exist $\gamma_0 \geq 1$ and $C > 0$ such that

$$\operatorname{Re} \langle T_a^\gamma T_\chi^\gamma u, T_\chi^\gamma u \rangle \geq \frac{c}{2} \|T_\chi^\gamma u\|_{m, \gamma}^2 - C \|u\|_{m-1, \gamma}^2$$

for all $u \in H^m(\mathbb{R}^2)$, $\gamma \geq \gamma_0$.

Here we have used the notation, $\operatorname{Re} B := (B + B^*)/2$, for any complex square matrix B with B^* being its conjugate transpose.

The reader may find the detailed proof of Lemma 5.2 (vii) in Métivier–Zumbrun [35, Theorem B.18].

5.1.2 Parilinearization of the interior equations

In view of (4.2) and (4.3b), we have the following estimate for the coefficients of \mathcal{L}_\pm^γ given in (5.2):

$$\|(\mathbf{A}_0^\pm, \mathbf{A}_1^\pm)\|_{W^{2, \infty}(\Omega)} + \|\mathbf{C}^\pm\|_{W^{1, \infty}(\Omega)} \leq C(K, \kappa_0).$$

It then follows from Lemma 5.2 (ii) that

$$\begin{aligned} \|\gamma \mathbf{A}_0^+ W^+ - T_{\gamma \mathbf{A}_0^+}^\gamma W^+\|_{1, \gamma}^2 &= \int_0^\infty \|\gamma \mathbf{A}_0^+ W^+(\cdot, x_2) - T_{\gamma \mathbf{A}_0^+}^\gamma W^+(\cdot, x_2)\|_{1, \gamma}^2 dx_2 \\ &\lesssim \int_0^\infty \|\mathbf{A}_0^+(\cdot, x_2)\|_{W^{2, \infty}(\mathbb{R}^2)}^2 \|W^+(\cdot, x_2)\|^2 dx_2 \\ &\leq C(K, \kappa_0) \|W^+\|^2. \end{aligned}$$

Similarly, we derive from Lemma 5.2 (i)–(ii) that

$$\begin{aligned} \|\mathbf{A}_0^+ \partial_t W^+ - T_{i\delta \mathbf{A}_0^+}^\gamma W^+\|_{1, \gamma} &\lesssim \|\mathbf{A}_0^+\|_{W^{2, \infty}(\Omega)} \|W^+\| \leq C(K, \kappa_0) \|W^+\|, \\ \|\mathbf{A}_1^+ \partial_1 W^+ - T_{i\eta \mathbf{A}_1^+}^\gamma W^+\|_{1, \gamma} &\lesssim \|\mathbf{A}_1^+\|_{W^{2, \infty}(\Omega)} \|W^+\| \leq C(K, \kappa_0) \|W^+\|, \\ \|\mathbf{C}^+ W^+ - T_{\mathbf{C}^+}^\gamma W^+\|_{1, \gamma} &\lesssim \|\mathbf{C}^+\|_{W^{1, \infty}(\Omega)} \|W^+\| \leq C(K, \kappa_0) \|W^+\|. \end{aligned}$$

Combining these estimates yields

$$\|\mathcal{L}_+^\gamma W^+ - \mathbf{I}_2 \partial_2 W^+ - T_{\tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+ + \mathbf{C}^+}^\gamma W^+\|_{1, \gamma} \leq C(K, \kappa_0) \|W^+\|, \quad (5.6)$$

where $\tau = \gamma + i\delta$, and \mathcal{L}_+^γ is the linearized operator defined by (5.2). We can also obtain the following estimate for the equations on W^- :

$$\|\mathcal{L}_-^\gamma W^- - \mathbf{I}_2 \partial_2 W^- - T_{\tau \mathbf{A}_0^- + i\eta \mathbf{A}_1^- + \mathbf{C}^-}^\gamma W^-\|_{1, \gamma} \leq C(K, \kappa_0) \|W^-\|. \quad (5.7)$$

The parilinearization for the interior equations is thus given as follows:

$$T_{\tau \mathbf{A}_0^\pm + i\eta \mathbf{A}_1^\pm + \mathbf{C}^\pm}^\gamma W^\pm + \mathbf{I}_2 \partial_2 W^\pm = \tilde{F}^\pm \quad \text{if } x_2 > 0. \quad (5.8)$$

Note that the above parilinearized equations do not involve the discontinuity function φ .

5.1.3 Paralinearization of the boundary conditions

According to (5.3a), we define

$$\begin{aligned} \mathbf{b}_0 &:= (0, 1, 0)^\top, \quad \mathbf{b}_1(t, x_1) := ((\hat{v}_1^+ - \hat{v}_1^-)|_{x_2=0}, \hat{v}_1^+|_{x_2=0}, 0)^\top, \\ \mathbf{b}(t, x_1, \delta, \eta, \gamma) &:= \tau \mathbf{b}_0 + i\eta \mathbf{b}_1(t, x_1) = (i\eta(\hat{v}_1^+ - \hat{v}_1^-), \tau + i\hat{v}_1^+ \eta, 0)^\top|_{x_2=0}. \end{aligned}$$

Since $\mathbf{b}_0, \mathbf{b}_1 \in W^{2,\infty}(\mathbb{R}^2)$, we obtain from Lemma 5.2 (iii) that

$$\begin{aligned} &\|\gamma \mathbf{b}_0 \psi + \hat{\mathbf{b}} \nabla \psi - T_{\mathbf{b}}^\gamma \psi\|_{1,\gamma} \\ &= \|\gamma \mathbf{b}_0 \psi + \mathbf{b}_0 \partial_t \psi + \mathbf{b}_1 \partial_1 \psi - T_{\mathbf{b}}^\gamma \psi\|_{1,\gamma} \\ &\lesssim \|(\mathbf{b}_0, \mathbf{b}_1)\|_{W^{2,\infty}(\mathbb{R}^2)} \|\psi\| \leq C(K) \|\psi\| \leq C(K) \gamma^{-1} \|\psi\|_{1,\gamma}. \end{aligned} \quad (5.9)$$

It follows from (4.2), (4.3b), and (4.13) that $\|\mathbf{b}_\sharp\|_{W^{1,\infty}(\mathbb{R}^2)} \leq C(K, \kappa_0)$. Employing Lemma 5.2 (ii)–(iii) yields

$$\begin{aligned} \|b_\sharp \psi\|_{1,\gamma} &\leq \|b_\sharp \psi - T_{b_\sharp}^\gamma \psi\|_{1,\gamma} + \|T_{b_\sharp}^\gamma \psi\|_{1,\gamma} \\ &\lesssim \|b_\sharp\|_{W^{1,\infty}(\mathbb{R}^2)} \|\psi\| + \|b_\sharp\|_{L^\infty(\mathbb{R}^2)} \|\psi\|_{1,\gamma} \leq C(K, \kappa_0) \|\psi\|_{1,\gamma}. \end{aligned} \quad (5.10)$$

In light of (4.21), we find that $\|\mathbf{B}\|_{W^{2,\infty}(\mathbb{R}^2)} \leq C(K, \kappa_0)$, and \mathbf{B} acts only on the noncharacteristic part W^{nc} of vector W . Hence, we deduce

$$\begin{aligned} &\|\mathbf{B}W|_{x_2=0} - T_{\mathbf{B}}^\gamma W|_{x_2=0}\|_{1,\gamma} \\ &\lesssim \gamma^{-1} \|\mathbf{B}\|_{W^{2,\infty}(\mathbb{R}^2)} \|W^{\text{nc}}|_{x_2=0}\| \leq C(K, \kappa_0) \gamma^{-1} \|W^{\text{nc}}|_{x_2=0}\|. \end{aligned} \quad (5.11)$$

Combine (5.9)–(5.11) together to find

$$\begin{aligned} &\|\mathcal{B}^\gamma(W|_{x_2=0}, \psi) - T_{\mathbf{b}}^\gamma \psi - T_{\mathbf{B}}^\gamma W|_{x_2=0}\|_{1,\gamma} \\ &\leq C(K, \kappa_0) (\|\psi\|_{1,\gamma} + \gamma^{-1} \|W^{\text{nc}}|_{x_2=0}\|). \end{aligned} \quad (5.12)$$

The paralinearization of the boundary conditions (5.3a) is then given as follows:

$$T_{\mathbf{b}}^\gamma \psi + T_{\mathbf{B}}^\gamma W = G \quad \text{if } x_2 = 0. \quad (5.13)$$

5.1.4 Eliminating the front

We can eliminate front ψ from the paralinearized boundary conditions (5.13) as in the constant coefficient case. For this purpose, we first notice that symbol \mathbf{b} is elliptic, which means that, for any $(t, x_1, \delta, \eta, \gamma) \in \mathbb{R}^4 \times (0, \infty)$,

$$|\mathbf{b}(t, x_1, \delta, \eta, \gamma)|^2 \geq c(K)(\gamma^2 + \delta^2 + \eta^2). \quad (5.14)$$

To show this estimate, by observing that \mathbf{b} is homogeneous of degree 1 with respect to (τ, η) and that Ξ_1 is compact, we only need to prove

$$|\mathbf{b}(t, x_1, \delta, \eta, \gamma)|^2 > 0 \quad \text{on } \Xi_1.$$

This estimate follows from the similar property for the constant coefficient case, by taking the perturbation, \hat{V} , small enough in $L^\infty(\Omega)$.

Using (5.14) and the Gårding inequality (Lemma 5.2 (vi)), we have

$$\operatorname{Re} \langle T_{\mathbf{b}^* \mathbf{b}}^\gamma \psi, \psi \rangle \geq c(K) \|\psi\|_{1,\gamma}^2 \quad \text{for all } \gamma \geq \gamma_0,$$

where γ_0 depends only on K . Since $\mathbf{b} \in \Gamma_2^1$, the operator:

$$T_{\mathbf{b}^* \mathbf{b}}^\gamma - (T_{\mathbf{b}}^\gamma)^* T_{\mathbf{b}}^\gamma = T_{\mathbf{b}^* \mathbf{b}}^\gamma - T_{\mathbf{b}^*}^\gamma T_{\mathbf{b}}^\gamma + \{T_{\mathbf{b}^*}^\gamma - (T_{\mathbf{b}}^\gamma)^*\} T_{\mathbf{b}}^\gamma$$

is of order ≤ 1 . Then

$$\|\psi\|_{1,\gamma}^2 \leq C(K) (\|T_{\mathbf{b}}^\gamma \psi\|^2 + \|\psi\|_{1,\gamma} \|\psi\|) \leq C(K) (\|T_{\mathbf{b}}^\gamma \psi\|^2 + \gamma^{-1} \|\psi\|_{1,\gamma}^2),$$

from which we take γ sufficiently large to derive

$$\|\psi\|_{1,\gamma} \leq C(K) \|T_{\mathbf{b}}^\gamma \psi\|.$$

Since the first and fourth columns of $\mathbf{B} \in W^{2,\infty}(\mathbb{R}^2)$ vanish, we apply Lemma 5.2 to obtain

$$\begin{aligned} \|\psi\|_{1,\gamma} &\leq C(K) (\|T_{\mathbf{b}}^\gamma \psi + T_{\mathbf{B}}^\gamma W|_{x_2=0}\| + \|W^{\text{nc}}|_{x_2=0}\|) \\ &\leq C(K) (\gamma^{-1} \|T_{\mathbf{b}}^\gamma \psi + T_{\mathbf{B}}^\gamma W|_{x_2=0}\|_{1,\gamma} + \|W^{\text{nc}}|_{x_2=0}\|). \end{aligned} \quad (5.15)$$

Combine this estimate with (5.12) and let γ large enough to deduce

$$\|\psi\|_{1,\gamma} \leq C(K) (\gamma^{-1} \|\mathcal{B}^\gamma(W|_{x_2=0}, \psi)\|_{1,\gamma} + \|W^{\text{nc}}|_{x_2=0}\|). \quad (5.16)$$

This last estimate indicates that it only remains to deduce an estimate of $W^{\text{nc}}|_{x_2=0}$ in terms of the source terms.

To eliminate ψ in the boundary conditions (5.13), we define the matrix:

$$Q(t, x_1, \delta, \eta, \gamma) := \left(\begin{array}{ccc} 0 & 0 & 1 \\ \tau + i\eta \hat{v}_1^+ & -i\eta(\hat{v}_1^+ - \hat{v}_1^-) & 0 \end{array} \right) \Big|_{x_2=0} \quad \text{for all } (\tau, \eta) \in \Xi_1.$$

Then we extend Q as a homogeneous mapping of degree 0 with respect to (τ, η) on Ξ . It follows that $Q \in \Gamma_2^0$ and $Q\mathbf{b} \equiv 0$. We define symbol β as

$$\beta(t, x_1, \delta, \eta, \gamma) := Q(t, x_1, \delta, \eta, \gamma) \mathbf{B}(t, x_1) \in \Gamma_2^0$$

for all $(t, x_1, \delta, \eta, \gamma) \in \mathbb{R}^4 \times \mathbb{R}_+$. After a direct calculation, we find that the first and fourth columns of β vanish, so that we consider β as a matrix with only four columns and two rows. More precisely, for all $(\tau, \eta) \in \Xi_1$, symbol β is given by

$$\begin{aligned} &\beta(t, x_1, \delta, \eta, \gamma) \\ &= \left(\begin{array}{cccc} \sqrt{\hat{\varrho}_+^2 + \xi_+^2} & \sqrt{\hat{\varrho}_+^2 + \xi_+^2} & -\sqrt{\hat{\varrho}_-^2 + \xi_-^2} & -\sqrt{\hat{\varrho}_-^2 + \xi_-^2} \\ \hat{a}_-(m_1^+ + m_2^+) & \hat{a}_-(m_1^+ - m_2^+) & -\hat{a}_+(m_1^- + m_2^-) & -\hat{a}_+(m_1^- - m_2^-) \end{array} \right) \Big|_{x_2=0}, \end{aligned} \quad (5.17)$$

where $\hat{\varrho}_\pm$ and $\hat{\varsigma}_\pm$ are given in (4.9), m_1^\pm and m_2^\pm are given in (4.22), and

$$\hat{a}_\pm := \tau + i\hat{v}_1^\pm(t, x)\eta. \quad (5.18)$$

Since $\mathbf{B} \in \Gamma_2^0$, $\mathbf{b} \in \Gamma_2^1$, and $Q\mathbf{b} \equiv 0$, we find from (5.12) and Lemma 5.2 that

$$\begin{aligned} & \|T_{\hat{\beta}}^\gamma W^{\text{nc}}|_{x_2=0}\|_{1,\gamma} \\ &= \|T_{Q\mathbf{B}}^\gamma W|_{x_2=0} - T_Q^\gamma T_{\mathbf{B}}^\gamma W|_{x_2=0} + T_Q^\gamma (T_{\mathbf{B}}^\gamma W|_{x_2=0} + T_{\mathbf{b}}^\gamma \psi) - T_Q^\gamma T_{\mathbf{b}}^\gamma \psi\|_{1,\gamma} \\ &\lesssim \|W^{\text{nc}}|_{x_2=0}\| + \|T_{\mathbf{B}}^\gamma W|_{x_2=0} + T_{\mathbf{b}}^\gamma \psi\|_{1,\gamma} + \|T_Q^\gamma T_{\mathbf{b}}^\gamma \psi - T_{Q\mathbf{b}}^\gamma \psi\|_{1,\gamma} \\ &\lesssim \|W^{\text{nc}}|_{x_2=0}\| + \|\mathcal{B}^\gamma(W|_{x_2=0}, \psi)\|_{1,\gamma} + \|\psi\|_{1,\gamma}. \end{aligned} \quad (5.19)$$

In view of (5.8) and (5.13), we obtain the following parilinearized problem with reduced boundary conditions:

$$T_{\tau\mathbf{A}_0^+ + i\eta\mathbf{A}_1^+ + \mathbf{C}^+}^\gamma W^+ + \mathbf{I}_2 \partial_2 W^+ = F^+ \quad \text{if } x_2 > 0, \quad (5.20a)$$

$$T_{\tau\mathbf{A}_0^- + i\eta\mathbf{A}_1^- + \mathbf{C}^-}^\gamma W^- + \mathbf{I}_2 \partial_2 W^- = F^- \quad \text{if } x_2 > 0, \quad (5.20b)$$

$$T_{\hat{\beta}}^\gamma W^{\text{nc}} = G \quad \text{if } x_2 = 0. \quad (5.20c)$$

We can deduce the following proposition for problem (5.20) by using the error estimates (5.6)–(5.7), (5.12), (5.15)–(5.16), and (5.19) (see also [17, Proposition 5.3]).

Proposition 5.2 *If there exist constants $K_0 > 0$ and $\gamma_0 \geq 1$ such that solution W to the parilinearized problem (5.20) satisfies*

$$\|W^{\text{nc}}|_{x_2=0}\|^2 \lesssim \gamma^{-3} \|F^\pm\|_{1,\gamma}^2 + \gamma^{-2} \|G\|_{1,\gamma}^2, \quad (5.21)$$

for $K \leq K_0$ and $\gamma \geq \gamma_0$, then Proposition 5.1 holds.

5.2 A Reduced Problem

In order to derive the energy estimate (5.21), we now derive a problem for the noncharacteristic variables W^{nc} from (5.20). This is possible since the coefficient matrix $\mathbf{I}_2 = \text{diag}(0, 1, 1)$ has constant rank. For convenience, we write

$$\tau\mathbf{A}_0^\pm + i\eta\mathbf{A}_1^\pm =: (b_{ij}^\pm) \in \Gamma_2^1, \quad (5.22)$$

where $\mathbf{A}_k^\pm = (A_{k,\pm}^{ij})$, $k = 0, 1$, are defined by (4.19). In particular, we compute

$$A_{1,\pm}^{11} = \hat{v}_1^\pm A_{0,\pm}^{11}, \quad A_{0,\pm}^{11} = \hat{I}_\pm \left\{ 1 - \frac{\epsilon^2 (\hat{\varsigma}_\pm \hat{v}_1^\pm + \hat{\varrho}_\pm \hat{v}_2^\pm)^2}{\hat{\varsigma}_\pm^2 + \hat{\varrho}_\pm^2} \right\} =: \mathbb{F}_1^\pm \in \mathbb{R}, \quad (5.23)$$

from which we obtain

$$b_{11}^\pm := \tau A_{0,\pm}^{11} + i\eta A_{1,\pm}^{11} = \mathbb{F}_1^\pm (\tau + i\hat{v}_1^\pm \eta). \quad (5.24)$$

In view of (4.9), $\mathbb{F}_1^\pm = \bar{F}(1 - \epsilon^2 \bar{v}^2) > 0$ when $(\mathring{V}^\pm, \mathring{\Psi}^\pm) = 0$ (cf. (3.24)). We use the continuity of \mathbb{F}_1^\pm and take K in (4.2) small enough to derive that $\mathbb{F}_1^\pm > 0$ for all $(t, x) \in \bar{\Omega}$. As a consequence, we have

$$b_{11}^\pm = 0 \quad \text{if and only if} \quad \mathring{a}_\pm = \tau + i\mathring{v}_1^\pm \eta = 0. \quad (5.25)$$

To represent the characteristic variables W_1^\pm in terms of W^{nc} , the singular points (t, x, τ, η) that are given by (5.25) should be excluded. We thus introduce two C^∞ -functions χ_+ and $\tilde{\chi}_+$ defined on $\bar{\Omega} \times \Xi$ such that

- Both χ_+ and $\tilde{\chi}_+$ are homogeneous of degree zero with respect to $(\tau, \eta) \in \Xi$;
- For all $(t, x, \tau, \eta) \in \bar{\Omega} \times \Xi_1$,

$$0 \leq \chi_+(t, x, \tau, \eta) \leq \tilde{\chi}_+(t, x, \tau, \eta) \leq 1, \quad (5.26)$$

$$\tilde{\chi}_+ \equiv 1 \quad \text{on} \quad \text{supp} \chi_+, \quad \text{supp} \tilde{\chi}_+ \subset \{\mathring{a}_+(t, x, \tau, \eta) \neq 0\}. \quad (5.27)$$

Since $\tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+ \in \Gamma_2^1$, $\mathbf{C}^+ \in \Gamma_1^0$, and $\chi_+ \in \Gamma_k^0$ for all $k \in \mathbb{N}$, we find from Lemma 5.2 (iv)–(v) that

$$T_{\chi_+}^\gamma T_{\tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+ + \mathbf{C}^+}^\gamma = T_{\tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+}^\gamma T_{\chi_+}^\gamma + T_{-i\{\chi_+, \tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+\}}^\gamma + T_{\mathbf{C}^+}^\gamma T_{\chi_+}^\gamma + \mathcal{R}_{-1},$$

where $\{a, b\}$ denotes the Poisson bracket of a and b :

$$\{a, b\} := \frac{\partial a}{\partial \delta} \frac{\partial b}{\partial t} + \frac{\partial a}{\partial \eta} \frac{\partial b}{\partial x_1} - \frac{\partial a}{\partial t} \frac{\partial b}{\partial \delta} - \frac{\partial a}{\partial x_1} \frac{\partial b}{\partial \eta}. \quad (5.28)$$

Setting

$$\mathbf{w}^+ := T_{\chi_+}^\gamma W^+$$

and applying operator $T_{\chi_+}^\gamma$ to (5.20a), we obtain

$$T_{\tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+}^\gamma \mathbf{w}^+ + T_{\mathbf{C}^+}^\gamma \mathbf{w}^+ + \mathbf{I}_2 \partial_2 \mathbf{w}^+ = T_r^\gamma W^+ + T_{\chi_+}^\gamma F^+ + \mathcal{R}_{-1} W^+, \quad (5.29)$$

where $r = i\{\chi_+, \tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+\} + \partial_2 \chi_+ \mathbf{I}_2$. We will employ letter r_+ to denote a generic symbol that belongs to Γ_1^0 and vanishes on $\{\chi_+ \equiv 1\} \cup \{\chi_+ \equiv 0\}$. Since $b_{11}^+ \neq 0$ on $\text{supp} \tilde{\chi}_+$, we infer that $\frac{\tilde{\chi}_+}{b_{11}^+} \in \Gamma_2^{-1}$ and

$$T_{\tilde{\chi}_+/b_{11}^+}^\gamma T_{b_{1j}^+}^\gamma \mathbf{w}_j^+ = T_{\tilde{\chi}_+ b_{1j}^+/b_{11}^+}^\gamma \mathbf{w}_j^+ + T_{\alpha_{-1}}^\gamma \mathbf{w}_j^+ + \mathcal{R}_{-2} W^+.$$

Applying operator $T_{\tilde{\chi}_+/b_{11}^+}^\gamma$ to the first equation in (5.29) yields

$$\begin{aligned} & T_{\tilde{\chi}_+}^\gamma \mathbf{w}_1^+ + T_{\tilde{\chi}_+ b_{12}^+/b_{11}^+}^\gamma \mathbf{w}_2^+ + T_{\tilde{\chi}_+ b_{13}^+/b_{11}^+}^\gamma \mathbf{w}_3^+ \\ &= \sum_{j=1}^3 T_{\alpha_{-1}}^\gamma \mathbf{w}_j^+ + T_{r_+ \alpha_{-1}}^\gamma W^+ + T_{\alpha_{-1}}^\gamma T_{\chi_+}^\gamma F_1^+ + \mathcal{R}_{-2} W^+. \end{aligned} \quad (5.30)$$

By virtue of (5.27), we have the identities: $\tilde{\chi}_+\chi_+ \equiv \chi_+$ and

$$\frac{\partial \tilde{\chi}_+}{\partial \delta} \frac{\partial \chi_+}{\partial t} + \frac{\partial \tilde{\chi}_+}{\partial x_1} \frac{\partial \chi_+}{\partial \eta} \equiv 0,$$

which imply

$$T_{\tilde{\chi}_+}^\gamma \mathbf{w}_1^+ = T_{\tilde{\chi}_+}^\gamma T_{\chi_+}^\gamma W_1^+ = T_{\tilde{\chi}_+\chi_+}^\gamma W_1^+ + \mathcal{R}_{-2}W_1^+ = \mathbf{w}_1^+ + \mathcal{R}_{-2}W_1^+.$$

Plug this identity into (5.30) to obtain

$$\begin{aligned} \mathbf{w}_1^+ &= -T_{\tilde{\chi}_+b_{12}^+/b_{11}^+}^\gamma \mathbf{w}_2^+ - T_{\tilde{\chi}_+b_{13}^+/b_{11}^+}^\gamma \mathbf{w}_3^+ \\ &\quad + \sum_{j=1}^3 T_{\alpha_{-1}}^\gamma \mathbf{w}_j^+ + T_{r_+\alpha_{-1}}^\gamma W^+ + T_{\alpha_{-1}}^\gamma T_{\chi_+}^\gamma F_1^+ + \mathcal{R}_{-2}W^+. \end{aligned} \quad (5.31)$$

The second equation of (5.29) reads

$$\sum_{j=1}^3 T_{b_{2j}^+}^\gamma \mathbf{w}_j^+ + \sum_{j=1}^3 T_{\alpha_0}^\gamma \mathbf{w}_j^+ + \partial_2 \mathbf{w}_2^+ = T_{r_+}^\gamma W^+ + T_{\chi_+}^\gamma F_2^+ + \mathcal{R}_{-1}W^+.$$

Since $b_{j1}^+ \in \Gamma_2^1$ and $\tilde{\chi}_+b_{1j}^+/b_{11}^+ \in \Gamma_2^0$, we then apply operator $T_{b_{21}^+}^\gamma$ to expression (5.31) and obtain

$$\begin{aligned} T_{b_{21}^+}^\gamma \mathbf{w}_1^+ &= -T_{\tilde{\chi}_+b_{21}^+b_{12}^+/b_{11}^+}^\gamma \mathbf{w}_2^+ - T_{\tilde{\chi}_+b_{21}^+b_{13}^+/b_{11}^+}^\gamma \mathbf{w}_3^+ \\ &\quad + \sum_{j=1}^3 T_{\alpha_0}^\gamma \mathbf{w}_j^+ + T_{r_+}^\gamma W^+ + \mathcal{R}_0 T_{\chi_+}^\gamma F^+ + \mathcal{R}_{-1}W^+. \end{aligned}$$

Consequently, we have

$$\begin{aligned} \partial_2 \mathbf{w}_2^+ &= T_{\mathbb{A}_{\tilde{\chi}_+}^{11}}^\gamma \mathbf{w}_2^+ + T_{\mathbb{A}_{\tilde{\chi}_+}^{12}}^\gamma \mathbf{w}_3^+ \\ &\quad + \sum_{j=1}^3 T_{\alpha_0}^\gamma \mathbf{w}_j^+ + T_{r_+}^\gamma W^+ + \mathcal{R}_0 T_{\chi_+}^\gamma F^+ + \mathcal{R}_{-1}W^+, \end{aligned} \quad (5.32)$$

where

$$\mathbb{A}_{\tilde{\chi}_+}^{11} = -b_{22}^+ + \tilde{\chi}_+b_{21}^+b_{12}^+/b_{11}^+, \quad \mathbb{A}_{\tilde{\chi}_+}^{12} = -b_{23}^+ + \tilde{\chi}_+b_{21}^+b_{13}^+/b_{11}^+.$$

Note that \mathbf{w}_1^+ appears in a zero-th order term in (5.32). We thus apply $T_{\alpha_0}^\gamma$ to expression (5.31) and deduce

$$T_{\alpha_0}^\gamma \mathbf{w}_1^+ = T_{\alpha_0}^\gamma \mathbf{w}_2^+ + T_{\alpha_0}^\gamma \mathbf{w}_3^+ + T_{r_+\alpha_{-1}}^\gamma W^+ + \mathcal{R}_{-1} T_{\chi_+}^\gamma F^+ + \mathcal{R}_{-1}W^+,$$

which, together with (5.32), implies the following equation for \mathbf{w}_2^+ :

$$\begin{aligned} \partial_2 \mathbf{w}_2^+ &= T_{\mathbb{A}_{\tilde{\chi}_+}^{11}}^\gamma \mathbf{w}_2^+ + T_{\mathbb{A}_{\tilde{\chi}_+}^{12}}^\gamma \mathbf{w}_3^+ \\ &\quad + \sum_{j=2}^3 T_{\alpha_0}^\gamma \mathbf{w}_j^+ + T_{r_+}^\gamma W^+ + \mathcal{R}_0 T_{\chi_+}^\gamma F^+ + \mathcal{R}_{-1}W^+. \end{aligned} \quad (5.33)$$

In this equation, the first and zero-th order terms in \mathbf{w}_1^+ have been eliminated. Performing a similar computation to the third equation of (5.29), we obtain the following equations for $\mathbf{w}_+^{\text{nc}} := (\mathbf{w}_2^+, \mathbf{w}_3^+)^{\text{T}}$:

$$\partial_2 \mathbf{w}_+^{\text{nc}} = T_{\mathbb{A}_{\tilde{\chi}_+}}^\gamma \mathbf{w}_+^{\text{nc}} + T_{\mathbb{E}^+}^\gamma \mathbf{w}_+^{\text{nc}} + T_{r_+}^\gamma W^+ + \mathcal{R}_0 T_{\chi_+}^\gamma F^+ + \mathcal{R}_{-1} W^+, \quad (5.34)$$

where $\mathbb{E}^+ \in \Gamma_1^0$, symbol $r_+ \in \Gamma_1^0$ vanishes on region $\{\chi_+ \equiv 1\} \cup \{\chi_+ \equiv 0\}$, and

$$\mathbb{A}_{\tilde{\chi}_+} = \begin{pmatrix} \mathbb{A}_{\tilde{\chi}_+}^{11} & \mathbb{A}_{\tilde{\chi}_+}^{12} \\ \mathbb{A}_{\tilde{\chi}_+}^{21} & \mathbb{A}_{\tilde{\chi}_+}^{22} \end{pmatrix} \in \Gamma_2^1, \quad \mathbb{A}_{\tilde{\chi}_+}^{ij} = -b_{i+1,j+1}^+ + \tilde{\chi}_+ b_{i+1,1}^+ b_{1,j+1}^+ / b_{11}^+.$$

Let us define χ_- and $\tilde{\chi}_-$ as χ_+ and $\tilde{\chi}_+$ by changing index “+” into “−”. We set $\mathbf{w}^- := T_{\chi_-}^\gamma W^-$ and employ a similar analysis to find that $\mathbf{w}_-^{\text{nc}} := (\mathbf{w}_2^-, \mathbf{w}_3^-)^{\text{T}}$ satisfies the same system as (5.34) with index “+” replaced by “−”. Applying the rule of symbolic calculus (Lemma 5.2(iv)) to (5.20c) yields the boundary condition for $\mathbf{w}^{\text{nc}} := (\mathbf{w}_2^+, \mathbf{w}_3^+, \mathbf{w}_2^-, \mathbf{w}_3^-)^{\text{T}}$:

$$T_\beta^\gamma \mathbf{w}^{\text{nc}}|_{x_2=0} = G + \mathcal{R}_{-1} W^{\text{nc}}.$$

We combine this last relation with the systems for $\mathbf{w}_\pm^{\text{nc}}$ to obtain the reduced problem:

$$\begin{cases} \partial_2 \mathbf{w}^{\text{nc}} = T_{\mathbb{A}_r}^\gamma \mathbf{w}^{\text{nc}} + T_{\mathbb{E}}^\gamma \mathbf{w}^{\text{nc}} + T_r^\gamma W + \mathcal{R}_0 T_\chi^\gamma F + \mathcal{R}_{-1} W & \text{if } x_2 > 0, \\ T_\beta^\gamma \mathbf{w}^{\text{nc}} = G + \mathcal{R}_{-1} W^{\text{nc}} & \text{if } x_2 = 0, \end{cases} \quad (5.35)$$

where β is given by (5.17) for $(\tau, \eta) \in \Xi_1$. The symbol matrix $\mathbb{A}_r \in \Gamma_2^1$ is given by

$$\begin{cases} \mathbb{A}_r = \begin{pmatrix} \mathbb{A}_{\tilde{\chi}_+} & 0 \\ 0 & \mathbb{A}_{\tilde{\chi}_-} \end{pmatrix}, \quad \mathbb{A}_{\tilde{\chi}_\pm} = (\mathbb{A}_{\tilde{\chi}_\pm}^{ij}), \\ \text{with } \mathbb{A}_{\tilde{\chi}_\pm}^{ij} = -b_{i+1,j+1}^\pm + \tilde{\chi}_\pm b_{i+1,1}^\pm b_{1,j+1}^\pm / b_{11}^\pm. \end{cases} \quad (5.36)$$

Matrices \mathbb{E} and r both belong to Γ_1^0 and have the same block diagonal structure as \mathbb{A}_r . Moreover, symbol r vanishes on region $\{\chi_+ = \chi_- \equiv 1\} \cup \{\chi_+ = \chi_- \equiv 0\}$.

5.3 Microlocalization

We now construct the degenerate Kreiss’ symmetrizers that are microlocal (*i.e.* local in the frequency space) in order to derive our energy estimate. The whole space $\bar{\Omega} \times \Xi$ will be divided into three disjoint parts according to the poles of the “non-cutoff” symbol \mathbb{A} and the zeros of the associated Lopatinskiĭ determinant, where

$$\mathbb{A} = \begin{pmatrix} \mathbb{A}^+ & 0 \\ 0 & \mathbb{A}^- \end{pmatrix}, \quad \mathbb{A}^\pm = (a_{ij}^\pm), \quad a_{ij}^\pm = -b_{i+1,j+1}^\pm + b_{i+1,1}^\pm b_{1,j+1}^\pm / b_{11}^\pm. \quad (5.37)$$

Notice that $\mathbb{A}_{\tilde{\chi}_{\pm}} = \mathbb{A}^{\pm}$ in region $\{\tilde{\chi}_{\pm} \equiv 1\}$. In light of (5.25), we obtain that the poles of \mathbb{A} belong to the set $\Upsilon_p := \Upsilon_p^+ \cup \Upsilon_p^-$, with

$$\Upsilon_p^{\pm} := \{(t, x, \tau, \eta) \in \bar{\Omega} \times \Xi : \tau = -i\eta\dot{v}_1^{\pm}(t, x, \tau, \eta)\}.$$

For the eigenvalues and the stable subspace of $\mathbb{A}(t, x, \tau, \eta)$, we have the following lemma.

Lemma 5.3 *Assume that $(\dot{V}, \nabla\dot{\Psi})$ is sufficiently small in $W^{2,\infty}(\Omega)$.*

- (a) *If $(\tau, \eta) \in \Xi_1$ with $\operatorname{Re}\tau > 0$, then the eigenvalues of $\mathbb{A}^{\pm}(t, x, \tau, \eta)$ are roots ω of*

$$\left(\omega - \frac{a_{11}^{\pm} + a_{22}^{\pm}}{2}\right)^2 = (\dot{C}_0^{\pm})^2 \left\{ (\dot{C}_1^{\pm})^2 (\tau \pm i\dot{C}_2^{\pm}\eta)^2 + \eta^2 \right\}, \quad (5.38)$$

where $\dot{C}_j^{\pm}, j = 0, 1, 2$, are positive smooth functions of $(\dot{V}^{\pm}, \nabla\dot{\Psi}^{\pm})$ such that $\dot{C}_j^{\pm} = \bar{C}_j$ when $(\dot{V}^{\pm}, \dot{\Psi}^{\pm}) = 0$, with \bar{C}_j given by (3.31). Moreover, \mathbb{A}^{\pm} has a unique eigenvalue ω_{\pm} (resp. ω'_{\pm}) of negative (resp. positive) real part.

- (b) *If $(\tau, \eta) \in \Xi_1$ with $\operatorname{Re}\tau > 0$, then the stable subspace $\mathcal{E}^-(t, x, \tau, \eta)$ of $\mathbb{A}(t, x, \tau, \eta)$ has dimension two and is spanned by*

$$\begin{cases} E_+(t, x, \tau, \eta) := \left(-(\tau + i\dot{v}_1^+ \eta)a_{12}^+, (\tau + i\dot{v}_1^+ \eta)(a_{11}^+ - \omega_+), 0, 0\right)^{\top}, \\ E_-(t, x, \tau, \eta) := \left(0, 0, (\tau + i\dot{v}_1^- \eta)(a_{22}^- - \omega_-), -(\tau + i\dot{v}_1^- \eta)a_{21}^-\right)^{\top}. \end{cases} \quad (5.39)$$

- (c) *Both ω_+ and ω_- admit a continuous extension to any point $(\tau, \eta) \in \Xi_1$ with $\operatorname{Re}\tau = 0$. If $(\tau, \eta) \in \Xi_1$ with $\tau = i\delta \in i\mathbb{R}$, then*

$$\begin{aligned} \tilde{\omega}_{\pm}(t, x, \tau, \eta) &:= \omega_{\pm}(t, x, \tau, \eta) - \frac{a_{11}^{\pm} + a_{22}^{\pm}}{2}(t, x, \tau, \eta) \\ &= \begin{cases} -\dot{C}_0^{\pm} \sqrt{\eta^2 - (\dot{C}_1^{\pm})^2 (\delta \pm \dot{C}_2^{\pm} \eta)^2} & \text{if } \eta^2 \geq (\dot{C}_1^{\pm})^2 (\delta \pm \dot{C}_2^{\pm} \eta)^2, \\ -i \operatorname{sgn}(\delta \pm \dot{C}_2^{\pm} \eta) \dot{C}_0^{\pm} \sqrt{(\dot{C}_1^{\pm})^2 (\delta \pm \dot{C}_2^{\pm} \eta)^2 - \eta^2} & \text{elsewise.} \end{cases} \end{aligned} \quad (5.40)$$

- (d) *Both $E_+(t, x, \tau, \eta)$ and $E_-(t, x, \tau, \eta)$ can be extended continuously to any point $(\tau, \eta) \in \Xi_1$ with $\operatorname{Re}\tau = 0$. These two vectors are linearly independent on the whole hemisphere Ξ_1 .*

- (e) *If $(t, x, \tau, \eta) \notin \Upsilon_{nd}$, where Υ_{nd} is given by*

$$\Upsilon_{nd} := \left\{ \tau \in \left\{ i(-\dot{C}_2^+ \pm (\dot{C}_1^+)^{-1})\eta, i(-\dot{C}_2^- \pm (\dot{C}_1^-)^{-1})\eta \right\} \right\}, \quad (5.41)$$

then matrix $\mathbb{A}(t, x, \tau, \eta)$ is diagonalizable.

Proof We just need to deduce that relations (5.38) hold and that $a_{11}^\pm + a_{22}^\pm$ are well-defined for any point $(\tau, \eta) \in \Xi_1$, since the other assertions can be proved similarly to the proof of Lemma 3.3.

By definition, we know that the eigenvalues of \mathbb{A}^\pm are roots ω of

$$\omega^2 - (a_{11}^\pm + a_{22}^\pm)\omega + a_{11}^\pm a_{22}^\pm - a_{12}^\pm a_{21}^\pm = 0, \quad (5.42)$$

from which we have

$$\tilde{\omega}_+^2 = \left(\frac{a_{11}^+ - a_{22}^+}{2} + a_{12}^+ \right) \left(\frac{a_{11}^+ - a_{22}^+}{2} - a_{12}^+ \right) + a_{12}^+(a_{12}^+ + a_{21}^+), \quad (5.43a)$$

$$\tilde{\omega}_-^2 = \left(\frac{a_{22}^- - a_{11}^-}{2} + a_{21}^- \right) \left(\frac{a_{22}^- - a_{11}^-}{2} - a_{21}^- \right) + a_{21}^-(a_{12}^- + a_{21}^-). \quad (5.43b)$$

We now deduce the expressions for $a_{11}^\pm - a_{22}^\pm - 2a_{12}^\pm$ and $a_{12}^\pm + a_{21}^\pm$. Recall that a_{ij}^\pm and b_{ij}^\pm are given by (5.37) and (5.22), respectively. Entries $A_{0,\pm}^{11}$ and $A_{1,\pm}^{11}$ are given by (5.23). For notational simplicity, we ignore indices “ \pm ” and “ $^\circ$ ” in the following expressions. We calculate coefficients A_j^\pm defined in (4.19) by using the computer algebra system “Maxima” to obtain the relations:

$$A_1^{23} = v_1 A_0^{23}, \quad A_1^{32} = v_1 A_0^{32}, \quad A_1^{22} - A_1^{33} = v_1 (A_0^{22} - A_0^{33}), \quad (5.44a)$$

$$A_1^{12} - A_1^{13} = v_1 (A_0^{12} - A_0^{13}), \quad A_1^{21} + A_1^{31} = v_1 (A_0^{21} + A_0^{31}), \quad (5.44b)$$

$$\begin{aligned} & A_1^{21}(A_0^{12} - A_0^{13}) - A_1^{13}(A_0^{21} + A_0^{31}) \\ &= v_1 \{ A_0^{21}(A_0^{12} - A_0^{13}) - A_0^{13}(A_0^{21} + A_0^{31}) \}, \end{aligned} \quad (5.44c)$$

$$\begin{aligned} & A_1^{13}(A_0^{21} + A_0^{31}) + A_1^{31}(A_0^{12} - A_0^{13}) \\ &= v_1 \{ A_0^{13}(A_0^{21} + A_0^{31}) + A_0^{31}(A_0^{12} - A_0^{13}) \}. \end{aligned} \quad (5.44d)$$

Then it follows from (5.44a) that

$$-b_{22} + b_{33} + 2b_{23} = (-A_0^{22} + A_0^{33} + 2A_0^{23})(\tau + iv_1\eta).$$

By virtue of (5.44b)–(5.44c), we obtain

$$\begin{aligned} b_{21}b_{12} - b_{31}b_{13} - 2b_{21}b_{13} &= b_{21}(b_{12} - b_{13}) - b_{13}(b_{21} + b_{31}) \\ &= \{ b_{21}(A_0^{12} - A_0^{13}) - b_{13}(A_0^{21} + A_0^{31}) \} (\tau + iv_1\eta) \\ &= \{ A_0^{21}(A_0^{12} - A_0^{13}) - A_0^{13}(A_0^{21} + A_0^{31}) \} (\tau + iv_1\eta)^2. \end{aligned}$$

Then

$$a_{11} - a_{22} - 2a_{12} = \mathbb{F}_2(\tau + iv_1\eta), \quad (5.45)$$

where

$$\begin{aligned} \mathbb{F}_2 &= (A_0^{11})^{-1} \{ A_0^{11}(-A_0^{22} + A_0^{33} + 2A_0^{23}) + A_0^{21}(A_0^{12} - A_0^{13}) - A_0^{13}(A_0^{21} + A_0^{31}) \} \\ &= \frac{2\partial_2 \Phi \{ \Gamma(\zeta^2 + \varrho^2)^{1/2}(\epsilon^2|v|^2 - 1) + \epsilon^2 c(\varsigma v_2 - \varrho v_1) \}}{c(\epsilon^2(\varrho v_2 + \varsigma v_1)^2 - \zeta^2 - \varrho^2)}. \end{aligned} \quad (5.46)$$

In particular, $\mathbb{F}_2^\pm = \pm 2\bar{\Gamma}/\bar{c} \neq 0$ when the perturbation $(\mathring{V}^\pm, \mathring{\Psi}^\pm)$ vanishes. As a consequence, \mathbb{F}_2^\pm never vanish by taking K in (4.2) small enough. Using (5.44a)–(5.44b) and (5.44d), we can deduce from a similar calculation that

$$a_{12} + a_{21} = \mathbb{F}_3(\tau + iv_1\eta), \quad (5.47)$$

where

$$\begin{aligned} \mathbb{F}_3 &= (A_0^{11})^{-1} \{A_0^{11}(-A_0^{23} - A_0^{32}) + A_0^{21}A_0^{13} + A_0^{31}A_0^{12}\} \\ &= -\frac{2\partial_2\Phi\epsilon^2(\zeta v_2 - \varrho v_1)}{\epsilon^2(\varrho v_2 + \zeta v_1)^2 - \zeta^2 - \varrho^2}. \end{aligned} \quad (5.48)$$

Relations (5.38) follows by plugging (5.45) and (5.47) into (5.43a)–(5.43b).

We now show that $a_{11}^\pm + a_{22}^\pm$ are well-defined. Use (5.44b) to derive

$$\begin{aligned} b_{21}b_{12} + b_{31}b_{13} &= b_{12}(b_{21} + b_{31}) + b_{31}(b_{13} - b_{12}) \\ &= \{b_{12}(A_0^{21} + A_0^{31}) + b_{31}(A_0^{13} - A_0^{12})\}(\tau + iv_1\eta), \end{aligned}$$

which implies

$$a_{11} + a_{22} = -b_{22} - b_{33} + \frac{b_{21}b_{12} + b_{31}b_{13}}{\mathbb{F}_1(\tau + iv_1\eta)} = \mathbb{F}_4\tau + i\eta\mathbb{F}_5, \quad (5.49)$$

where \mathbb{F}_4 and \mathbb{F}_5 are some smooth functions of $(\mathring{V}, \nabla\mathring{\Psi})$ that vanish when $(\mathring{V}, \mathring{\Psi}) = 0$ (cf. (3.27) with $a_{11}^\pm = \pm\mu_\pm$ and $a_{22}^\pm = \mp\mu_\pm$). The proof of the lemma can be completed by using the fact that \mathbf{A}_0^\pm and \mathbf{A}_1^\pm are smooth with respect to $(\mathring{V}^\pm, \nabla\mathring{\Psi}^\pm)$. \square

As in the constant coefficient case, we define the Lopatinskiĭ determinant associated with \mathbb{A} and β as

$$\Delta(t, x_1, \tau, \eta) := \det[\beta(t, x_1, \tau, \eta) (E_+(t, x_1, 0, \tau, \eta) \ E_-(t, x_1, 0, \tau, \eta))], \quad (5.50)$$

where β and E_\pm are given by (5.17) and (5.39), respectively. For the zeros of $\Delta(t, x_1, \tau, \eta)$, we have the following lemma.

Lemma 5.4 *Assume that $(\mathring{V}, \nabla\mathring{\Psi})$ is sufficiently small in $W^{2,\infty}(\Omega)$. Then*

$$\Delta(t, x_1, \tau, \eta) = 0 \quad \text{if and only if} \quad (t, x_1, \tau, \eta) \in \mathcal{Y}_c,$$

where $\mathcal{Y}_c := \mathcal{Y}_c^{-1} \cup \mathcal{Y}_c^0 \cup \mathcal{Y}_c^1$ is called the critical set with

$$\mathcal{Y}_c^q := \{(t, x_1, \tau, \eta) \in \mathbb{R}^2 \times \Xi : \tau = i\eta\dot{z}_q(t, x_1)\},$$

where \dot{z}_0 and $\dot{z}_{\pm 1}$ are real-valued functions of $(\mathring{V}^\pm|_{x_2=0}, \nabla\mathring{\varphi})$ satisfying

$$\mathcal{Y}_c \cap ((\mathcal{Y}_{nd} \cup \mathcal{Y}_p) \cap \{x_2 = 0\}) = \emptyset.$$

Moreover, each of these roots is simple in the sense that, if $q \in \{0, \pm 1\}$, then there exist a neighborhood \mathcal{V} of $(i\dot{z}_q\eta, \eta)$ in Ξ_1 and a C^∞ -function h_q defined on $\mathbb{R}^2 \times \mathcal{V}$ such that

$$\Delta(t, x_1, \tau, \eta) = (\tau - i\dot{z}_q\eta)h_q(t, x_1, \tau, \eta), \quad h_q(t, x_1, \tau, \eta) \neq 0 \quad (5.51)$$

for all $(\tau, \eta) \in \mathcal{V}$.

Proof Thanks to (5.17) and (5.39), we obtain that, for $(\tau, \eta) \in \Xi_1$,

$$\beta(t, x_1, \tau, \eta)(E_+(t, x_1, 0, \tau, \eta) - E_-(t, x_1, 0, \tau, \eta)) = \begin{pmatrix} \mathring{\zeta}_1 & \mathring{\zeta}_2 \\ \mathring{\zeta}_3 & \mathring{\zeta}_4 \end{pmatrix} \Big|_{x_2=0}, \quad (5.52)$$

where

$$\begin{aligned} \mathring{\zeta}_1 &:= \mathring{a}_+ \sqrt{\mathring{\varrho}_+^2 + \mathring{\varsigma}_+^2} (a_{11}^+ - a_{12}^+ - \omega_+), \quad \mathring{\zeta}_2 := -\mathring{a}_- \sqrt{\mathring{\varrho}_-^2 + \mathring{\varsigma}_-^2} (a_{22}^- - a_{21}^- - \omega_-), \\ \mathring{\zeta}_3 &:= \mathring{a}_+ \mathring{a}_- \left\{ -a_{12}^+ (m_1^+ + m_2^+) + (a_{11}^+ - \omega_+) (m_1^+ - m_2^+) \right\}, \\ \mathring{\zeta}_4 &:= \mathring{a}_+ \mathring{a}_- \left\{ -(a_{22}^- - \omega_-) (m_1^- + m_2^-) + a_{21}^- (m_1^- - m_2^-) \right\}. \end{aligned}$$

Recall that \mathring{a}_\pm and m_j^\pm are defined by (5.18) and (4.22), respectively.

Thanks to (5.45), we have

$$\begin{aligned} \mathring{\zeta}_1 &= \mathring{a}_+ \sqrt{\mathring{\varrho}_+^2 + \mathring{\varsigma}_+^2} \left(\frac{a_{11}^+ - a_{22}^+ - 2a_{12}^+}{2} - \tilde{\omega}_+ \right) \\ &= \mathring{a}_+ \sqrt{\mathring{\varrho}_+^2 + \mathring{\varsigma}_+^2} \left(\frac{\mathbb{F}_2^+ \mathring{a}_+}{2} - \tilde{\omega}_+ \right). \end{aligned} \quad (5.53)$$

Using (5.45) and (5.47) yields

$$\mathring{\zeta}_2 = \mathring{a}_- \sqrt{\mathring{\varrho}_-^2 + \mathring{\varsigma}_-^2} \left(\frac{\mathbb{F}_2^- + 2\mathbb{F}_3^-}{2} \mathring{a}_- + \tilde{\omega}_- \right). \quad (5.54)$$

It follows from (5.43a) that

$$\frac{a_{11}^+ - a_{22}^+ + 2a_{12}^+}{2} = \frac{2\tilde{\omega}_+^2 - 2a_{12}^+ (a_{12}^+ + a_{21}^+)}{a_{11}^+ - a_{22}^+ - 2a_{12}^+}.$$

By virtue of this last identity, we obtain

$$\begin{aligned} \mathring{\zeta}_3 &= \mathring{a}_+ \mathring{a}_- \left\{ \left(\tilde{\omega}_+ - \frac{a_{11}^+ - a_{22}^+ + 2a_{12}^+}{2} \right) (m_2^+ - m_1^+) - 2a_{12}^+ m_1^+ \right\} \\ &= \mathring{a}_+ \mathring{a}_- (m_2^+ - m_1^+) \tilde{\omega}_+ \left(1 - \frac{2\tilde{\omega}_+}{a_{11}^+ - a_{22}^+ - 2a_{12}^+} \right) \\ &\quad + \frac{2\mathring{a}_+ \mathring{a}_- a_{12}^+}{a_{11}^+ - a_{22}^+ - 2a_{12}^+} \left\{ (m_2^+ - m_1^+) (a_{12}^+ + a_{21}^+) - m_1^+ (a_{11}^+ - a_{22}^+ - 2a_{12}^+) \right\}. \end{aligned}$$

Use (5.45) and (5.47) to deduce

$$\begin{aligned} &(m_2^+ - m_1^+) (a_{12}^+ + a_{21}^+) - m_1^+ (a_{11}^+ - a_{22}^+ - 2a_{12}^+) \\ &= \mathring{a}_+ (m_2^+ \mathbb{F}_3^+ - m_1^+ (\mathbb{F}_2^+ + \mathbb{F}_3^+)) = 0, \end{aligned}$$

which, combined with (5.45), yields

$$\mathring{\zeta}_3 = \mathring{a}_+ \mathring{a}_- (m_2^+ - m_1^+) \tilde{\omega}_+ \left(1 - \frac{2\tilde{\omega}_+}{\mathring{a}_+ \mathbb{F}_2^+} \right). \quad (5.55)$$

Similar to the derivation of (5.55), we can infer from (5.43b), (5.45), and (5.47) that

$$\dot{\zeta}_4 = \dot{a}_+ \dot{a}_- (m_2^- + m_1^-) \tilde{\omega}_- \left(1 + \frac{2\tilde{\omega}_-}{\dot{a}_- (\mathbb{F}_2^- + 2\mathbb{F}_3^-)} \right). \quad (5.56)$$

Therefore, we find that $\Delta = \Delta_1 \Delta_2 \Delta_3|_{x_2=0}$, where

$$\begin{aligned} \Delta_1 &:= \frac{\mathbb{F}_2^+ \dot{a}_+}{2} - \tilde{\omega}_+, & \Delta_2 &:= \frac{\mathbb{F}_2^- + 2\mathbb{F}_3^-}{2} \dot{a}_- + \tilde{\omega}_-, \\ \Delta_3 &:= \det \begin{pmatrix} \dot{a}_+ \sqrt{\dot{\varrho}_+^2 + \dot{\varsigma}_+^2} & \dot{a}_- \sqrt{\dot{\varrho}_-^2 + \dot{\varsigma}_-^2} \\ \frac{2(m_2^+ - m_1^+)}{\mathbb{F}_2^+} \dot{a}_- \tilde{\omega}_+ & \frac{2(m_2^- + m_1^-)}{\mathbb{F}_2^- + 2\mathbb{F}_3^-} \dot{a}_+ \tilde{\omega}_- \end{pmatrix}. \end{aligned}$$

If $(\dot{V}^\pm, \dot{\Psi}^\pm) = 0$, then $\mathbb{F}_2^\pm = \pm 2\bar{\Gamma}/\bar{c}$, $\mathbb{F}_3^\pm = 0$, $m_1^\pm = 0$, and $m_2^\pm = 1/(\bar{\Gamma}\bar{c}\bar{h})$. Thus,

$$\begin{aligned} \Delta_1|_{(\dot{V}, \dot{\Psi})=0} &= \bar{c}^{-1} \bar{\Gamma} (\tau + i\bar{v}\eta) - \omega_+, & \Delta_2|_{(\dot{V}, \dot{\Psi})=0} &= -\bar{c}^{-1} \bar{\Gamma} (\tau - i\bar{v}\eta) + \omega_-, \\ \Delta_3|_{(\dot{V}, \dot{\Psi})=0} &= -\frac{1}{\bar{\Gamma}^2 \bar{h}} \{ \omega_- (\tau + i\bar{v}\eta)^2 + \omega_+ (\tau - i\bar{v}\eta)^2 \}. \end{aligned}$$

Recalling the proof of Lemma 3.4 and using the continuity of Δ_k with respect to $(\dot{V}, \nabla \dot{\Psi})$, we find that, if perturbation $(\dot{V}, \nabla \dot{\Psi})$ is suitably small in $W^{2,\infty}(\Omega)$, then Δ_1 and Δ_2 never vanish on $\mathbb{R}^2 \times \Xi$, and $\Delta_3(t, x_1, \tau, \eta) \neq 0$ for $\eta = 0$. Consequently, $\Delta(t, x_1, \tau, \eta) = 0$ if and only if $\Delta_3(t, x_1, \tau, \eta) = 0$ and $\eta \neq 0$.

Let $\eta \neq 0$. Setting $z := \tau/(i\eta)$, we obtain that $\Delta_3/(i\eta)^3 = \dot{Q}_1(z) + \dot{Q}_2(z)$, where

$$\begin{aligned} \dot{Q}_1(z) &:= \frac{2(m_2^- + m_1^-)}{\mathbb{F}_2^- + 2\mathbb{F}_3^-} \sqrt{\dot{\varrho}_+^2 + \dot{\varsigma}_+^2} \frac{\dot{a}_+ \tilde{\omega}_-}{(i\eta)^3}, \\ \dot{Q}_2(z) &:= \frac{2(m_1^+ - m_2^+)}{\mathbb{F}_2^+} \sqrt{\dot{\varrho}_-^2 + \dot{\varsigma}_-^2} \frac{\dot{a}_- \tilde{\omega}_+}{(i\eta)^3}. \end{aligned}$$

As in the proof of Lemma 3.4, we define

$$\dot{P}(z) := \bar{c}^2 \bar{h}^2 \bar{\Gamma}^2 (\dot{Q}_1(z)^2 - \dot{Q}_2(z)^2).$$

When $(\dot{V}, \dot{\Psi}) = 0$, $\dot{P}(z)$ is exactly a polynomial $P(z)$ of degree 5, given by (3.44). As a consequence, if K in (4.2) is suitably small, then $\dot{P}(z)$ is a polynomial function with degree 5 or 6, and there are functions \dot{z}_k , $k \in \{0, \pm 1, \pm 2, \pm 3\}$, of $(\dot{V}, \dot{\Psi})$ such that

$$\dot{P}(z) = (\dot{z}_3 z + 1) \dot{P}_1(z), \quad \dot{P}_1(z) = \dot{z}_{-3} \prod_{k \in \{0, \pm 1, \pm 2\}} (z - \dot{z}_k),$$

where \dot{z}_3 and $\dot{P}_1(z)$ satisfy that

$$\dot{z}_3 = 0, \quad \dot{P}_1(z) = P(z) \quad \text{when } (\dot{V}, \dot{\Psi}) = 0.$$

Under condition (2.25), we can compute that the discriminants of $\frac{d^j P(z)}{dz^j}$, $j \in \{0, 1, 2, 3\}$, are all positive. Since the discriminant for a polynomial is continuous with respect to the coefficients of the polynomial, we take K suitably small to conclude that the discriminants of $\frac{d^j \hat{P}_1(z)}{dz^j}$, $j \in \{0, 1, 2, 3\}$, are all positive. Consequently, roots \hat{z}_k , $k \in \{0, \pm 1, \pm 2\}$, of $\hat{P}_1(z)$ are real and distinct.

Noting that the coefficients of $\hat{P}(z)$ are all real, we obtain that $\hat{z}_{\pm 3}$ are both real. Since \hat{z}_k , $k \in \{0, \pm 1, \pm 2\}$, are all different, we infer that \hat{z}_k , $k \in \{0, \pm 1, \pm 2, \pm 3\}$, can be expressed as continuous functions of the coefficients of $\hat{P}(z)$. Choosing K in (4.2) sufficiently small, we see that \hat{z}_{-3} is always nonzero, \hat{z}_3 and \hat{z}_0 are in a small neighborhood of 0, and \hat{z}_k , $k \in \{\pm 1, \pm 2\}$, are respectively in a small neighborhood of $\pm z_{|k|}$ with z_1 and z_2 given by (3.48).

We then use (5.40) and employ an entirely similar argument as in the proof of Lemma 3.4 to conclude the result as expected. \square

In view of Lemma 5.4, we can obtain the following result by using the continuity of A_k^\pm and the fact that the perturbation, $(\hat{V}, \hat{\Psi})$, has a compact support (see [15, Page 423] for the proof of Proposition 5.3 (c)).

Proposition 5.3 *Assume that $(\hat{V}, \hat{\Psi})$ satisfies (4.1)–(4.2) with K being sufficiently small. Then we can find neighborhoods \mathcal{V}_c^q of \mathcal{T}_c^q , $q \in \{0, \pm 1\}$, in $\bar{\Omega} \times \Xi$ such that*

- (a) $\mathcal{V}_c^q \cap (\mathcal{T}_p \cup \mathcal{T}_{nd}) = \emptyset$.
- (b) Matrix \mathbb{A} defined by (5.37) is diagonalizable on \mathcal{V}_c^q . In particular, there exist matrices $Q_0^\pm \in \Gamma_2^0$ such that

$$Q_0^\pm(z) \mathbb{A}^\pm(z) Q_0^\pm(z)^{-1} = \text{diag}(\omega_\pm(z), \omega'_\pm(z)) =: \mathbb{D}_1^\pm \quad (5.57)$$

for all $z = (t, x, \tau, \eta) \in \mathcal{V}_c^q$, where $\omega_+(z) \neq \omega'_+(z)$ and $\omega_-(z) \neq \omega'_-(z)$.

- (c) Let \hbar be $\text{Im} \omega_+$ or $\text{Im} \omega_-$. Then the solution of the system:

$$\begin{cases} \frac{dt}{dx_2} = \frac{\partial \hbar}{\partial \delta}(t, x_1, x_2, \tau, \eta), & \frac{dx_1}{dx_2} = \frac{\partial \hbar}{\partial \eta}(t, x_1, x_2, \tau, \eta), \\ \frac{d\delta}{dx_2} = -\frac{\partial \hbar}{\partial t}(t, x_1, x_2, \tau, \eta), & \frac{d\eta}{dx_2} = -\frac{\partial \hbar}{\partial x_1}(t, x_1, x_2, \tau, \eta), \\ (t, x_1, \gamma + i\delta, \eta)|_{x_2=0} \in \mathcal{V}_c^q \cap \{x_2 = 0\} \end{cases} \quad (5.58)$$

defines a curve $(t, x_1, \gamma + i\delta, \eta)$ for all $x_2 \geq 0$, which remains in \mathcal{V}_c^q and is called the bicharacteristic curve.

In order to absorb the error terms caused by microlocalization, as in [15, 17], we will construct the weight functions that vanish on the bicharacteristic curves originating from \mathcal{T}_c and that are nonzero far from these curves.

We define the complex-valued functions: For all $z = (t, x_1, \tau, \eta) \in \mathbb{R}^2 \times \Xi_1$ with $\tau = \gamma + i\delta$,

$$\sigma_q(z) := -i\gamma + \tilde{\sigma}_q(z), \quad \tilde{\sigma}_q(z) := \delta - \eta \hat{z}_q(t, x_1), \quad q \in \{0, \pm 1\}, \quad (5.59)$$

and extend σ_q to $\mathbb{R}^2 \times \Xi$ as a homogeneous mapping of degree 1 with respect to (τ, η) . Functions $\hat{z}_0(t, x_1)$ and $\hat{z}_{\pm 1}(t, x_1)$ are given by Lemma 5.4 and correspond to the points where the Lopatinskiĭ determinant vanishes. Symbol σ_q thus belongs to $\Gamma_{\frac{1}{2}}^1$ such that

$$\mathcal{Y}_c^q = \{z = (t, x_1, \tau, \eta) \in \mathbb{R}^2 \times \Xi : \sigma_q(z) = 0\}.$$

In view of Proposition 5.3(c), we can construct solutions σ_{\pm}^q of the linear transport equations:

$$\begin{cases} \partial_2 \sigma_+^q + \{\sigma_+^q, \text{Im } \omega_+\} = 0 & \text{if } x_2 > 0, \\ \partial_2 \sigma_-^q + \{\sigma_-^q, \text{Im } \omega_-\} = 0 & \text{if } x_2 > 0, \\ \sigma_+^q = \sigma_-^q = \sigma_q & \text{if } x_2 = 0, \end{cases} \quad (5.60)$$

where σ_q is given in (5.59), and $\{\cdot, \cdot\}$ is the Poisson bracket defined by (5.28). Then we infer that σ_+^q (resp. σ_-^q) is constant along each bicharacteristic curve defined by (5.58) with $\bar{h} = \text{Im } \omega_+$ (resp. $\bar{h} = \text{Im } \omega_-$). In particular, function σ_+^q (resp. σ_-^q) vanishes *only* on the bicharacteristic curves originated from \mathcal{Y}_c^q with $\bar{h} = \text{Im } \omega_+$ (resp. $\bar{h} = \text{Im } \omega_-$). By shrinking \mathcal{V}_c^q if necessary, we may assume that σ_{\pm}^q are defined in the whole set \mathcal{V}_c^q . We will see that functions σ_{\pm}^q are appropriate to deal with the error terms appearing in the energy estimates.

From the above analysis, the whole space $\bar{\Omega} \times \Xi$ is naturally divided into three disjoint subsets: \mathcal{Y}_p , \mathcal{V}_c , and $\bar{\Omega} \times \Xi \setminus (\mathcal{Y}_p \cup \mathcal{V}_c)$, where $\mathcal{V}_c := \mathcal{V}_c^{-1} \cup \mathcal{V}_c^0 \cup \mathcal{V}_c^1$. To derive our energy estimate (5.21), we introduce smooth cut-off functions according to this division. More precisely, we introduce nonnegative functions χ_p^{\pm} and χ_c^q (with values in $[0, 1]$), $q \in \{0, \pm 1\}$, such that

- χ_p^{\pm} and χ_c^q are C^∞ and homogeneous of degree 0 with respect to (τ, η) so that they belong to Γ_k^0 for all integer k ;
- $\text{supp } \chi_c^q \subset \mathcal{V}_c^q$ and $\chi_c^q \equiv 1$ in a neighborhood of the bicharacteristic curves originated from the critical set \mathcal{Y}_c^q ;
- $\chi_p^{\pm} \equiv 1$ in a neighborhood of \mathcal{Y}_p^{\pm} , $\text{supp } \chi_p^+ \cap \text{supp } \chi_p^- = \emptyset$, and $\text{supp } \chi_p^{\pm} \cap \text{supp } \chi_c^q = \emptyset$ for all $q \in \{0, \pm 1\}$.

Since σ_+^q and σ_-^q vanish *only* on the bicharacteristic curves originated from \mathcal{Y}_c^q , there exists a constant c such that

$$|\sigma_{\pm}^q| \geq c > 0 \quad \text{in } \{\chi_c^q < 1\} \cap \mathcal{V}_c^q. \quad (5.61)$$

We also define

$$\chi_p := \chi_p^+ + \chi_p^-, \quad \chi_c := \chi_c^{-1} + \chi_c^0 + \chi_c^1, \quad \chi_u := 1 - \chi_p - \chi_c. \quad (5.62)$$

Then χ_u has support far from the poles and the bicharacteristic curves originated from \mathcal{Y}_c . We observe that the Lopatinskiĭ determinant does not vanish on $\text{supp } \chi_u \cap \{x_2 = 0\}$. This enables us to apply the standard Kreiss' symmetrizers to derive the energy estimate for $T_{\chi_u}^\gamma W^{\text{nc}}$, which will be shown in §5.4. After that, we will show how the traces of $T_{\chi_p}^\gamma W^{\text{nc}}$ and $T_{\chi_c}^\gamma W^{\text{nc}}$ can be estimated. At the end of this section, we will complete the proof of Theorem 5.1 by using a weighted energy estimate with the weight functions σ_{\pm}^q given by (5.60). In particular, we will prove that the microlocalization error terms can be absorbed by such a weighted estimate.

5.4 Estimate at Good Frequencies

In this subsection, we show how the solutions of problem (5.20) can be estimated for the frequencies that are far from both the poles \mathcal{Y}_p and the critical set \mathcal{Y}_c . We define

$$W_u^{\text{nc}} := (T_{\chi_u}^\gamma W_2^+, T_{\chi_u}^\gamma W_3^+, T_{\chi_u}^\gamma W_2^-, T_{\chi_u}^\gamma W_3^-)^\top$$

and introduce a smooth cut-off function $\tilde{\chi}_u$ with values in $[0, 1]$ such that

$$\tilde{\chi}_u \equiv 1 \text{ on } \text{supp } \chi_u, \quad \text{supp } \tilde{\chi}_u \cap \mathcal{Y}_p = \emptyset, \quad (\text{supp } \tilde{\chi}_u \cap \{x_2 = 0\}) \cap \mathcal{Y}_c = \emptyset,$$

where χ_u is given by (5.62). Employing the same analysis as in § 5.2, we derive that W_u^{nc} satisfies

$$\begin{cases} \partial_2 W_u^{\text{nc}} = T_{\mathbb{A}_u}^\gamma W_u^{\text{nc}} + T_{\mathbb{E}}^\gamma W_u^{\text{nc}} + T_r^\gamma W + \mathcal{R}_0 T_{\chi_u}^\gamma F + \mathcal{R}_{-1} W, & x_2 > 0, \\ T_\beta^\gamma W_u^{\text{nc}} = G + \mathcal{R}_{-1} W^{\text{nc}}, & x_2 = 0, \end{cases} \quad (5.63)$$

where $\beta \in \Gamma_2^0$ is given by (5.17) for $(\tau, \eta) \in \Xi_1$. The symbol matrices \mathbb{A}_u is defined as \mathbb{A}_r in (5.36) with $\tilde{\chi}_\pm$ replaced by $\tilde{\chi}_u$. Both \mathbb{E} and r have the same block diagonal structure as \mathbb{A}_u and belong to Γ_1^0 . We note that $\mathbb{A}_u \equiv \mathbb{A}$ on region $\{\tilde{\chi}_u \equiv 1\}$, and r is identically zero on region $\{\chi_u \equiv 1\} \cup \{\chi_u \equiv 0\}$.

In view of Lemma 5.4, we find that the Lopatinskiĭ determinant never vanishes on $\text{supp } \chi_u \cap \{x_2 = 0\}$. Note that the perturbation, $(\check{V}, \check{\Psi})$, is assumed in (4.1) to have a compact support. In the following lemma, we construct the Kreiss' symmetrizers that are microlocalized at all frequencies in the compact set \mathbb{K} , where

$$\mathbb{K} := \text{supp } \chi_u \cap \{-T \leq t \leq 2T, x_2 \geq 0, |x| \leq R, (\tau, \eta) \in \Xi_1\}.$$

Lemma 5.5 *Assume that (4.1)–(4.2) hold for a sufficiently small positive constant K . Then, for each $z_0 \in \mathbb{K}$, there exist a neighborhood \mathcal{V} of z_0 in \mathbb{K} and C^∞ -mappings $r(z)$ and $T(z)$ defined on \mathcal{V} such that*

- (a) *Matrix $r(z)$ is Hermitian, and $T(z)$ is invertible for all $z \in \mathcal{V}$;*
- (b) *There exists $c > 0$ so that*

$$\text{Re}(r(z)T(z)\mathbb{A}(z)T(z)^{-1}) \geq c\gamma I \quad \text{for all } z \in \mathcal{V} \text{ with } \gamma = \text{Re } \tau; \quad (5.64)$$

- (c) *If $z_0 \in \mathbb{K} \cap \{x_2 = 0\}$, then there exists a positive constant C so that*

$$r(z) + C(\beta(z)T(z)^{-1})^* \beta(z)T(z)^{-1} \geq I \quad (5.65)$$

for all $z \in \mathcal{V} \cap \{x_2 = 0\}$,

where \mathbb{A} and β are given by (5.37) and (5.17), respectively.

To prove Lemma 5.5, we first establish the following result.

Lemma 5.6 *Let $z_0 = (t_0, x_0, \tau_0, \eta_0) \in \bar{\Omega} \times \Xi_1$ so that $\operatorname{Re} \tau_0 = 0$ and $\tau_0 \neq i(-\dot{C}_2^+ \pm (\dot{C}_1^+)^{-1})\eta_0$. Assume that K given in (4.2) is sufficiently small. Then there exists a neighborhood \mathcal{V} of z_0 in $\bar{\Omega} \times \Xi_1$ such that*

$$\operatorname{Re} \omega_+(z) \lesssim -\gamma, \quad \operatorname{Re} \omega'_+(z) \gtrsim \gamma \quad (5.66)$$

for all $z = (t, x, \tau, \eta) \in \mathcal{V}$ with $\gamma = \operatorname{Re} \tau$. A similar result holds for ω_- and ω'_- near $z_0 = (t_0, x_0, \tau_0, \eta_0) \in \bar{\Omega} \times \Xi_1$ so that $\operatorname{Re} \tau_0 = 0$ and $\tau_0 \neq i(-\dot{C}_2^- \pm (\dot{C}_1^-)^{-1})\eta_0$.

Proof This proof is divided into two steps.

1. If $\operatorname{Re} \omega_+(z_0) < 0$, then one uses the identities:

$$\tilde{\omega}_+ = \omega_+ - \frac{a_{11}^+ + a_{22}^+}{2}, \quad -\tilde{\omega}_+ = \omega'_+ - \frac{a_{11}^+ + a_{22}^+}{2}, \quad (5.67)$$

and (5.49) to infer that $\operatorname{Re} \omega'_+(z_0) > 0$ for sufficiently small K . Since $\gamma = \operatorname{Re} \tau \leq 1$ for $(\tau, \eta) \in \Xi_1$, estimates (5.66) follow directly from the continuity of $\omega_+(z)$ and $\omega'_+(z)$ with respect to z .

2. Assume that $\operatorname{Re} \omega_+(z_0) = 0$. It follows from $\operatorname{Re} \tau_0 = 0$ and (5.49) that $\operatorname{Re} \tilde{\omega}_+(z_0) = 0$. Thanks to (5.40), $\eta_0^2 < (\dot{C}_1^+)^2(\delta_0 + \dot{C}_2^+\eta_0)^2$ for $\delta_0 = \operatorname{Im} \tau_0$ so that $\delta_0 \neq -\dot{C}_2^+\eta_0$. For all $(i\delta, \eta)$ with $(\dot{C}_1^+)^2(\delta + \dot{C}_2^+\eta)^2 > \eta^2$, we apply (5.40) again to derive that

$$\delta \neq -\dot{C}_2^+\eta, \quad \tilde{\omega}_+(t, x, i\delta, \eta) = -i \operatorname{sgn}(\delta + \dot{C}_2^+\eta) \dot{C}_0^+ \sqrt{(\dot{C}_1^+)^2(\delta + \dot{C}_2^+\eta)^2 - \eta^2}. \quad (5.68)$$

Since $\tau_0 \neq i(-\dot{C}_2^+ \pm (\dot{C}_1^+)^{-1})\eta_0$, $\tilde{\omega}_+$ depends analytically on (τ, η) by applying the implicit functions theorem to (5.38). In particular, we obtain that, for z near z_0 ,

$$\tilde{\omega}_+(t, x, \tau, \eta) \partial_\gamma \tilde{\omega}_+(t, x, \tau, \eta) = (\dot{C}_0^+ \dot{C}_1^+)^2 (\tau + i\dot{C}_2^+\eta). \quad (5.69)$$

From (5.68)–(5.69), $\partial_\gamma \tilde{\omega}_+(t, x, i\delta, \eta)$ is real and negative for $(t, x, i\delta, \eta)$ in a suitable neighborhood \mathcal{V} of z_0 . Using the Taylor expansion yields that, for all $(\tau, \eta) \in \mathcal{V}$,

$$\tilde{\omega}_+(\tau, \eta) = \tilde{\omega}_+(i\delta, \eta) + \partial_\gamma \tilde{\omega}_+(i\delta, \eta) \gamma + O(\gamma^2) \quad (\gamma \rightarrow 0).$$

Then we deduce that $\operatorname{Re} \tilde{\omega}_+ \lesssim -\gamma$, up to shrinking \mathcal{V} . In view of (5.49) and (5.67), estimates (5.66) follow by taking K small enough. \square

Proof of Lemma 5.5 The proof is divided into two cases.

Case 1. Let $z_0 \in \mathbb{K} \setminus \mathcal{Y}_{\text{nd}}$ with \mathcal{Y}_{nd} given in (5.41). In light of Lemmas 5.4 and 5.6, we can find a neighborhood \mathcal{V} of z_0 in \mathbb{K} such that

$$\operatorname{Re} \omega_\pm(z) \lesssim -\gamma, \quad \operatorname{Re} \omega'_\pm(z) \gtrsim \gamma \quad \text{for all } z \in \mathcal{V}, \quad (5.70)$$

and

$$\Delta(z) \neq 0 \quad \text{for all } z \in \mathcal{V} \cap \{x_2 = 0\}. \quad (5.71)$$

According to Lemma 5.3, matrix \mathbb{A} is diagonalizable in \mathcal{V} . Indeed, a smooth basis of the eigenvectors is given by

$$E_+(z), \quad Y_+(z) := ((\tau + i\dot{v}_1^+ \eta)(-a_{22}^+ + \omega'_+), (\tau + i\dot{v}_1^+ \eta)a_{21}^+, 0, 0)^\top, \quad (5.72)$$

$$E_-(z), \quad Y_-(z) := (0, 0, (\tau + i\dot{v}_1^- \eta)a_{12}^-, (\tau + i\dot{v}_1^- \eta)(-a_{11}^- + \omega'_-))^\top, \quad (5.73)$$

where E_\pm are given by (5.39). Notice that E_\pm and Y_\pm are linearly independent in \mathcal{V} . We can thus define the smooth and invertible matrix $T(z) := (E_+ \ E_- \ Y_+ \ Y_-)^{-1}$ in \mathcal{V} so that

$$T(z)\mathbb{A}(z)T(z)^{-1} = \text{diag}(\omega_+, \omega_-, \omega'_+, \omega'_-) \quad \text{for all } z \in \mathcal{V}.$$

Construct the symmetrizer, $r(z)$, as

$$r(z) := \text{diag}(-1, -1, K', K') \quad \text{for all } z \in \mathcal{V}, \quad (5.74)$$

where $K' \geq 1$ is a constant to be chosen. Then we can obtain (5.64) from (5.70) directly.

Thanks to (5.71), we have

$$|\beta(z)(E_+(z) \ E_-(z))Z^-| \gtrsim |Z^-| \quad \text{for all } z \in \mathcal{V} \cap \{x_2 = 0\}, \quad Z^- \in \mathbb{C}^2. \quad (5.75)$$

This implies

$$|Z^-|^2 \leq C_0(|Z^+|^2 + |\beta(z)T(z)^{-1}Z|^2)$$

for all $Z = (Z^-, Z^+)^\top \in \mathbb{C}^4$ with $Z^\pm \in \mathbb{C}^2$, where C_0 is some positive constant independent of $z \in \mathcal{V}$. Then we have

$$\begin{aligned} & \langle (r(z) + 2C_0(\beta(z)T(z)^{-1})^* \beta(z)T(z)^{-1})Z, Z \rangle_{\mathbb{C}^4} \\ &= -|Z^-|^2 + K|Z^+|^2 + 2C_0|\beta(z)T(z)^{-1}Z|^2 \\ &\geq |Z^-|^2 + (K' - 2C_0)|Z^+|^2 \geq |Z|^2 \end{aligned}$$

by choosing $K' \geq 2C_0 + 1$, which implies (5.65).

Case 2. Let $z_0 \in \mathbb{K} \cap \mathcal{Y}_{\text{nd}}$. Then symbol \mathbb{A} is not diagonalizable at z_0 . We consider without loss of generality that $z_0 = (t_0, x_0, \tau_0, \eta_0) \in \mathbb{K}$ satisfies $\tau_0 = -i(\dot{C}_2^+ \pm (\dot{C}_1^+)^{-1})\eta_0$. The case, $\tau_0 = -i(\dot{C}_2^- \pm (\dot{C}_1^-)^{-1})\eta_0$, can be dealt with in an entirely similar way. Using (3.36) and the continuity of \dot{C}_j^\pm in $(\dot{V}^\pm, \nabla \dot{\psi}^\pm)$, we take K sufficiently small to find

$$\tau_0 + i\dot{v}_1^+ \eta_0 \neq 0, \quad \tau_0 \neq -i(\dot{C}_2^- \pm (\dot{C}_1^-)^{-1})\eta_0,$$

which, combined with (5.38) and (5.45), imply

$$(a_{11}^+ - a_{22}^+ - 2a_{12}^+)(z_0) = \mathbb{F}_2^+(z_0)a_+(z_0) \neq 0, \quad \omega_-(z_0) \neq \omega'_-(z_0).$$

Since $\tilde{\omega}_+(z_0) = 0$, we use (5.43a), (5.45), and (5.47) to obtain

$$\frac{(a_{11}^+ - a_{22}^+)(z_0)}{2} = -\left(1 + \frac{2\mathbb{F}_3^+(z_0)}{\mathbb{F}_2^+(z_0)}\right)a_{12}^+(z_0), \quad (5.76)$$

and hence

$$(a_{11}^+ - \omega_+)(z_0) = \frac{a_{11}^+ - a_{22}^+}{2}(z_0) - \tilde{\omega}_+(z_0) = -\left(1 + \frac{2\mathbb{F}_3^+(z_0)}{\mathbb{F}_2^+(z_0)}\right)a_{12}^+(z_0) \neq 0.$$

In view of (5.39) and (5.72)–(5.73), we find that the vectors:

$$\tilde{E}_+ := (-a_{12}^+, -(\mathbb{F}_2^+ + 2\mathbb{F}_3^+)a_{12}^+/\mathbb{F}_2^+, 0, 0)^\top, \quad \tilde{Y}_+ := (i, 0, 0, 0)^\top, \quad E_-(z), \quad Y_-(z)$$

form a smooth basis in neighborhood \mathcal{V} . Define $T(z) := (\tilde{E}_+ \quad \tilde{Y}_+ \quad E_- \quad Y_-)^{-1}$. Then

$$T(z)\mathbb{A}(z)T(z)^{-1} = \begin{pmatrix} a_r & 0 & 0 \\ 0 & \omega_- & 0 \\ 0 & 0 & \omega'_- \end{pmatrix} \quad \text{for all } z \in \mathcal{V},$$

where a_r is the 2×2 matrix with (i, j) -entry $a_r^{ij}(z)$:

$$\begin{aligned} a_r^{11}(z) &= \frac{(2\mathbb{F}_3^+ + \mathbb{F}_2^+)a_{22}^+ + \mathbb{F}_2^+a_{21}^+}{2\mathbb{F}_3^+ + \mathbb{F}_2^+}, & a_r^{12}(z) &= \frac{-\mathbb{F}_2^+a_{21}^+i}{(2\mathbb{F}_3^+ + \mathbb{F}_2^+)a_{12}^+}, \\ a_r^{22}(z) &= \frac{(2\mathbb{F}_3^+ + \mathbb{F}_2^+)a_{11}^+ - \mathbb{F}_2^+a_{21}^+}{2\mathbb{F}_3^+ + \mathbb{F}_2^+}, & a_r^{21}(z) &= \frac{a_{12}^+\tilde{a}_r^{21}(z)}{(2\mathbb{F}_3^+ + \mathbb{F}_2^+)\mathbb{F}_2^+i} \end{aligned}$$

with

$$\tilde{a}_r^{21}(z) := \mathbb{F}_2^+(2\mathbb{F}_3^+ + \mathbb{F}_2^+)(a_{22}^+ - a_{11}^+ - 2a_{12}^+) - 4(\mathbb{F}_3^+)^2a_{12}^+ + (\mathbb{F}_2^+)^2(a_{12}^+ + a_{21}^+).$$

By virtue of (5.45), (5.47), and (5.76), we derive

$$\begin{aligned} 2\mathbb{F}_2^+a_{21}^+(z_0) &= -2\mathbb{F}_2^+a_{12}^+(z_0) + 2\mathbb{F}_3^+(a_{11}^+ - a_{22}^+ - 2a_{12}^+)(z_0) \\ &= (2\mathbb{F}_3^+ + \mathbb{F}_2^+)(a_{11}^+ - a_{22}^+)(z_0), \\ \tilde{a}_r^{21}(z_0) &= 4\mathbb{F}_3^+(\mathbb{F}_3^+ + \mathbb{F}_2^+)a_{12}^+(z_0) + (\mathbb{F}_2^+)^2(a_{12}^+ + a_{21}^+)(z_0) \\ &= 4\mathbb{F}_3^+(\mathbb{F}_3^+ + \mathbb{F}_2^+)a_{12}^+(z_0) + \mathbb{F}_2^+\mathbb{F}_3^+(a_{11}^+ - a_{22}^+ - 2a_{12}^+)(z_0) \\ &= 0, \end{aligned}$$

which implies

$$a_r^{11}(z_0) - a_r^{22}(z_0) = a_r^{21}(z_0) = 0. \quad (5.77)$$

We now look for a symmetrizer r with the form:

$$r(z) = \begin{pmatrix} s(z) & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & K' \end{pmatrix},$$

where $K' \geq 1$ is some real constant, and s is some 2×2 Hermitian matrix, depending smoothly on z . Both K' and s are to be fixed such that (5.65) holds for $z \in \mathcal{V} \cap \{x_2 = 0\}$ when $z_0 \in \{x_2 = 0\}$ and (5.64) holds for all $z \in \mathcal{V}$.

We recall that $\Delta(z) \neq 0$ in $\mathcal{V} \cap \{x_2 = 0\}$ so that (5.75) holds. Noting that the first and third columns of $T(z_0)^{-1}$ are $E_+(z_0)$ and $E_-(z_0)$, we can find a positive constant C_0 such that, if $z_0 \in \mathbb{K} \cap \{x_2 = 0\}$, then

$$|Z_1|^2 + |Z_3|^2 \leq C_0(|Z_2|^2 + |Z_4|^2 + |\beta(z_0)T(z_0)^{-1}Z|^2)$$

for all $Z = (Z_1, Z_2, Z_3, Z_4)^T \in \mathbb{C}^4$. Assume that the Hermitian matrix s satisfies

$$s(z_0) = \begin{pmatrix} 0 & e_1 \\ e_1 & e_2 \end{pmatrix} =: E, \quad (5.78)$$

where e_1 and e_2 are some real constants to be fixed. Then we have

$$\begin{aligned} & \langle (r(z_0) + C'C_0(\beta(z_0)T(z_0)^{-1})^* \beta(z_0)T(z_0)^{-1})Z, Z \rangle_{\mathbb{C}^4} \\ &= 2e_1 \operatorname{Re} \langle Z_1, Z_2 \rangle_{\mathbb{C}} + e_2 |Z_2|^2 - |Z_3|^2 + K'|Z_4|^2 + C'C_0 |\beta(z_0)T(z_0)^{-1}Z|^2 \\ &\geq (C' - \max\{|e_1|, 1\}) (|Z_1|^2 + |Z_3|^2) \\ &\quad + (e_2 - |e_1| - C'C_0)|Z_2|^2 + (K' - C'C_0)|Z_4|^2. \end{aligned}$$

We choose $C' = \max\{|e_1|, 1\} + 2$, $e_2 = |e_1| + C'C_0 + 2$, and $K' = C'C_0 + 2$ to obtain

$$r(z_0) + C'C_0(\beta(z_0)T(z_0)^{-1})^* \beta(z_0)T(z_0)^{-1} \geq 2I.$$

Using the continuity and shrinking \mathcal{V} if necessary, we derive estimate (5.65) for $C = C'C_0$.

It remains to choose a suitable Hermitian matrix $s(\tau, \eta)$ and $e_1 \in \mathbb{R}$ such that both (5.64) and (5.78) hold. Since $\tau_0 \neq -i(C_2^- \pm (C_1^-)^{-1})\eta_0$, we find that $\operatorname{Re} \omega_-(z) \lesssim -\gamma$ and $\operatorname{Re} \omega'_-(z) \gtrsim \gamma$ for all $z \in \mathcal{V}$ from Lemma 5.6. Consequently, it suffices to find $e_1 \in \mathbb{R}$ and a Hermitian matrix $s(z)$ satisfying (5.78) and

$$\operatorname{Re}(s(z)a_r(z)) \gtrsim \gamma I \quad \text{for all } z \in \mathcal{V}. \quad (5.79)$$

To this end, we let

$$s(z) = E + F(z) + \gamma G(z)$$

for some smooth 2×2 Hermitian matrices F and G satisfying $F(z_0) = 0$, where E is defined by (5.78). In light of Taylor's formula, we may write

$$a_r(z) = a_r(t, x, \gamma + i\delta, \eta) = a_r(t, x, i\delta, \eta) + \gamma \partial_\gamma a_r(t, x, i\delta, \eta) + \gamma^2 N_1(z)$$

for a suitable continuous function N_1 . Noting from (5.22) and (5.37) that $a_r^{ij}(t, x, i\delta, \eta)$ are purely imaginary, we may choose

$$F(z) := \operatorname{diag}(f(z), 0)$$

with

$$f(z) = \frac{e_1(a_r^{11} - a_r^{22})(t, x, i\delta, \eta) + e_2 a_r^{21}(t, x, i\delta, \eta)}{a_r^{12}(t, x, i\delta, \eta)},$$

so that matrix $(E + F(z))a_r(t, x, i\delta, \eta)$ is symmetric and purely imaginary for all $z \in \mathcal{V}$. It follows from (5.77) that $F(z_0) = 0$. Therefore, we have

$$\begin{aligned} & \operatorname{Re}(s(z)a_r(z)) \\ &= \operatorname{Re}\{\gamma(E + F(z))\partial_\gamma a_r(t, x, i\delta, \eta) + \gamma G(z)a_r(t, x, i\delta, \eta) + \gamma^2 N_2(z)\} \\ &= \gamma \operatorname{Re}\{E\partial_\gamma a_r(t, x, i\delta, \eta) + G(z)a_r(t, x, i\delta, \eta) + N_3(z)\} \end{aligned}$$

for some continuous functions N_2 and N_3 satisfying $N_3(z_0) = 0$, where we have used $F(z_0) = 0$. According to (3.28), we see that, for $\tilde{\tau} = -i(\bar{C}_2 \pm \bar{C}_1^{-1})\eta_0$,

$$\begin{aligned} & \partial_\gamma a_r^{21}(t_0, x_0, \tilde{\tau}, \eta_0)|_{(\check{V}, \check{\Psi})=0} = \partial_\gamma \{2i m_+(m_+ - \mu_+)\}(\tilde{\tau}, \eta_0) \\ &= \frac{2i\bar{I}^2(i\eta_0 + \epsilon^2\bar{v}\tilde{\tau})}{(\tilde{\tau} + i\bar{v}\eta_0)^3} \\ & \quad \times \left\{ 2\epsilon^2\bar{c}^2\bar{v}(\tilde{\tau} + i\bar{v}\eta_0)(i\eta_0 + \epsilon^2\bar{v}\tilde{\tau})^2 - \bar{c}^2(i\eta_0 + \epsilon^2\bar{v}\tilde{\tau})^3 - \epsilon^2\bar{v}(\tilde{\tau} + i\bar{v}\eta_0)^3 \right\} \\ &= 2\bar{I}^2(1 - \epsilon^2\bar{v}^2) \frac{\eta_0^2(i\eta_0 + \epsilon^2\bar{v}\tilde{\tau})}{\tilde{\tau} + i\bar{v}\eta_0} \in \mathbb{R} \setminus \{0\}, \end{aligned}$$

where we have used $\bar{c}^2(i\eta_0 + \epsilon^2\bar{v}\tilde{\tau})^2 = (\tilde{\tau} + i\bar{v}\eta_0)^2$ and condition (2.10). Then $\partial_\gamma a_r^{21}(z_0)$ is always non-zero by choosing K sufficiently small and using the continuity of \mathbf{A}_j^\pm and \check{C}_j^\pm . In order to obtain (5.64), we choose

$$e_1 := (\partial_\gamma a_r^{21}(z_0))^{-1}, \quad G(z) := \begin{pmatrix} 0 & ig \\ -ig & 0 \end{pmatrix}$$

for some positive constant g . This choice of e_1 and G yields

$$\operatorname{Re}\{E\partial_\gamma a_r(z_0) + G(z)a_r(z_0)\} = \begin{pmatrix} 1 & \star \\ \star & \star \end{pmatrix} + \begin{pmatrix} 0 & iga_r^{22}(z_0) \\ -iga_r^{11}(z_0) & -iga_r^{12}(z_0) \end{pmatrix},$$

where the entries with \star are the coefficients that depend only on z_0 , e_1 , and e_2 (which have been fixed earlier). Notice that, if $(\check{V}, \check{\Psi}) = 0$, then

$$a_r^{11}(z_0) = -a_r^{22}(z_0) = (\mu_+ - m_+)(\tilde{\tau}, \eta_0) = 0, \quad a_r^{12}(z_0) = i$$

for $\tilde{\tau} = -i(\bar{C}_2 \pm \bar{C}_1^{-1})\eta_0$. Then we can take K sufficiently small, g suitably large, and shrink \mathcal{V} to conclude (5.64). This completes the proof. \square

Thanks to Lemma 5.5, one can deduce the following lemma by using a partition of unity. We refer to [4, Theorem 9.1] and [52, §4.7.3] for a detailed derivation of the following “global” symmetrizer \mathbf{S} .

Lemma 5.7 *Assume that (4.1)–(4.2) hold for a sufficiently small positive constant K . Then there exists a mapping*

$$\mathbf{S} : \bar{\Omega} \times (\mathbb{R}^2 \times \mathbb{R}_+ \setminus \{0\}) \rightarrow \mathcal{M}_4(\mathbb{C})$$

that satisfies the following properties:

- (a) For all $z \in \bar{\Omega} \times (\mathbb{R}^2 \times \mathbb{R}_+ \setminus \{0\})$, matrix $\mathbf{S}(z)$ is Hermitian and $\mathbf{S} \in \Gamma_2^2$;
- (b) For all $z = (t, x_1, \delta, \eta, \gamma) \in \partial\Omega \times (\mathbb{R}^2 \times \mathbb{R}_+ \setminus \{0\})$,

$$\tilde{\chi}_u(z)^2 \mathbf{S}(z) + C \tilde{\chi}_u(z)^2 \lambda^{2,\gamma} \boldsymbol{\beta}(z)^* \boldsymbol{\beta}(z) \geq c \tilde{\chi}_u(z)^2 \lambda^{2,\gamma} I,$$

where $\lambda^{m,\gamma} := (\gamma^2 + \delta^2 + \eta^2)^{m/2}$;

- (c) There exists a finite set of matrix-valued mappings V_j , H_j , and E_j such that

$$\operatorname{Re}(\mathbf{S}(z) \mathbb{A}_u(z)) = \sum_j V_j(z)^* \begin{pmatrix} \gamma H_j(z) & 0 \\ 0 & E_j(z) \end{pmatrix} V_j(z),$$

where V_j and E_j belong to Γ_2^1 , $H_j \in \Gamma_2^0$, and the following estimates hold:

$$\sum_j V_j(z)^* V_j(z) \geq c \lambda^{2,\gamma} \tilde{\chi}_u(z)^2 I, \quad H_j(z) \geq cI, \quad E_j(z) \geq c \lambda^{1,\gamma} I.$$

With Lemma 5.7 in hand, we can choose \mathbf{S} as a symmetrizer for problem (5.63) to show the energy estimate as in [15, § 3.5], for which we only give the result here for brevity. We just recall that the components $T_{\chi_u}^\gamma W_1^\pm$ are given in terms of $T_{\chi_u}^\gamma W_{2,3}^\pm$ by relations similar to (5.31). The estimate for $T_{\chi_u}^\gamma W$ reads:

$$\begin{aligned} & \gamma \|T_{\chi_u}^\gamma W\|_{1,\gamma}^2 + \|T_{\chi_u}^\gamma W^{\text{nc}}|_{x_2=0}\|_{1,\gamma}^2 \\ & \lesssim \|G\|_{1,\gamma}^2 + \|W^{\text{nc}}|_{x_2=0}\|^2 + \gamma^{-1} (\|F\|_{1,\gamma}^2 + \|W\|^2 + \|T_r^\gamma W\|_{1,\gamma}^2), \end{aligned} \quad (5.80)$$

where symbol $r \in \Gamma_1^0$ vanishes on region $\{\chi_u \equiv 1\} \cup \{\chi_u \equiv 0\}$.

5.5 Estimate near the Poles

This subsection is devoted to deriving the energy estimate near poles $\Upsilon_p = \Upsilon_p^+ \cup \Upsilon_p^-$. Matrix \mathbb{A} is not defined at points in Υ_p , while the stable subspace \mathcal{E}^- of \mathbb{A} admits a continuous extension at these points, due to Lemma 5.3. We show the estimate near Υ_p^+ without loss of generality. For this purpose, we define two cut-off functions $\tilde{\chi}$ and $\tilde{\chi}_1$ with values in $[0, 1]$ that are both C^∞ and homogeneous of degree 0 with respect to (τ, η) and satisfy that

$$\tilde{\chi} \equiv 1 \quad \text{on } \operatorname{supp} \chi_p^+, \quad \tilde{\chi}_1 \equiv 1 \quad \text{on } \operatorname{supp} \tilde{\chi}, \quad \operatorname{supp} \tilde{\chi}_1 \cap \operatorname{supp} \chi_c = \emptyset, \quad (5.81)$$

where χ_p^+ and χ_c are introduced at the end of § 5.3. As in [17], we go back to the original problem (5.20) and set

$$W_p^+ := T_{\chi_p^+}^\gamma W^+, \quad W_p^- := T_{\chi_p^+}^\gamma W^-. \quad (5.82)$$

Then we employ the argument in the derivation of (5.29) to obtain

$$T_{\tau\mathbf{A}_0^+ + i\eta\mathbf{A}_1^+}^\gamma W_p^+ + T_{\mathbf{C}^+}^\gamma W_p^+ + \mathbf{I}_2 \partial_2 W_p^+ = T_r^\gamma W^+ + T_{\chi_p^+}^\gamma F^+ + \mathcal{R}_{-1} W^+, \quad (5.83)$$

where $r = i\{\chi_p^+, \tau\mathbf{A}_0^+ + i\eta\mathbf{A}_1^+\} + \partial_2 \chi_p^+ \mathbf{I}_2 \in \Gamma_1^0$. The equation for W_p^- is the same as (5.83) with index “+” replaced by “-”.

Let us introduce symbols R_+ and L_1^+ , both belonging to Γ_2^0 , such that, for $(\tau, \eta) \in \Xi_1$,

$$R_+ := \begin{pmatrix} 1 & k_1 & k_2 \\ 0 & -\dot{a}_+ a_{12}^+ & 0 \\ 0 & \dot{a}_+ (a_{11}^+ - \omega_+) & 1 \end{pmatrix}, \quad L_1^+ := \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ l_1 & l_2 & l_3 \end{pmatrix}.$$

Recalling from (5.22) and (5.24) that $(b_{ij}^+) = \tau\mathbf{A}_0^+ + i\eta\mathbf{A}_1^+$ and $b_{11}^+ = \mathbb{F}_1^+ \dot{a}_+$, we choose

$$k_1 = (\mathbb{F}_1^+)^{-1} \{b_{12}^+ a_{12}^+ - b_{13}^+ (a_{11}^+ - \omega_+)\}, \quad k_2 = -b_{23}^+ / b_{21}^+, \\ l_1 = -(\mathbb{F}_1^+)^{-1} \{b_{21}^+ (a_{11}^+ - \omega_+) + b_{31}^+ a_{12}^+\}, \quad l_2 = \dot{a}_+ (a_{11}^+ - \omega_+), \quad l_3 = \dot{a}_+ a_{12}^+,$$

so that $L_1^+ \mathbf{I}_2 R^+ = \text{diag}(0, -\dot{a}_+ a_{12}^+, \dot{a}_+ a_{12}^+)$ and

$$L_1^+ (\tau\mathbf{A}_0^+ + i\eta\mathbf{A}_1^+) R^+ = \begin{pmatrix} b_{11}^+ & 0 & k_2 b_{11}^+ + b_{13}^+ \\ b_{21}^+ & d_1 & 0 \\ 0 & d_3 & d_2 \end{pmatrix},$$

where

$$d_1 = b_{21}^+ k_1 - b_{22}^+ \dot{a}_+ a_{12}^+ + b_{23}^+ \dot{a}_+ (a_{11}^+ - \omega_+), \quad d_2 = l_j (k_2 b_{j1}^+ + b_{j3}^+) = l_j b_{j3}^+, \\ d_3 = l_j (k_1 b_{j1}^+ - b_{j2}^+ \dot{a}_+ a_{12}^+ + b_{j3}^+ \dot{a}_+ (a_{11}^+ - \omega_+)) \\ = l_j (-b_{j2}^+ \dot{a}_+ a_{12}^+ + b_{j3}^+ \dot{a}_+ (a_{11}^+ - \omega_+)).$$

We have used the relation: $l_j b_{j1}^+ = 0$. From (5.24) and definition (5.37) of \mathbb{A} , we have

$$d_1 = (\mathbb{F}_1^+)^{-1} \{ (b_{21}^+ b_{12}^+ - b_{22}^+ \mathbb{F}_1^+ \dot{a}_+) a_{12}^+ + (-b_{21}^+ b_{13}^+ + b_{23}^+ \mathbb{F}_1^+ \dot{a}_+) (a_{11}^+ - \omega_+) \} \\ = (\mathbb{F}_1^+)^{-1} \{ b_{11}^+ a_{11}^+ a_{12}^+ - b_{11}^+ a_{12}^+ (a_{11}^+ - \omega_+) \} = \dot{a}_+ a_{12}^+ \omega_+, \\ d_2 = (\mathbb{F}_1^+)^{-1} \{ (a_{11}^+ - \omega_+) (-b_{21}^+ b_{13}^+ + \mathbb{F}_1^+ \dot{a}_+ b_{23}^+) + a_{12}^+ (-b_{31}^+ b_{13}^+ + \mathbb{F}_1^+ \dot{a}_+ b_{33}^+) \} \\ = -(\mathbb{F}_1^+)^{-1} \{ (a_{11}^+ - \omega_+) b_{11}^+ a_{12}^+ + a_{12}^+ b_{11}^+ a_{22}^+ \} = \dot{a}_+ a_{12}^+ (\omega_+ - a_{11}^+ - a_{22}^+),$$

and

$$\mathbb{F}_1^+ d_3 = (a_{11}^+ - \omega_+)^2 (-b_{21}^+ \dot{a}_+ b_{13}^+ + \mathbb{F}_1^+ \dot{a}_+^2 b_{23}^+) + (a_{12}^+)^2 \dot{a}_+ (b_{12}^+ b_{31}^+ - \mathbb{F}_1^+ \dot{a}_+ b_{32}^+) \\ + (a_{11}^+ - \omega_+) a_{12}^+ \dot{a}_+ (b_{12}^+ b_{21}^+ - b_{13}^+ b_{31}^+ - \mathbb{F}_1^+ \dot{a}_+ b_{22}^+ + \mathbb{F}_1^+ \dot{a}_+ b_{33}^+) \\ = -\dot{a}_+ b_{11}^+ a_{12}^+ \{ (a_{11}^+ - \omega_+)^2 - a_{12}^+ a_{21}^+ - (a_{11}^+ - \omega_+) (a_{11}^+ - a_{22}^+) \} = 0.$$

We have used (5.42) for deriving the last identity. We notice from (3.27)–(3.28) that

$$\mathring{a}_+ a_{12}^+ |_{(\mathring{V}, \mathring{\Psi})=0} = -(\tau + i\bar{v}\eta)m_+ = \bar{\Gamma}\bar{c}(i\eta + \epsilon^2\bar{v}\tau)^2/2$$

does not vanish on $\text{supp } \tilde{\chi}_1$ by shrinking $\text{supp } \tilde{\chi}_1$ if necessary. Since $\mathring{a}_+ a_{12}^+$ is smooth in (τ, η) and $(\mathring{V}, \nabla\mathring{\Psi})$, matrix $L_2^+ := \text{diag}(1, -(\mathring{a}_+ a_{12}^+)^{-1}, (\mathring{a}_+ a_{12}^+)^{-1})$ is a smooth and invertible mapping on $\text{supp } \tilde{\chi}_1$. We then derive

$$L_+ \mathbf{I}_2 R_+ = \mathbf{I}_2, \quad L_+ := L_2^+ L_1^+, \quad (5.84)$$

$$\begin{aligned} & L_+(\tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+) R_+ \\ &= \begin{pmatrix} b_{11}^+ & 0 & k_2 b_{11}^+ + b_{13}^+ \\ -b_{21}^+ / (\mathring{a}_+ a_{12}^+) & -\omega_+ & 0 \\ 0 & 0 & \omega_+ - a_{11}^+ - a_{22}^+ \end{pmatrix} =: A_+^d. \end{aligned} \quad (5.85)$$

We also introduce symbols R_- and L_1^- that belong to Γ_2^0 and satisfy that, for $(\tau, \eta) \in \Xi_1$,

$$\begin{aligned} R_- &:= \begin{pmatrix} 1 & k_1^- & k_2^- \\ 0 & \mathring{a}_-(a_{22}^- - \omega_-) & \mathring{a}_- a_{12}^- \\ 0 & -\mathring{a}_- a_{21}^- & \mathring{a}_-(-a_{11}^- + \omega'_-) \end{pmatrix}, \\ L_1^- &:= \begin{pmatrix} 1 & 0 & 0 \\ l_1^- & \mathring{a}_-(a_{11}^- - \omega'_-) & \mathring{a}_- a_{12}^- \\ l_2^- & \mathring{a}_- a_{21}^- & \mathring{a}_-(a_{22}^- - \omega_-) \end{pmatrix}, \end{aligned}$$

with

$$\begin{aligned} k_1^- &= (\mathbb{F}_1^-)^{-1} \{-(a_{22}^- - \omega_-)b_{12}^- + a_{21}^- b_{13}^-\}, \\ k_2^- &= (\mathbb{F}_1^-)^{-1} \{-a_{12}^- b_{12}^- + (a_{11}^- - \omega'_-)b_{13}^-\}, \\ l_1^- &= (\mathbb{F}_1^-)^{-1} \{-(a_{11}^- - \omega'_-)b_{21}^- - a_{12}^- b_{31}^-\}, \\ l_2^- &= (\mathbb{F}_1^-)^{-1} \{-a_{21}^- b_{21}^- - (a_{22}^- - \omega_-)b_{31}^-\}, \end{aligned}$$

so that $L_1^- \mathbf{I}_2 R_- = \text{diag}(0, d_4, -d_4)$ and

$$L_1^-(\tau \mathbf{A}_0^- + i\eta \mathbf{A}_1^-) R_- = \begin{pmatrix} b_{11}^- & 0 & 0 \\ 0 & d_5 & d_6 \\ 0 & d_7 & d_8 \end{pmatrix},$$

where $d_4 = \mathring{a}_-^2 \{(a_{11}^- - \omega'_-)(a_{22}^- - \omega_-) - a_{12}^- a_{21}^-\}$ and

$$\begin{aligned} d_5 &= \mathring{a}_-(a_{11}^- - \omega'_-) \{b_{21}^- k_1^- + b_{22}^- \mathring{a}_-(a_{22}^- - \omega_-) - b_{23}^- \mathring{a}_- a_{21}^-\} \\ &\quad + \mathring{a}_- a_{12}^- \{b_{31}^- k_1^- + b_{32}^- \mathring{a}_-(a_{22}^- - \omega_-) - b_{33}^- \mathring{a}_- a_{21}^-\}. \end{aligned}$$

We omit the expressions for d_j , $j \in \{6, 7, 8\}$, since they are quite similar to that of d_5 . By virtue of the identity: $a_{22}^- - \omega_- = -a_{11}^- + \omega'_-$ and (5.42), we deduce

$$\begin{aligned} (a_{11}^- - \omega'_-)(a_{22}^- - \omega_-) - a_{12}^- a_{21}^- &= -(a_{22}^- - \omega_-)^2 - a_{12}^- a_{21}^- \\ &= (\omega_- - a_{22}^-)(-2\omega_- + a_{11}^- + a_{22}^-), \end{aligned} \quad (5.86)$$

which yields $d_4 = \hat{a}_-^2 (\omega_- - a_{22}^-)(-2\omega_- + a_{11}^- + a_{22}^-)$. Using (5.22), (5.24), and the identity: $a_{22}^- - \omega_- = -a_{11}^- + \omega'_-$, we compute

$$d_5 = \hat{a}_-^2 \{a_{11}^- (a_{22}^- - \omega_-)^2 - 2a_{12}^- a_{21}^- (a_{22}^- - \omega_-) + a_{12}^- a_{21}^- a_{22}^-\},$$

which, combined with (5.42) and (5.86), implies

$$d_5 = \hat{a}_-^2 \omega_- (\omega_- - a_{22}^-)(2\omega_- - a_{11}^- - a_{22}^-) = -\omega_- d_4.$$

Performing the similar calculations to d_j for $j = 6, 7, 8$, we can discover that $d_6 = d_7 = 0$ and $d_8 = \omega'_- d_4$ so that

$$L_1^- (\tau \mathbf{A}_0^- + i\eta \mathbf{A}_1^-) R_- = \text{diag}(0, -\omega_- d_4, \omega'_- d_4).$$

Note that d_4 does not vanish in neighborhood $\text{supp } \tilde{\chi}_1$ of \mathcal{Y}_p^+ up to shrinking $\text{supp } \tilde{\chi}_1$. Setting $L_2^- := \text{diag}(1, d_4^{-1}, -d_4^{-1})$ and $L_- := L_2^- L_1^-$, we obtain

$$\begin{cases} L_- \mathbf{I}_2 R_- = \mathbf{I}_2, \\ L_- (\tau \mathbf{A}_0^- + i\eta \mathbf{A}_1^-) R_- = \text{diag}(b_{11}^-, -\omega_-, -\omega'_-) =: A_-^d. \end{cases} \quad (5.87)$$

Let us define

$$Z^+ := T_{\tilde{\chi} R_+}^\gamma W_p^+, \quad Z^- := T_{\tilde{\chi} R_-}^\gamma W_p^-.$$

Applying operator $T_{\tilde{\chi} L_+}^\gamma$ to (5.83) and using (5.84) yield

$$\begin{aligned} T_{\tilde{\chi} L_+}^\gamma (\tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+) W_p^+ &+ T_{-i \sum_{j=0}^1 \partial_{\xi_j} (\tilde{\chi} L_+) \partial_{x_j} (\tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+)}^\gamma W_p^+ + T_{\tilde{\chi} L_+ C^+}^\gamma W_p^+ \\ &= -T_{\tilde{\chi} L_+ \mathbf{I}_2}^\gamma \partial_2 W_p^+ + T_{\tilde{\chi} L_+}^\gamma T_r^\gamma W^+ + T_{\tilde{\chi} L_+}^\gamma T_{\chi_p^+}^\gamma F^+ + \mathcal{R}_{-1} W^+ \\ &= -\mathbf{I}_2 \partial_2 Z^+ + T_{\mathbf{I}_2 \partial_2 (\tilde{\chi} R_+^{-1})}^\gamma W_p^+ + T_{\tilde{\chi} L_+}^\gamma T_r^\gamma W^+ + T_{\tilde{\chi} L_+}^\gamma T_{\chi_p^+}^\gamma F^+ + \mathcal{R}_{-1} W^+, \end{aligned}$$

where $x_0 := t$, $\xi_0 := \delta$, and $\xi_1 := \eta$ to avoid overloaded equations. On the other hand, it follows from (5.82) and (5.85) that

$$T_{A_+^d}^\gamma Z^+ = T_{\tilde{\chi} L_+ (\tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+)}^\gamma W_p^+ + T_{-i \sum_{j=0}^1 \partial_{\xi_j} A_+^d \partial_{x_j} (\tilde{\chi} R_+^{-1})}^\gamma W_p^+ + \mathcal{R}_{-1} W^+.$$

Then we have

$$\mathbf{I}_2 \partial_2 Z^+ + T_{A_+^d}^\gamma Z^+ + T_{\mathbb{D}_+^p}^\gamma Z^+ = T_r^\gamma W^+ + \mathcal{R}_0 T_{\chi_p^+}^\gamma F^+ + \mathcal{R}_{-1} W^+ \quad (5.88)$$

for new $r \in \Gamma_1^0$ vanishing on $\{\chi_p^+ \equiv 1\} \cup \{\chi_p^+ \equiv 0\}$, where \tilde{A}_+^d is an extension of A_+^d to the whole set $\bar{\Omega} \times \Xi$, and $\mathbb{D}_+^p \in \Gamma_1^0$ is given by

$$\begin{aligned} \mathbb{D}_+^p &:= \tilde{\chi} L_+ C^+ R_+ - \mathbf{I}_2 \partial_2 (\tilde{\chi} R_+^{-1}) R_+ \\ &\quad + i \sum_{j=0,1} \{ \partial_{\xi_j} \tilde{A}_+^d \partial_{x_j} (\tilde{\chi} R_+^{-1}) - \partial_{\xi_j} (\tilde{\chi} L_+) \partial_{x_j} (\tau \mathbf{A}_0^+ + i\eta \mathbf{A}_1^+) \} R_+. \end{aligned}$$

Similarly, we have

$$\mathbf{I}_2 \partial_2 Z^- + T_{\tilde{A}_-^d}^\gamma Z^- + T_{\mathbb{D}_-^p}^\gamma Z^- = T_r^\gamma W^- + \mathcal{R}_0 T_{\chi_p^+}^\gamma F^- + \mathcal{R}_{-1} W^-, \quad (5.89)$$

where $r \in \Gamma_1^0$ vanishes on $\{\chi_p^+ \equiv 1\} \cup \{\chi_p^+ \equiv 0\}$, \tilde{A}_-^d is an extension to the whole set $\bar{\Omega} \times \Xi$ of A_-^d , and $\mathbb{D}_-^p \in \Gamma_1^0$. According to the definitions of R_\pm , we have

$$Z^{\text{nc}} = T_{\tilde{\chi} \tilde{R}^{-1}}^\gamma T_{\chi_p^+}^\gamma W^{\text{nc}}, \quad (5.90)$$

where $Z^{\text{nc}} := (Z_2^+, Z_3^+, Z_2^-, Z_3^-)^\top$ and

$$\tilde{R} := \begin{pmatrix} -\dot{a}_+ a_{12}^+ & 0 & 0 & 0 \\ \dot{a}_+ (a_{11}^+ - \omega_+) & 1 & 0 & 0 \\ 0 & 0 & \dot{a}_- (a_{22}^- - \omega_-) & \dot{a}_- a_{12}^- \\ 0 & 0 & -\dot{a}_- a_{21}^- & \dot{a}_- (-a_{11}^- + \omega_-') \end{pmatrix}.$$

Note from (5.39) that the first and third columns of \tilde{R} are E_+ and E_- . By virtue of (5.20c), we obtain the following boundary conditions in terms of Z^{nc} :

$$T_{\beta(E_+ E_-)}^\gamma \begin{pmatrix} Z_2^+ \\ Z_2^- \end{pmatrix} + \mathcal{R}_0 \begin{pmatrix} Z_3^+ \\ Z_3^- \end{pmatrix} = \mathcal{R}_0 G + \mathcal{R}_{-1} W^{\text{nc}} \quad \text{if } x_2 = 0. \quad (5.91)$$

For problem (5.88)–(5.89) and (5.91), we obtain the following energy estimate:

Lemma 5.8 *There exists constants $K_0 \leq 1$ and $\gamma_0 \geq 1$ such that, if $\gamma \geq \gamma_0$ and $K \leq K_0$ for K given in (4.2), then*

$$\begin{aligned} &\gamma \|Z^\pm\|_{1,\gamma}^2 + \|(Z_2^\pm, Z_3^\pm)|_{x_2=0}\|_{1,\gamma}^2 \\ &\lesssim \|G\|_{1,\gamma}^2 + \|W^{\text{nc}}|_{x_2=0}\|^2 + \gamma^{-1} (\|T_{\chi_p^+}^\gamma F\|_{1,\gamma}^2 + \|W\|^2 + \|T_r^\gamma W\|_{1,\gamma}^2), \end{aligned} \quad (5.92)$$

where symbol $r \in \Gamma_1^0$ vanishes in region $\{\chi_p^+ \equiv 1\} \cup \{\chi_p^+ \equiv 0\}$.

Proof We divide the proof into five steps.

1. Estimate for Z_3^+ . According to the form of A_+^d given by (5.85), the third equation in (5.88) for Z_3^+ reads

$$\partial_2 Z_3^+ = T_{-\omega_+ + a_{11}^+ + a_{22}^+}^\gamma Z_3^+ + T_{\alpha_0}^\gamma Z^+ + T_r^\gamma W^+ + \mathcal{R}_0 T_{\chi_p^+}^\gamma F_3^+ + \mathcal{R}_{-1} W^+. \quad (5.93)$$

Take the scalar product in $L^2(\Omega)$ of (5.93) with $A^{2,\gamma}Z_3^+$ to obtain

$$\|Z_3^+|_{x_2=0}\|_{1,\gamma}^2 + 2 \operatorname{Re} \langle A^{1,\gamma}Z_3^+, A^{1,\gamma}T_{-\omega_+ + a_{11}^+ + a_{22}^+}^\gamma Z_3^+ \rangle = \sum_{j=1}^4 \mathcal{H}_j, \quad (5.94)$$

where each term \mathcal{H}_j in the decomposition is defined in the following:

$$\begin{aligned} \mathcal{H}_1 &:= -2 \operatorname{Re} \langle A^{1,\gamma}Z_3^+, A^{1,\gamma}T_{\alpha_0}^\gamma Z^+ \rangle \lesssim \|Z^+\|_{1,\gamma}^2, \\ \mathcal{H}_2 &:= -2 \operatorname{Re} \langle A^{1,\gamma}Z_3^+, A^{1,\gamma}T_r^\gamma W^+ \rangle \lesssim \varepsilon \gamma \|Z_3^+\|_{1,\gamma}^2 + \frac{1}{\varepsilon \gamma} \|T_r^\gamma W^+\|_{1,\gamma}^2, \\ \mathcal{H}_3 &:= -2 \operatorname{Re} \langle A^{1,\gamma}Z_3^+, A^{1,\gamma}\mathcal{R}_0 T_{\chi_p^+}^\gamma F_3^+ \rangle \lesssim \varepsilon \gamma \|Z_3^+\|_{1,\gamma}^2 + \frac{1}{\varepsilon \gamma} \|T_{\chi_p^+}^\gamma F_3^+\|_{1,\gamma}^2, \\ \mathcal{H}_4 &:= -2 \operatorname{Re} \langle A^{1,\gamma}Z_3^+, A^{1,\gamma}\mathcal{R}_{-1} W^+ \rangle \lesssim \varepsilon \gamma \|Z_3^+\|_{1,\gamma}^2 + \frac{1}{\varepsilon \gamma} \|W^+\|^2. \end{aligned}$$

For the second term on the left-hand side of (5.94), we employ Lemma 5.2 (iv) to deduce

$$\begin{aligned} &\operatorname{Re} \langle A^{1,\gamma}Z_3^+, A^{1,\gamma}T_{-\omega_+ + a_{11}^+ + a_{22}^+}^\gamma Z_3^+ \rangle \\ &\geq \operatorname{Re} \langle A^{1,\gamma}Z_3^+, T_{-\omega_+ + a_{11}^+ + a_{22}^+}^\gamma A^{1,\gamma}Z_3^+ \rangle - C \|Z_3^+\|_{1,\gamma}^2. \end{aligned}$$

Thanks to (5.49), $\operatorname{Re}(a_{11}^+ + a_{22}^+) = \mathbb{F}_4^+ \gamma$, where \mathbb{F}_4^+ is a smooth function of $(\mathring{V}, \nabla \mathring{\Psi})$ that vanishes at the origin. We then employ Lemma 5.6 and take K in (4.2) sufficiently small to obtain that $\operatorname{Re}(-\omega_+ + a_{11}^+ + a_{22}^+) \gtrsim \gamma$. Apply Gårding's inequality (Lemma 5.2 (vi)) to obtain

$$\operatorname{Re} \langle A^{1,\gamma}Z_3^+, T_{-\omega_+ + a_{11}^+ + a_{22}^+}^\gamma A^{1,\gamma}Z_3^+ \rangle \gtrsim \gamma \|A^{1,\gamma}Z_3^+\|^2 \gtrsim \gamma \|Z_3^+\|_{1,\gamma}^2,$$

from which we have

$$\operatorname{Re} \langle A^{1,\gamma}Z_3^+, A^{1,\gamma}T_{-\omega_+ + a_{11}^+ + a_{22}^+}^\gamma Z_3^+ \rangle \gtrsim (\gamma - C) \|Z_3^+\|_{1,\gamma}^2.$$

Choosing ε small and γ large, we derive from (5.94) that

$$\begin{aligned} &\gamma \|Z_3^+\|_{1,\gamma}^2 + \|Z_3^+|_{x_2=0}\|_{1,\gamma}^2 \\ &\lesssim \|Z^+\|_{1,\gamma}^2 + \gamma^{-1} (\|T_r^\gamma W^+\|_{1,\gamma}^2 + \|T_{\chi_p^+}^\gamma F_3^+\|_{1,\gamma}^2 + \|W^+\|^2). \end{aligned} \quad (5.95)$$

2. Estimate for Z_1^+ . The equation for Z_1^+ in (5.88) is as follows:

$$T_{b_{11}^+}^\gamma Z_1^+ + T_{k_2 b_{11}^+ + b_{13}^+}^\gamma Z_3^+ = T_{\alpha_0}^\gamma Z^+ + T_r^\gamma W^+ + \mathcal{R}_0 T_{\chi_p^+}^\gamma F_1^+ + \mathcal{R}_{-1} W^+. \quad (5.96)$$

Recall from (5.24) that $\operatorname{Re} b_{11}^+ = \mathbb{F}_1^+ \gamma$ and $\operatorname{Re}(k_2 b_{11}^+ + b_{13}^+) = \gamma \alpha_0$. Similar to Step 1, we take the scalar product in $L^2(\Omega)$ of (5.96) with $A^{2,\gamma}Z_1^+$ and use (5.95) to obtain

$$\begin{aligned} &\gamma \|Z_1^+\|_{1,\gamma}^2 \\ &\lesssim \|Z^+\|_{1,\gamma}^2 + \gamma \|Z_3^+\|_{1,\gamma}^2 + \gamma^{-1} (\|T_r^\gamma W^+\|_{1,\gamma}^2 + \|T_{\chi_p^+}^\gamma F_1^+\|_{1,\gamma}^2 + \|W^+\|^2) \\ &\lesssim \|Z^+\|_{1,\gamma}^2 + \gamma^{-1} (\|T_r^\gamma W^+\|_{1,\gamma}^2 + \|T_{\chi_p^+}^\gamma F_1^+\|_{1,\gamma}^2 + \|W^+\|^2) \end{aligned} \quad (5.97)$$

for γ sufficiently large.

3. Estimate for Z_2^+ . The equation for Z_2^+ in (5.88) reads

$$\begin{aligned} \partial_2 Z_2^+ &= T_{(\tilde{a}_+ a_{12}^+)^{-1} b_{21}^+}^\gamma Z_1^+ + T_{\omega_+}^\gamma Z_2^+ + T_{\alpha_0}^\gamma Z^+ \\ &\quad + T_r^\gamma W^+ + \mathcal{R}_0 T_{\chi_p^+}^\gamma F_2^+ + \mathcal{R}_{-1} W^+. \end{aligned}$$

We note that $\operatorname{Re} \omega_+ \lesssim -\gamma$ and $\operatorname{Re}((\tilde{a}_+ a_{12}^+)^{-1} b_{21}^+) = \gamma \alpha_0$. Employing a similar analysis as Step 1 and using (5.97) yield

$$\begin{aligned} &\gamma \|Z_2^+\|_{1,\gamma}^2 - C \|Z_2^+|_{x_2=0}\|_{1,\gamma}^2 \\ &\lesssim \|Z^+\|_{1,\gamma}^2 + \gamma \|Z_1^+\|_{1,\gamma}^2 + \gamma^{-1} (\|T_r^\gamma W^+\|_{1,\gamma}^2 + \|T_{\chi_p^+}^\gamma F_2^+\|_{1,\gamma}^2 + \|W^+\|^2) \\ &\lesssim \|Z^+\|_{1,\gamma}^2 + \gamma^{-1} (\|T_r^\gamma W^+\|_{1,\gamma}^2 + \|T_{\chi_p^+}^\gamma F^+\|_{1,\gamma}^2 + \|W^+\|^2). \end{aligned} \quad (5.98)$$

4. Combine estimates (5.95) and (5.97)–(5.98), and take γ suitably large to find

$$\begin{aligned} &\gamma \|Z^+\|_{1,\gamma}^2 + \|Z_3^+|_{x_2=0}\|_{1,\gamma}^2 \\ &\lesssim \|Z_2^+|_{x_2=0}\|_{1,\gamma}^2 + \gamma^{-1} (\|T_r^\gamma W^+\|_{1,\gamma}^2 + \|T_{\chi_p^+}^\gamma F^+\|_{1,\gamma}^2 + \|W^+\|^2). \end{aligned} \quad (5.99)$$

The derivation for the estimate of Z^- is entirely similar so that

$$\begin{aligned} &\gamma \|Z^\pm\|_{1,\gamma}^2 + \|Z_3^\pm|_{x_2=0}\|_{1,\gamma}^2 \\ &\lesssim \|Z_2^\pm|_{x_2=0}\|_{1,\gamma}^2 + \gamma^{-1} (\|T_r^\gamma W\|_{1,\gamma}^2 + \|T_{\chi_p^\pm}^\gamma F\|_{1,\gamma}^2 + \|W\|^2). \end{aligned} \quad (5.100)$$

5. Estimate on the boundary. It remains to make an estimate for $\|Z_2^\pm|_{x_2=0}\|_{1,\gamma}$. Using the boundary conditions (5.91), we have

$$\|T_{\tilde{\beta}}^\gamma Z_2|_{x_2=0}\|_{1,\gamma}^2 \lesssim \|Z_3^\pm|_{x_2=0}\|_{1,\gamma}^2 + \|G\|_{1,\gamma}^2 + \|W^{\text{nc}}|_{x_2=0}\|^2, \quad (5.101)$$

where $\tilde{\beta} := \beta(E_+ E_-)$ and $Z_2 := (Z_2^+, Z_2^-)^\top$. Setting $V^\pm := T_{\tilde{R}^{-1} \chi_p^+}^\gamma W_\pm^{\text{nc}}$, we see from (5.90) that

$$Z_2^\pm = T_{\tilde{\chi}}^\gamma V_1^\pm + \mathcal{R}_{-1} W^{\text{nc}}. \quad (5.102)$$

Since $\tilde{\beta} \in \Gamma_2^0$, we apply the rule of symbolic calculus (Lemma 5.2 (iv)) to find that

$$T_{\tilde{\beta}}^\gamma \mathcal{R}_{-1} = \mathcal{R}_{-1}, \quad (A^{1,\gamma} T_{\tilde{\beta}}^\gamma)^* A^{1,\gamma} T_{\tilde{\beta}}^\gamma - T_{\lambda^{2,\gamma} \tilde{\beta}^* \tilde{\beta}}^\gamma = \mathcal{R}_{-1}.$$

Thus, we have

$$\begin{aligned} &\|T_{\tilde{\beta}}^\gamma Z_2|_{x_2=0}\|_{1,\gamma}^2 \\ &\gtrsim \operatorname{Re} \langle (A^{1,\gamma} T_{\tilde{\beta}}^\gamma)^* A^{1,\gamma} T_{\tilde{\beta}}^\gamma T_{\tilde{\chi}}^\gamma V_1|_{x_2=0}, T_{\tilde{\chi}}^\gamma V_1|_{x_2=0} \rangle - C \|W^{\text{nc}}|_{x_2=0}\|^2 \\ &\gtrsim \operatorname{Re} \langle T_{\lambda^{2,\gamma} \tilde{\beta}^* \tilde{\beta}}^\gamma T_{\tilde{\chi}}^\gamma V_1|_{x_2=0}, T_{\tilde{\chi}}^\gamma V_1|_{x_2=0} \rangle \\ &\quad - C \|T_{\tilde{\chi}}^\gamma V_1|_{x_2=0}\|_{1,\gamma} \|T_{\tilde{\chi}}^\gamma V_1|_{x_2=0}\| - C \|W^{\text{nc}}|_{x_2=0}\|^2 \end{aligned}$$

for $V_1 := (V_1^+, V_1^-)^\top$, which, combined with (5.101), implies

$$\begin{aligned} & \operatorname{Re} \langle T_{\lambda^2, \gamma}^\gamma \tilde{\beta}^* \tilde{\beta} T_{\tilde{\chi}}^\gamma V_1|_{x_2=0}, T_{\tilde{\chi}}^\gamma V_1|_{x_2=0} \rangle \\ & \lesssim \|Z_3^\pm|_{x_2=0}\|_{1, \gamma}^2 + \|G\|_{1, \gamma}^2 + \gamma^{-1} \|Z_2^\pm|_{x_2=0}\|_{1, \gamma}^2 + \|W^{\text{nc}}|_{x_2=0}\|^2. \end{aligned} \quad (5.103)$$

Recall from Lemma 5.4 that the Lopatinskiĭ determinant Δ does not vanish on $\operatorname{supp} \tilde{\chi}_1$, owing to (5.81). It then follows from definition (5.50) of Δ that

$$\tilde{\chi}_1^2 \operatorname{Re}(\lambda^{2, \gamma} \tilde{\beta}^* \tilde{\beta}) \gtrsim \tilde{\chi}_1^2 \lambda^{2, \gamma} I.$$

Then we can employ the localized Gårding's inequality (Lemma 5.2 (vii)) and utilize (5.102) to derive

$$\begin{aligned} & \operatorname{Re} \langle T_{\lambda^2, \gamma}^\gamma \tilde{\beta}^* \tilde{\beta} T_{\tilde{\chi}}^\gamma V_1|_{x_2=0}, T_{\tilde{\chi}}^\gamma V_1|_{x_2=0} \rangle \\ & \gtrsim \|T_{\tilde{\chi}}^\gamma V_1|_{x_2=0}\|_{1, \gamma}^2 - C \|V_1|_{x_2=0}\|^2 \gtrsim \|Z_2|_{x_2=0}\|_{1, \gamma}^2 - C \|W^{\text{nc}}|_{x_2=0}\|^2. \end{aligned} \quad (5.104)$$

Combine (5.103) with (5.104) and take γ small to infer that

$$\|Z_2^\pm|_{x_2=0}\|_{1, \gamma}^2 \lesssim \|Z_3^\pm|_{x_2=0}\|_{1, \gamma}^2 + \|G\|_{1, \gamma}^2 + \|W^{\text{nc}}|_{x_2=0}\|^2. \quad (5.105)$$

We combine (5.105) with (5.95) to eliminate the first term on the right hand side of (5.105), and then use (5.100) to conclude estimate (5.92). This completes the proof. \square

Recall that $\chi_p = \chi_p^+ + \chi_p^-$ and $\operatorname{supp} \chi_p^+ \cap \operatorname{supp} \chi_p^- = \emptyset$. Shrinking the support of χ_p if necessary, we obtain the following result from Lemma 5.8.

Proposition 5.4 *There exist constants $K_0 \leq 1$ and $\gamma_0 \geq 1$ such that, if $\gamma \geq \gamma_0$ and $K \leq K_0$ for K given in (4.2), then*

$$\begin{aligned} & \gamma \| \|T_{\chi_p}^\gamma W\|_{1, \gamma}^2 + \|T_{\chi_p}^\gamma W^{\text{nc}}|_{x_2=0}\|_{1, \gamma}^2 \\ & \lesssim \|G\|_{1, \gamma}^2 + \|W^{\text{nc}}|_{x_2=0}\|^2 + \gamma^{-1} (\| \|T_{\chi_p}^\gamma F\|_{1, \gamma}^2 + \| \|W\|^2 + \| \|T_r^\gamma W\|_{1, \gamma}^2), \end{aligned} \quad (5.106)$$

where symbol $r \in \Gamma_1^0$ vanishes in region $\{\chi_p \equiv 1\} \cup \{\chi_p \equiv 0\}$.

5.6 Estimate near Bad Frequencies

We now show the energy estimate near the points in $\Upsilon_c = \cup_{q \in \{0, \pm 1\}} \Upsilon_c^q$, i.e. near the zeros of the Lopatinskiĭ determinant. We consider the case near set Υ_c^0 , without loss of generality. To this end, we introduce two smooth cut-off functions χ_1 and χ_2 with values in $[0, 1]$ such that

- $\chi_1 \equiv 1$ on the support of χ_c^0 , $\chi_2 \equiv 1$ on the support of χ_c^0 , and $\operatorname{supp} \chi_2 \subset \mathcal{V}_c^0$;
- χ_1 and χ_2 are both C^∞ and homogeneous of degree 0 with respect to (τ, η) ,

where χ_c^0 is given at the end of § 5.3. Defining

$$\mathbf{w}^\pm := T_{\chi_c^0}^\gamma W^\pm, \quad \mathbf{w}_\pm^{\text{nc}} := (\mathbf{w}_2^\pm, \mathbf{w}_3^\pm)^\top,$$

we perform similar calculations as we have done in § 5.2 to obtain the following system:

$$\partial_2 \mathbf{w}_\pm^{\text{nc}} = T_{\mathbb{A}_{\chi_2}^\pm}^\gamma \mathbf{w}_\pm^{\text{nc}} + T_{\mathbb{E}^\pm}^\gamma \mathbf{w}_\pm^{\text{nc}} + T_r^\gamma W^\pm + \mathcal{R}_0 T_{\chi_c^0}^\gamma F^\pm + \mathcal{R}_{-1} W^\pm, \quad (5.107)$$

where $\mathbb{E}^\pm \in \Gamma_1^0$, $\mathbb{A}_{\chi_2}^\pm \in \Gamma_2^1$ is given in (5.36) with $\tilde{\chi}_\pm$ replaced by χ_2 , and $r \in \Gamma_1^0$ vanishes in region $\{\chi_c^0 \equiv 1\} \cup \{\chi_c^0 \equiv 0\}$.

Since matrix $\mathbb{A}_{\chi_2}^\pm \equiv \mathbb{A}_\pm$ in region $\{\chi_2 \equiv 1\}$, we obtain from Proposition 5.3 that

$$Q_0^\pm \mathbb{A}_{\chi_2}^\pm = \mathbb{D}_1^\pm Q_0^\pm \quad \text{in } \{\chi_2 \equiv 1\}. \quad (5.108)$$

More precisely, we have

$$\begin{pmatrix} Q_0^+ & 0 \\ 0 & Q_0^- \end{pmatrix}^{-1} = (E_+ \ Y_+ \ E_- \ Y_-), \quad (5.109)$$

where E_\pm , Y_+ , and Y_- are defined by (5.39) and (5.72)–(5.73), respectively. Then the following lemma can be proved as in [15, Page 425] by using (5.108).

Lemma 5.9 *There exist symbols $Q_{-1}^\pm \in \Gamma_1^{-1}$ and diagonal symbols $\mathbb{D}_0^\pm \in \Gamma_1^0$, which are defined in region $\{\chi_2 \equiv 1\}$, such that*

$$\begin{aligned} & (Q_0^\pm + Q_{-1}^\pm)(\mathbb{A}_{\chi_2}^\pm + \mathbb{E}^\pm) - (\mathbb{D}_1^\pm + \mathbb{D}_0^\pm)(Q_0^\pm + Q_{-1}^\pm) + \partial_2 Q_0^\pm \\ & - i(\partial_\delta Q_0^\pm \partial_t \mathbb{A}_{\chi_2}^\pm + \partial_\eta Q_0^\pm \partial_{x_1} \mathbb{A}_{\chi_2}^\pm - \partial_\delta \mathbb{D}_1^\pm \partial_t Q_0^\pm - \partial_\eta \mathbb{D}_1^\pm \partial_{x_1} Q_0^\pm) \in \Gamma_1^{-1}. \end{aligned}$$

We now prove the estimates for

$$Z^\pm := T_{\chi_1(Q_0^\pm + Q_{-1}^\pm)}^\gamma \mathbf{w}_\pm^{\text{nc}}, \quad (5.110)$$

which will be shown to satisfy the paradifferential equations with diagonal principle symbols.

In fact, using Lemmas 5.2 and 5.9, we see from (5.107) that

$$\partial_2 Z^+ = T_{\tilde{\mathbb{D}}_1^+}^\gamma Z^+ + T_{\mathbb{D}_0^+}^\gamma Z^+ + T_r^\gamma W^+ + \mathcal{R}_0 T_{\chi_c^0}^\gamma F^+ + \mathcal{R}_{-1} W^+, \quad (5.111)$$

where $\tilde{\mathbb{D}}_1^+$ (resp. $\tilde{\mathbb{D}}_0^+$) is an extension of \mathbb{D}_1^+ (resp. \mathbb{D}_0^+) to the whole set $\bar{\Omega} \times \Xi$. Thanks to Lemma 5.6, these extension can be chosen such that

$$\tilde{\mathbb{D}}_1^+ = \begin{pmatrix} \omega_+ & 0 \\ 0 & \omega'_+ \end{pmatrix} = \begin{pmatrix} \gamma e_+ + i\hbar_+ & 0 \\ 0 & \gamma e'_+ + i\hbar'_+ \end{pmatrix}, \quad \tilde{\mathbb{D}}_0^+ = \text{diag}(d_+, d'_+), \quad (5.112)$$

where $e_+, e'_+ \in \Gamma_2^0$ and $\hbar_+, \hbar'_+ \in \Gamma_2^1$ are real-valued symbols, and $d_+, d'_+ \in \Gamma_1^0$ such that

$$e_+ \lesssim -1, \quad e'_+ \gtrsim 1.$$

We obtain the following result for functions Z^\pm that are given in (5.110):

Lemma 5.10 *There exist constants $K_0 \leq 1$ and $\gamma_0 \geq 1$ such that, if $\gamma \geq \gamma_0$ and $K \leq K_0$ for K given in (4.2), then*

$$\begin{aligned} & \gamma^3 \|Z_1\|^2 + \gamma \|Z_2\|_{1,\gamma}^2 + \|Z_2|_{x_2=0}\|_{1,\gamma}^2 + \gamma^2 \|Z_1|_{x_2=0}\|^2 + \|T_{\tilde{\sigma}_0}^\gamma Z_1|_{x_2=0}\|^2 \\ & \lesssim \gamma^{-1} (\|T_r^\gamma W\|_{1,\gamma}^2 + \|T_{\chi_c^0}^\gamma F\|_{1,\gamma}^2 + \|W\|^2) + \|G\|_{1,\gamma}^2 + \|W^{\text{nc}}|_{x_2=0}\|^2, \end{aligned} \quad (5.113)$$

where $Z_j := (Z_j^+, Z_j^-)^\top$, $j = 1, 2$, $\tilde{\sigma}_0$ is the scalar real symbol given by (5.59), and $r \in \Gamma_1^0$ vanishes in region $\{\chi_c^0 \equiv 1\} \cup \{\chi_c^0 \equiv 0\}$.

Proof The proof is divided into two steps.

1. Estimate in domain Ω . The first equation in (5.111) reads

$$\partial_2 Z_1^+ = T_{\omega_+}^\gamma Z_1^+ + T_{d_+}^\gamma Z_1^+ + T_r^\gamma W^+ + \mathcal{R}_0 T_{\chi_c^0}^\gamma F^+ + \mathcal{R}_{-1} W^+. \quad (5.114)$$

Recalling that $\text{Re } \omega_+ = \gamma e_+ \lesssim -\gamma$, we choose the identity as a symmetrizer and obtain the following L^2 estimate:

$$\begin{aligned} & \gamma \|Z_1^+\|^2 \lesssim \|Z_1^+|_{x_2=0}\|^2 + \gamma^{-1} \| (T_r^\gamma W^+, \mathcal{R}_0 T_{\chi_c^0}^\gamma F^+, \mathcal{R}_{-1} W^+) \|^2 \\ & \lesssim \|Z_1^+|_{x_2=0}\|^2 + \gamma^{-3} (\|T_r^\gamma W^+\|_{1,\gamma}^2 + \|T_{\chi_c^0}^\gamma F^+\|_{1,\gamma}^2 + \|W^+\|^2) \end{aligned} \quad (5.115)$$

for sufficiently large γ .

The second equation in (5.111) reads

$$\partial_2 Z_2^+ = T_{\omega_+}^\gamma Z_2^+ + T_{d_+}^\gamma Z_2^+ + T_r^\gamma W^+ + \mathcal{R}_0 T_{\chi_c^0}^\gamma F^+ + \mathcal{R}_{-1} W^+.$$

Recalling that $\text{Re } \omega_+' = \gamma e_+' \gtrsim \gamma$, we perform a similar calculation as Step 1 in the proof of Lemma 5.8 to deduce

$$\begin{aligned} & \gamma \|Z_2^+\|_{1,\gamma}^2 + \|Z_2^+|_{x_2=0}\|_{1,\gamma}^2 \\ & \lesssim \gamma^{-1} (\|T_r^\gamma W^+\|_{1,\gamma}^2 + \|T_{\chi_c^0}^\gamma F^+\|_{1,\gamma}^2 + \|W^+\|^2) \end{aligned} \quad (5.116)$$

for sufficiently large γ .

A similar analysis enables us to deduce the energy estimates for Z_1^- and Z_2^- as (5.115) and (5.116). The combination of all these estimates is

$$\begin{aligned} & \gamma^3 \|Z_1\|^2 + \gamma \|Z_2\|_{1,\gamma}^2 + \|Z_2|_{x_2=0}\|_{1,\gamma}^2 \\ & \lesssim \gamma^2 \|Z_1|_{x_2=0}\|^2 + \gamma^{-1} (\|T_r^\gamma W\|_{1,\gamma}^2 + \|T_{\chi_c^0}^\gamma F\|_{1,\gamma}^2 + \|W\|^2). \end{aligned} \quad (5.117)$$

2. Estimate for the boundary terms. We now estimate the traces of the incoming modes Z_1 in terms of the outgoing models Z_2 and the source term G . Using the boundary condition (5.20c) yields

$$T_\beta^\gamma \mathbf{w}^{\text{nc}} = G + \mathcal{R}_{-1} W^{\text{nc}} \quad \text{if } x_2 = 0.$$

From the proof of Lemma 5.4, we find that $\zeta_1 \neq 0$, and

$$(\zeta_1 \zeta_4 - \zeta_2 \zeta_3)|_{x_2=0} = \Delta = (\tau - i\zeta_0 \eta) h_0(t, x_1, \tau, \eta), \quad h_0(t, x_1, \tau, \eta) \neq 0$$

in a neighborhood of $(i\dot{z}_0\eta, \eta) \in \Xi_1$. According to identity (3.53), we define the following invertible matrices in a suitably small neighborhood of \mathcal{Y}_c^0 :

$$P_1 = \begin{pmatrix} 1/\dot{\zeta}_1 & 0 \\ -\dot{\zeta}_3/(\dot{\zeta}_1\dot{\zeta}_5) & 1/\dot{\zeta}_5 \end{pmatrix}, \quad P_2 = \begin{pmatrix} 1 & -\dot{\zeta}_2 \\ 0 & \dot{\zeta}_1 \end{pmatrix},$$

with $\zeta_5 := h_0(t, x_1, \tau, \eta)$ such that P_1 and P_2 belong to Γ_2^0 . Shrinking \mathcal{Y}_c^0 if necessary, we have

$$\beta_{\text{in}} := P_1 \beta(E_+ \ E_-) P_2 = \begin{pmatrix} 1 & 0 \\ 0 & \lambda^{-1, \gamma}(\gamma + i\tilde{\sigma}_0) \end{pmatrix} \quad \text{in } \mathcal{Y}_c^0, \quad (5.118)$$

where $\tilde{\sigma}_0 = \delta - \dot{z}_0\eta$ is the scalar real symbol in Γ_2^1 . We recall from (5.59) that $\sigma_0 = -i\gamma + \tilde{\sigma}_0$.

We then fix the four cut-off functions χ_{c_1} , χ_{c_2} , χ_{c_3} , and χ_{c_4} such that

- $\chi_{c_1} \equiv 1$ in a neighborhood of $\text{supp } \chi_1 \cap \{x_2 = 0\}$;
- $\chi_{c_j} \equiv 1$ in a neighborhood of $\text{supp } \chi_{c_{j-1}}$ for $j = 2, 3, 4$;
- $\text{supp } \chi_{c_4} \subset \mathcal{Y}_c^0 \cap \{x_2 = 0\}$.

As in [15, § 3.4.3], the following estimate can be obtained by using the localized Gårding's inequality:

$$\begin{aligned} & \|T_{\chi_{c_2} \lambda^{1, \gamma} \beta_{\text{in}}}^\gamma T_{\chi_{c_1}}^\gamma T_{\chi_{c_4} P_2^{-1}}^\gamma Z_1|_{x_2=0}\| \\ & \lesssim \|G\|_{1, \gamma} + \|Z_2|_{x_2=0}\|_{1, \gamma} + \|W^{\text{nc}}|_{x_2=0}\|. \end{aligned} \quad (5.119)$$

Now we utilize the special structure of β_{in} to derive a lower bound for the term on the left-hand side of (5.119). Setting

$$(v_1, v_2)^\top := T_{\chi_{c_4} P_2^{-1}}^\gamma Z_1|_{x_2=0}, \quad (5.120)$$

we obtain from (5.118) that

$$\begin{aligned} & \|T_{\chi_{c_2} \lambda^{1, \gamma} \beta_{\text{in}}}^\gamma T_{\chi_{c_1}}^\gamma T_{\chi_{c_4} P_2^{-1}}^\gamma Z_1|_{x_2=0}\|^2 \\ & = \|T_{\chi_{c_2} \lambda^{1, \gamma}}^\gamma T_{\chi_{c_1}}^\gamma v_1\|^2 + \|T_{\chi_{c_2} (\gamma + i\tilde{\sigma}_0)}^\gamma T_{\chi_{c_1}}^\gamma v_2\|^2. \end{aligned} \quad (5.121)$$

Use Lemma 5.2 (iv) and apply the localized Gårding's inequality (Lemma 5.2 (vii)) to obtain

$$\begin{aligned} & \|T_{\chi_{c_2} \lambda^{1, \gamma}}^\gamma T_{\chi_{c_1}}^\gamma v_1\|^2 = \langle (T_{\chi_{c_2} \lambda^{1, \gamma}}^\gamma)^* T_{\chi_{c_2} \lambda^{1, \gamma}}^\gamma T_{\chi_{c_1}}^\gamma v_1, T_{\chi_{c_1}}^\gamma v_1 \rangle \\ & \geq \text{Re} \langle T_{\chi_{c_2} \lambda^{2, \gamma}}^\gamma T_{\chi_{c_1}}^\gamma v_1, T_{\chi_{c_1}}^\gamma v_1 \rangle - C \|T_{\chi_{c_1}}^\gamma v_1\| \|T_{\chi_{c_1}}^\gamma v_1\|_{1, \gamma} \\ & \geq c \|T_{\chi_{c_1}}^\gamma v_1\|_{1, \gamma}^2 - C \|v_1\|^2 - C \|T_{\chi_{c_1}}^\gamma v_1\|^2 \\ & \gtrsim \|v_1\|_{1, \gamma}^2 - C \|Z_1|_{x_2=0}\|^2 \gtrsim \gamma^2 \|v_1\|^2 + \|T_{\tilde{\sigma}_0}^\gamma v_1\|^2 - C \|Z_1|_{x_2=0}\|^2 \end{aligned}$$

for large enough γ . Similarly, we obtain that, for sufficiently large γ ,

$$\|T_{\chi_{c_2} (\gamma + i\tilde{\sigma}_0)}^\gamma T_{\chi_{c_1}}^\gamma v_2\|^2 \gtrsim \gamma^2 \|v_2\|^2 + \|T_{\tilde{\sigma}_0}^\gamma v_2\|^2 - C \|Z_1|_{x_2=0}\|^2.$$

Plug the above two estimates into (5.121) to infer

$$\begin{aligned} & \|T_{\chi_{c_2}}^\gamma \lambda^{1,\gamma} \beta_{\text{in}} T_{\chi_{c_1}}^\gamma T_{\chi_{c_4} P_2^{-1}}^\gamma Z_1|_{x_2=0}\|^2 \\ & \gtrsim \gamma^2 \|(v_1, v_2)\|^2 + \|T_{\tilde{\sigma}_0}^\gamma(v_1, v_2)\|^2 - C\|Z_1|_{x_2=0}\|^2. \end{aligned} \quad (5.122)$$

Since $\chi_{c_3} \chi_1 \equiv \chi_1$, we see from (5.110) and (5.120) that

$$T_{\chi_{c_3}}^\gamma T_{\tilde{\sigma}_0}^\gamma Z_1 = T_{\tilde{\sigma}_0}^\gamma T_{\chi_{c_3}}^\gamma Z_1 + \mathcal{R}_0 Z_1 = T_{\tilde{\sigma}_0}^\gamma Z_1 + \mathcal{R}_0 Z_1,$$

so that

$$\begin{aligned} T_{\tilde{\sigma}_0}^\gamma(v_1, v_2) &= T_{\chi_{c_4} P_2^{-1}}^\gamma T_{\tilde{\sigma}_0}^\gamma Z_1|_{x_2=0} + \mathcal{R}_0 Z_1|_{x_2=0} \\ &= T_{\chi_{c_4} P_2^{-1}}^\gamma T_{\chi_{c_3}}^\gamma T_{\tilde{\sigma}_0}^\gamma Z_1|_{x_2=0} + \mathcal{R}_0 Z_1|_{x_2=0}. \end{aligned}$$

Thanks to the ellipticity of $(P_2^{-1})^* P_2^{-1}$ on the support of χ_{c_4} , we apply the localized Gårding's inequality to obtain

$$\begin{aligned} & \|T_{\tilde{\sigma}_0}^\gamma(v_1, v_2)\|^2 \\ & \geq \langle (T_{\chi_{c_4} P_2^{-1}}^\gamma)^* T_{\chi_{c_4} P_2^{-1}}^\gamma T_{\chi_{c_3}}^\gamma T_{\tilde{\sigma}_0}^\gamma Z_1|_{x_2=0}, T_{\chi_{c_3}}^\gamma T_{\tilde{\sigma}_0}^\gamma Z_1|_{x_2=0} \rangle - C\|Z_1|_{x_2=0}\|^2 \\ & \gtrsim \|T_{\chi_{c_3}}^\gamma T_{\tilde{\sigma}_0}^\gamma Z_1|_{x_2=0}\|^2 - C\|T_{\tilde{\sigma}_0}^\gamma Z_1|_{x_2=0}\|_{-1,\gamma}^2 \\ & \quad - C\|T_{\chi_{c_3}}^\gamma T_{\tilde{\sigma}_0}^\gamma Z_1|_{x_2=0}\|_{-1,\gamma}^2 - C\|Z_1|_{x_2=0}\|^2 \end{aligned}$$

for large enough γ . Then we take γ sufficiently large to deduce

$$\|T_{\tilde{\sigma}_0}^\gamma(v_1, v_2)\|^2 \gtrsim \|T_{\tilde{\sigma}_0}^\gamma Z_1|_{x_2=0}\|^2 - C\|Z_1|_{x_2=0}\|^2. \quad (5.123)$$

Similarly, we have

$$\begin{aligned} \|(v_1, v_2)\|^2 & \gtrsim \|Z_1|_{x_2=0}\|^2 - C\|Z_1|_{x_2=0}\|_{-1,\gamma}^2 \\ & \gtrsim \|Z_1|_{x_2=0}\|^2 - \frac{C}{\gamma^2} \|Z_1|_{x_2=0}\|^2. \end{aligned} \quad (5.124)$$

Combining estimates (5.117), (5.119), and (5.122)–(5.124), we take γ large enough to derive (5.113) and conclude the proof. \square

Recall that vectors Z^\pm are defined by (5.110) and that matrices $Q_0^\pm \in \Gamma_2^0$ are invertible in a neighborhood of the support of χ_1 and $Q_{-1}^\pm \in \Gamma_1^{-1}$. It then follows from Lemma 5.10 that

$$\begin{aligned} & \gamma^3 \|T_{\chi_c^0}^\gamma W^{\text{nc}}\|^2 + \gamma^2 \|T_{\chi_c^0}^\gamma W^{\text{nc}}|_{x_2=0}\|^2 + \|T_{\tilde{\sigma}_0}^\gamma T_{\chi_c^0}^\gamma W^{\text{nc}}|_{x_2=0}\|^2 \\ & \lesssim \gamma^{-1} (\|T_r^\gamma W\|_{1,\gamma}^2 + \|T_{\chi_c^0}^\gamma F\|_{1,\gamma}^2 + \|W\|^2) + \|G\|_{1,\gamma}^2 + \|W^{\text{nc}}|_{x_2=0}\|^2. \end{aligned} \quad (5.125)$$

Noting that components $T_{\chi_c^0}^\gamma W_1^\pm$ are given in terms of $T_{\chi_c^0}^\gamma W_2^\pm$ and $T_{\chi_c^0}^\gamma W_3^\pm$ by relation (5.31), we can deduce an L^2 estimate for $T_{\chi_c^0}^\gamma W_1^\pm$, that is, we can add the terms, $\gamma^3 \|T_{\chi_c^0}^\gamma W_1^\pm\|^2$, on the left-hand side of (5.125).

The following proposition then follows by combining the estimates for the three cases \mathcal{T}_c^q with $q \in \{0, \pm 1\}$.

Proposition 5.5 *There exist constants $K_0 \leq 1$ and $\gamma_0 \geq 1$ such that, if $\gamma \geq \gamma_0$ and $K \leq K_0$ for K given in (4.2), then*

$$\begin{aligned} & \gamma^3 \|T_{\chi_c}^\gamma W\|^2 + \gamma^2 \|T_{\chi_c}^\gamma W^{\text{nc}}|_{x_2=0}\|^2 + \sum_{q \in \{0, \pm 1\}} \|T_{\tilde{\sigma}_q}^\gamma T_{\chi_c}^\gamma W^{\text{nc}}|_{x_2=0}\|^2 \\ & \lesssim \gamma^{-1} (\|T_r^\gamma W\|_{1,\gamma}^2 + \|T_{\chi_c}^\gamma F\|_{1,\gamma}^2 + \|W\|^2) + \|G\|_{1,\gamma}^2 + \|W^{\text{nc}}|_{x_2=0}\|^2, \end{aligned} \quad (5.126)$$

where $\tilde{\sigma}_q \in \Gamma_2^1$ is given in (5.59) and $r \in \Gamma_1^0$ vanishes in region $\{\chi_c \equiv 1\} \cup \{\chi_c \equiv 0\}$.

5.7 Proof of Theorem 5.1

We now patch the microlocalized energy estimates (5.80), (5.106), and (5.126) together to deduce estimate (5.21). Since $\chi_u + \chi_p + \chi_c \equiv 1$,

$$\begin{aligned} & \gamma^3 \|W\|^2 + \gamma^2 \|W^{\text{nc}}|_{x_2=0}\|^2 \\ & \lesssim \gamma \| (T_{\chi_u}^\gamma W, T_{\chi_p}^\gamma W) \|_{1,\gamma}^2 + \gamma^3 \|T_{\chi_c}^\gamma W\|^2 \\ & \quad + \| (T_{\chi_u}^\gamma W^{\text{nc}}, T_{\chi_p}^\gamma W^{\text{nc}}) |_{x_2=0} \|_{1,\gamma}^2 + \gamma^2 \|T_{\chi_c}^\gamma W^{\text{nc}}|_{x_2=0}\|^2. \end{aligned} \quad (5.127)$$

Adding estimates (5.80), (5.106), and (5.126), we use (5.127) and take γ large enough to deduce

$$\begin{aligned} & \gamma \| (T_{\chi_u}^\gamma W, T_{\chi_p}^\gamma W) \|_{1,\gamma}^2 + \| (T_{\chi_u}^\gamma W^{\text{nc}}, T_{\chi_p}^\gamma W^{\text{nc}}) |_{x_2=0} \|_{1,\gamma}^2 \\ & \quad + \gamma^3 \|T_{\chi_c}^\gamma W\|^2 + \gamma^2 \|T_{\chi_c}^\gamma W^{\text{nc}}|_{x_2=0}\|^2 \\ & \lesssim \|G\|_{1,\gamma}^2 + \gamma^{-1} \|F\|_{1,\gamma}^2 + \gamma^{-1} \|T_r^\gamma W\|_{1,\gamma}^2. \end{aligned} \quad (5.128)$$

In order to absorb the *microlocalization error* term $\|T_r^\gamma W\|_{1,\gamma}$, we decompose symbol r in terms of χ_u , χ_p , and σ_q ($q \in \{0, \pm 1\}$). Notice that symbol $r \in \Gamma_1^0$ vanishes in the region:

$$\{\chi_c \equiv 1\} \cap \{\chi_p \equiv 0\} \cap \{\chi_u \equiv 0\} \equiv \{\chi_c \equiv 1\}.$$

In region $\{\chi_c \leq 1/2\}$, $\chi_u + \chi_p \geq \frac{1}{2}$, so that we can write

$$r = \alpha_u \chi_u + \alpha_p \chi_p,$$

where matrices α_u and α_p belong to Γ_1^0 and have the same block diagonal structure as \mathbb{A} . In region $\{\frac{1}{2} \leq \chi_c < 1\}$, we can utilize (5.61) to write

$$r = \sum_{q \in \{0, \pm 1\}} \alpha_c^q \begin{pmatrix} \sigma_+^q I_3 & 0 \\ 0 & \sigma_-^q I_3 \end{pmatrix} \chi_c^q,$$

where $\alpha_c^q \in \Gamma_1^{-1}$ has the same block diagonal structure as \mathbb{A} , and σ_\pm^q are solutions to (5.60). Thus we obtain

$$\|T_r^\gamma W\|_{1,\gamma} \lesssim \| (T_{\chi_u}^\gamma W, T_{\chi_p}^\gamma W) \|_{1,\gamma} + \sum_{q \in \{0, \pm 1\}} \|T_{\sigma_\pm^q}^\gamma T_{\chi_c^q}^\gamma W\|. \quad (5.129)$$

We now make the estimate for the last term in (5.129) in the following lemma:

Lemma 5.11 *There exist constants $K_0 \leq 1$ and $\gamma_0 \geq 1$ such that, if $\gamma \geq \gamma_0$ and $K \leq K_0$ for K given in (4.2), then, for $q \in \{0, \pm 1\}$,*

$$\begin{aligned} \gamma \|T_{\sigma_{\pm}^q}^{\gamma} T_{\chi_c^q}^{\gamma} W\|^2 &\lesssim \gamma \|W\|^2 + \gamma^{-1} (\|T_r^{\gamma} W\|_{1,\gamma}^2 + \|T_{\chi_c^q}^{\gamma} F\|_{1,\gamma}^2) \\ &\quad + \|G\|_{1,\gamma}^2 + \gamma^2 \|W^{\text{nc}}|_{x_2=0}\|^2, \end{aligned} \quad (5.130)$$

where $r \in \Gamma_1^0$ vanishes in region $\{\chi_c^q \equiv 1\}$.

Proof Let us show an estimate for $T_{\sigma_+^0}^{\gamma} Z_1^+$ with Z_1^+ defined by (5.110). Recall from (5.60) that symbol σ_+^0 satisfies the transport equation:

$$\begin{cases} \partial_2 \sigma_+^0 + \{\sigma_+^0, \text{Im} \omega_+\} = 0 & \text{if } x_2 > 0, \\ \sigma_+^0 = -i\gamma + \tilde{\sigma}_0 & \text{if } x_2 = 0. \end{cases} \quad (5.131)$$

Setting $S := (T_{\sigma_+^0}^{\gamma})^* T_{\sigma_+^0}^{\gamma}$, we take the scalar product in $L^2(\Omega)$ of (5.114) with SZ_1^+ and apply integration by parts to derive

$$\|(T_{\sigma_+^0}^{\gamma} Z_1^+)|_{x_2=0}\|^2 + \sum_{j=1}^6 \mathcal{I}_j = 0, \quad (5.132)$$

where each term \mathcal{I}_j in the decomposition will be defined and estimated below. First, noting that $\sigma_+^0 \in \Gamma_{\frac{1}{2}}^1$ and $T_{\sigma_+^0}^{\gamma} T_{d_+}^{\gamma} = T_{d_+}^{\gamma} T_{\sigma_+^0}^{\gamma} + \mathcal{R}_0$, we obtain

$$\begin{aligned} \mathcal{I}_1 &:= 2 \text{Re} \langle T_{\sigma_+^0}^{\gamma} T_{d_+}^{\gamma} Z_1^+, T_{\sigma_+^0}^{\gamma} Z_1^+ \rangle \\ &\lesssim \|T_{\sigma_+^0}^{\gamma} Z_1^+\|^2 + \|Z_1^+\| \|T_{\sigma_+^0}^{\gamma} Z_1^+\| \lesssim \|T_{\sigma_+^0}^{\gamma} Z_1^+\|^2 + \|Z_1^+\|^2. \end{aligned} \quad (5.133)$$

Moreover, we have

$$\begin{aligned} \mathcal{I}_2 &:= 2 \text{Re} \langle T_{\sigma_+^0}^{\gamma} T_r^{\gamma} W^+, T_{\sigma_+^0}^{\gamma} Z_1^+ \rangle \\ &\lesssim (\varepsilon\gamma)^{-1} \|T_r^{\gamma} W\|_{1,\gamma}^2 + \varepsilon\gamma \|T_{\sigma_+^0}^{\gamma} Z_1^+\|^2, \end{aligned} \quad (5.134)$$

$$\begin{aligned} \mathcal{I}_3 &:= 2 \text{Re} \langle T_{\sigma_+^0}^{\gamma} \mathcal{R}_0 T_{\chi_c^0}^{\gamma} F^+, T_{\sigma_+^0}^{\gamma} Z_1^+ \rangle \\ &\lesssim (\varepsilon\gamma)^{-1} \|T_{\chi_c^0}^{\gamma} F^+\|_{1,\gamma}^2 + \varepsilon\gamma \|T_{\sigma_+^0}^{\gamma} Z_1^+\|^2, \end{aligned} \quad (5.135)$$

$$\begin{aligned} \mathcal{I}_4 &:= 2 \text{Re} \langle T_{\sigma_+^0}^{\gamma} \mathcal{R}_{-1} W^+, T_{\sigma_+^0}^{\gamma} Z_1^+ \rangle \\ &\lesssim (\varepsilon\gamma)^{-1} \|W^+\|^2 + \varepsilon\gamma \|T_{\sigma_+^0}^{\gamma} Z_1^+\|^2. \end{aligned} \quad (5.136)$$

For the terms:

$$\mathcal{I}_5 := \text{Re} \langle (\partial_2 S) Z_1^+, Z_1^+ \rangle, \quad \mathcal{I}_6 := 2 \text{Re} \langle S T_{\omega_+}^{\gamma} Z_1^+, Z_1^+ \rangle,$$

we use the identity: $\partial_2 S = (T_{\partial_2 \sigma_+^0}^{\gamma})^* T_{\sigma_+^0}^{\gamma} + (T_{\sigma_+^0}^{\gamma})^* T_{\partial_2 \sigma_+^0}^{\gamma}$ to obtain

$$\mathcal{I}_5 + \mathcal{I}_6 = 2 \text{Re} \langle T_{\partial_2 \sigma_+^0}^{\gamma} Z_1^+ + T_{\sigma_+^0}^{\gamma} T_{\omega_+}^{\gamma} Z_1^+, T_{\sigma_+^0}^{\gamma} Z_1^+ \rangle. \quad (5.137)$$

We write $\omega_+ = \gamma e_+ + i\hbar_+$ with real-valued symbols $e_+ \in \Gamma_2^0$ and $\hbar_+ \in \Gamma_2^1$ as in (5.112). Employing the rule of symbolic calculus (Lemma 5.2), we have

$$\begin{aligned} T_{\sigma_+^0}^\gamma T_{\omega_+}^\gamma &= \gamma T_{e_+}^\gamma T_{\sigma_+^0}^\gamma + \gamma T_{-i\{\sigma_+^0, e_+\}}^\gamma \\ &\quad + \gamma \mathcal{R}_{-1} + T_{i\hbar_+}^\gamma T_{\sigma_+^0}^\gamma + T_{\{\sigma_+^0, \hbar_+\}}^\gamma + \mathcal{R}_0. \end{aligned}$$

It follows from (5.131) and (5.137) that

$$\begin{aligned} \mathcal{I}_5 + \mathcal{I}_6 &\lesssim \operatorname{Re} \langle\langle (\gamma T_{e_+}^\gamma T_{\sigma_+^0}^\gamma + \gamma T_{-i\{\sigma_+^0, e_+\}}^\gamma + T_{i\hbar_+}^\gamma T_{\sigma_+^0}^\gamma) Z_1^+, T_{\sigma_+^0}^\gamma Z_1^+ \rangle\rangle \\ &\quad + \|\| Z_1^+ \|\| \| T_{\sigma_+^0}^\gamma Z_1^+ \|\|. \end{aligned}$$

Since $\sigma_+^0 \in \Gamma_2^1$ and $i\hbar_+ \in i\mathbb{R}$, operators $T_{-i\{\sigma_+^0, e_+\}}^\gamma$ and $\operatorname{Re} T_{i\hbar_+}^\gamma$ are both of order ≤ 0 . It then follows that

$$\begin{aligned} &\operatorname{Re} \langle\langle (\gamma T_{-i\{\sigma_+^0, e_+\}}^\gamma + T_{i\hbar_+}^\gamma T_{\sigma_+^0}^\gamma) Z_1^+, T_{\sigma_+^0}^\gamma Z_1^+ \rangle\rangle \\ &\lesssim \gamma \|\| Z_1^+ \|\| \| T_{\sigma_+^0}^\gamma Z_1^+ \|\| + \| T_{\sigma_+^0}^\gamma Z_1^+ \|^2 \lesssim \varepsilon^{-1} \gamma \|\| Z_1^+ \|\|^2 + \varepsilon \gamma \| T_{\sigma_+^0}^\gamma Z_1^+ \|^2, \end{aligned}$$

which implies

$$\mathcal{I}_5 + \mathcal{I}_6 \lesssim \operatorname{Re} \langle\langle \gamma T_{e_+}^\gamma T_{\sigma_+^0}^\gamma Z_1^+, T_{\sigma_+^0}^\gamma Z_1^+ \rangle\rangle + \varepsilon^{-1} \gamma \|\| Z_1^+ \|\|^2 + \varepsilon \gamma \| T_{\sigma_+^0}^\gamma Z_1^+ \|^2. \quad (5.138)$$

Since $e_+ \leq -c < 0$, we apply Gårding's inequality to deduce

$$-\operatorname{Re} \langle\langle \gamma T_{e_+}^\gamma T_{\sigma_+^0}^\gamma Z_1^+, T_{\sigma_+^0}^\gamma Z_1^+ \rangle\rangle \gtrsim \gamma \| T_{\sigma_+^0}^\gamma Z_1^+ \|^2 \quad (5.139)$$

for sufficiently large γ . Plugging estimates (5.133)–(5.136) and (5.138)–(5.139) into (5.132), we take ε small enough to deduce

$$\begin{aligned} \gamma \| T_{\sigma_+^0}^\gamma Z_1^+ \|^2 &\lesssim \| (T_{\sigma_+^0}^\gamma Z_1^+) |_{x_2=0} \|^2 \\ &\quad + \gamma^{-1} \| T_r^\gamma W \|_{1,\gamma}^2 + \gamma^{-1} \| T_{\chi_e}^\gamma F^+ \|_{1,\gamma}^2 + \gamma \| W^+ \|^2. \end{aligned} \quad (5.140)$$

For the first term on the right-hand side, we use the fact that $\sigma_+^0|_{x_2=0} = -i\gamma + \tilde{\sigma}_0$ to obtain

$$\| (T_{\sigma_+^0}^\gamma Z_1^+) |_{x_2=0} \| \lesssim \gamma \| Z_1^+ |_{x_2=0} \| + \| T_{\tilde{\sigma}_0}^\gamma Z_1^+ |_{x_2=0} \|. \quad (5.141)$$

We plug (5.141) into (5.140) and use (5.113) to find

$$\begin{aligned} \gamma \| T_{\sigma_+^0}^\gamma Z_1^+ \|^2 &\lesssim \gamma^{-1} (\| T_r^\gamma W \|_{1,\gamma}^2 + \| T_{\chi_e}^\gamma F^+ \|_{1,\gamma}^2) \\ &\quad + \gamma \| W^+ \|^2 + \| G \|_{1,\gamma}^2 + \gamma^2 \| W^{\text{nc}} |_{x_2=0} \|^2. \end{aligned} \quad (5.142)$$

Recall that Z_1^+ is defined by (5.110) with $Q_0 \in \Gamma_2^0$ being invertible. We then use (5.116) and (5.31) to conclude (5.130). The proof is completed. \square

Combining estimates (5.128)–(5.129) together, using (5.127), and taking γ suitably large, we obtain (5.21). In view of Proposition 5.2, estimate (5.1) also holds. This completes the proof of Theorem 5.1.

6 Well-posedness for the Linearized Problem

In this section, we establish a well-posedness result for the linearized problem (4.12) in the usual Sobolev space H^s with s large enough. The essential point is to deduce a tame estimate in H^s . For a hyperbolic problem with a characteristic boundary, there is a loss of derivatives in *a priori* energy estimates. To overcome this difficulty, it is natural to introduce Sobolev spaces with conormal regularity, where two tangential derivatives count as one normal derivative (see Secchi [43] and the references therein). However, for our problem (4.12), we can manage to compensate the loss of derivatives and deduce *a priori* estimates in the usual Sobolev spaces. This is achieved by employing the idea in [19] and estimating the missing derivatives through the equation of the linearized vorticity.

The main result in this section is stated as follows:

Theorem 6.1 *Let $T > 0$ and $s \in [3, \tilde{\alpha}] \cap \mathbb{N}$ with any integer $\tilde{\alpha} \geq 3$. Assume that the background state (2.24) satisfies (2.10) and (2.25), and that perturbations $(\dot{V}^\pm, \dot{\Psi}^\pm)$ belong to $H_\gamma^{s+3}(\Omega_T)$ for all $\gamma \geq 1$ and satisfy (4.1)–(4.3), and*

$$\|(\dot{V}^\pm, \nabla \dot{\Psi}^\pm)\|_{H_\gamma^s(\Omega_T)} + \|(\dot{V}^\pm, \partial_2 \dot{V}^\pm, \nabla \dot{\Psi}^\pm)|_{x_2=0}\|_{H_\gamma^s(\omega_T)} \leq K. \quad (6.1)$$

Assume further that the source terms $(f, g) \in H^{s+1}(\Omega_T) \times H^{s+1}(\omega_T)$ vanish in the past. Then there exists a positive constant K_0 , which is independent of s and T , and there exist two constants $C > 0$ and $\gamma \geq 1$, which depend solely on K_0 , such that, if $K \leq K_0$, then problem (4.12) admits a unique solution $(\dot{V}^\pm, \psi) \in H^s(\Omega_T) \times H^{s+1}(\omega_T)$ that vanishes in the past and obeys the following tame estimate:

$$\begin{aligned} & \|\dot{V}\|_{H_\gamma^s(\Omega_T)} + \|\mathbb{P}^\pm(\dot{\varphi})\dot{V}^\pm|_{x_2=0}\|_{H_\gamma^s(\omega_T)} + \|\psi\|_{H_\gamma^{s+1}(\omega_T)} \\ & \leq C \{ \|f\|_{H_\gamma^{s+1}(\Omega_T)} + \|g\|_{H_\gamma^{s+1}(\omega_T)} \\ & \quad + (\|f\|_{H_\gamma^4(\Omega_T)} + \|g\|_{H_\gamma^4(\omega_T)}) \|(\dot{V}^\pm, \dot{\Psi}^\pm)\|_{H_\gamma^{s+3}(\Omega_T)} \}. \end{aligned} \quad (6.2)$$

We consider the case where the source terms f and g vanish in the past, which corresponds to the case with zero initial data. The case of general initial data is postponed to the nonlinear analysis which involves the construction of a so-called approximate solution. Before estimating the higher order derivatives of solutions, we first prove that the linearized problem (4.12) is well-posed in L^2 .

6.1 First Well-Posedness Result

In this subsection, we apply the well-posedness result in L^2 of Coulombel [16] to the effective linear problem (4.12). We recall that system (4.12a) is symmetrizable hyperbolic and observe that the coefficients of the linearized

operators satisfy the regularity assumptions of [16]. We also recall that problem (4.12) satisfies the energy estimate (5.1), which exhibits a loss of one tangential derivative. According to the result in [16], we only need to find a dual problem that obeys an appropriate energy estimate.

Let us define a dual problem for (4.12). We introduce the following matrices:

$$\mathring{B}_1 := \left(\begin{array}{cccccc} 0 & 0 & 0 & -\mathring{\ell}_1^- & 0 & 0 \\ \mathring{\ell}_1^+ & 0 & 0 & \mathring{\ell}_1^- & 0 & 0 \\ 0 & \mathring{\ell}_2^+ & \mathring{\ell}_3^+ & 0 & -\mathring{\ell}_2^- & -\mathring{\ell}_3^- \end{array} \right) \Big|_{x_2=0}, \quad \mathring{D}_1 := \left(\begin{array}{cccccc} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \mathring{\ell}_2^+ & \mathring{\ell}_3^+ & 0 & \mathring{\ell}_2^- & \mathring{\ell}_3^- \end{array} \right) \Big|_{x_2=0}, \quad (6.3)$$

and

$$\mathring{D} := \left(\begin{array}{cccccc} 0 & \mathring{r}_+^{-1} \mathring{\varrho}_+ & -\mathring{r}_+^{-1} \mathring{\varsigma}_+ & 0 & \mathring{r}_-^{-1} \mathring{\varrho}_- & -\mathring{r}_-^{-1} \mathring{\varsigma}_- \\ 0 & \mathring{r}_+^{-1} \mathring{\varrho}_+ & -\mathring{r}_+^{-1} \mathring{\varsigma}_+ & 0 & 0 & 0 \\ 1 & 0 & 0 & 1 & 0 & 0 \end{array} \right) \Big|_{x_2=0},$$

where

$$\mathring{\ell}_1^\pm := -\frac{\mathring{N}^\pm \mathring{c}_\pm^2 \mathring{r}_\pm}{\partial_2 \mathring{\Phi}^\pm}, \quad \mathring{\ell}_2^\pm := -\frac{\mathring{\varrho}_\pm}{2\mathring{N}^\pm \partial_2 \mathring{\Phi}^\pm}, \quad \mathring{\ell}_3^\pm := \frac{\mathring{\varsigma}_\pm}{2\mathring{N}^\pm \partial_2 \mathring{\Phi}^\pm}.$$

Thanks to (4.15), we compute that these matrices satisfy the relation:

$$\mathring{B}_1^\top \mathring{B} + \mathring{D}_1^\top \mathring{D} = \text{diag}(\tilde{A}_2(\mathring{U}^+, \mathring{\Phi}^+), \tilde{A}_2(\mathring{U}^-, \mathring{\Phi}^-)) \Big|_{x_2=0}, \quad (6.4)$$

where \mathring{B} is defined by (4.8). Moreover, we infer from (4.2) that all matrices \mathring{B} , \mathring{B}_1 , \mathring{D} , and \mathring{D}_1 belong to $W^{2,\infty}(\mathbb{R}^2)$. Following [34, § 3.2], we define a dual problem for (4.12) as:

$$\begin{cases} \mathbb{L}'_e(\mathring{U}^\pm, \mathring{\Phi}^\pm)^* U^\pm = \tilde{f}_\pm, & x_2 > 0, \\ \mathring{D}_1 U = 0, & x_2 = 0, \\ \text{div}(b^\top \mathring{B}_1 U) - b_\sharp^\top \mathring{B}_1 U = 0, & x_2 = 0, \end{cases} \quad (6.5)$$

where b , b_\sharp , \mathring{B}_1 , and \mathring{D}_1 are defined in (4.7), (4.13), and (6.3), div denotes the divergence operator in \mathbb{R}^2 with respect to (t, x_1) , and the dual operators $\mathbb{L}'_e(\mathring{U}^\pm, \mathring{\Phi}^\pm)^*$ are the formal adjoints of $\mathbb{L}'_e(\mathring{U}^\pm, \mathring{\Phi}^\pm)$. More precisely, we have

$$\begin{aligned} \mathbb{L}'_e(V, \Psi)^* U &= -A_0(V)^\top \partial_t U - A_1(V)^\top \partial_1 U - \tilde{A}_2(V, \Psi)^\top \partial_2 U \\ &\quad + \mathcal{C}(V, \Psi)^* U - (\partial_t A_0(V)^\top + \partial_1 A_1(V)^\top + \partial_2 \tilde{A}_2(V, \Psi)^\top) U, \end{aligned}$$

where $\mathcal{C}(V, \Psi)^*$, the adjoint of $\mathcal{C}(V, \Psi)$, is a zero-th order operator. We refer to [34, § 3.2] for the derivation of the dual problem by using integration by parts and identity (6.4).

Since the first two rows of matrix \mathring{D}_1 given in (6.3) are zero, we see that the number of the boundary conditions in (6.5) is exactly two. This is compatible with the number of incoming characteristics, that is, the number of negative eigenvalues of the boundary matrix for (6.5). In fact, the boundary matrix

of operator $\mathbb{L}'_e(V, \Psi)^*$ in the half-space Ω is $\tilde{A}_2(V, \Psi)^\top$. Then we infer from (4.15) that problem (6.5) has two incoming characteristics and two outgoing characteristics.

We can define and analyze the Lopatinskiĭ determinant associated with the boundary conditions in (6.5) as we have done in §5. Then we have the following result, which is an analogue of Lemma 5.4 by changing γ into $-\gamma$.

Lemma 6.1 *Assume that (4.1)–(4.2) hold for a sufficiently small $K > 0$. Then the dual problem (6.5) satisfies the backward Lopatinskiĭ condition. Moreover, the roots of the associated Lopatinskiĭ determinant are simple and coincide with the roots of the Lopatinskiĭ determinant (5.50) for the original problem (4.12).*

One can reproduce the same analysis as we have done in §5 to show that the dual problem satisfies an *a priori* estimate that is similar to (5.1). The linearized problem (4.12) thus satisfies all the assumptions (*i.e.* symmetrizability, regularity, and weak stability) listed in [16]. We therefore obtain the following well-posedness result.

Theorem 6.2 *Let $T > 0$ be any fixed constant. Assume that the background state (2.24) satisfies (2.10) and (2.25). Assume further that the basic state $(\dot{V}^\pm, \dot{\Psi}^\pm)$ satisfies (4.1)–(4.3). Then there exist positive constants $K_0 > 0$ and $\gamma_0 \geq 1$, independent of T , such that, if $K \leq K_0$, then, for the source terms $f^\pm \in L^2(\mathbb{R}_+; H^1(\omega_T))$ and $g \in H^1(\omega_T)$ that vanish for $t < 0$, the problem:*

$$\begin{cases} \mathbb{L}'_e(\dot{U}^\pm, \dot{\Phi}^\pm)\dot{V}^\pm = f^\pm & \text{for } t < T, \ x_2 > 0, \\ \mathbb{B}'_e(\dot{U}^\pm, \dot{\Phi}^\pm)(\dot{V}, \psi) = g & \text{for } t < T, \ x_2 = 0, \end{cases}$$

has a unique solution $(\dot{V}^+, \dot{V}^-, \psi) \in L^2(\Omega_T) \times L^2(\Omega_T) \times H^1(\omega_T)$ that vanishes for $t < 0$ and satisfies $\mathbb{P}^\pm(\dot{\varphi})\dot{V}^\pm|_{x_2=0} \in L^2(\omega_T)$. Moreover, the following estimate holds for all $\gamma \geq \gamma_0$ and for all $t \in [0, T]$:

$$\begin{aligned} & \gamma \|\dot{V}\|_{L^2_\gamma(\Omega_t)}^2 + \|\mathbb{P}^\pm(\dot{\varphi})\dot{V}^\pm|_{x_2=0}\|_{L^2_\gamma(\omega_t)}^2 + \|\psi\|_{H^1_\gamma(\omega_t)}^2 \\ & \lesssim \gamma^{-3} \|f^\pm\|_{L^2(H^1_\gamma(\omega_t))}^2 + \gamma^{-2} \|g\|_{H^1_\gamma(\omega_t)}^2. \end{aligned} \quad (6.6)$$

Theorem 6.2 shows the well-posedness of problem (4.12) in L^2 when the source terms (f, g) belong to $L^2(H^1) \times H^1$. We now turn to the energy estimates for the higher-order derivatives of solutions.

6.2 *A Priori* Tame Estimates

To obtain the estimates for the higher-order derivatives of solutions to (4.12), it is convenient to deal with the reformulated problem (4.18) and (4.20) for the new unknowns W . Until the end of this section, we always assume that $\gamma \geq \gamma_0$

and $K \leq K_0$, where γ_0 and K_0 are given by Theorem 6.2. Then estimate (6.6) can be rewritten as

$$\begin{aligned} & \sqrt{\gamma} \|W\|_{L^2_\gamma(\Omega_T)} + \|W^{\text{nc}}|_{x_2=0}\|_{L^2_\gamma(\omega_T)} + \|\psi\|_{H^1_\gamma(\omega_T)} \\ & \lesssim \gamma^{-3/2} \|F^\pm\|_{L^2(H^1_\gamma(\omega_T))} + \gamma^{-1} \|g\|_{H^1_\gamma(\omega_T)}. \end{aligned} \quad (6.7)$$

We first derive the estimate of the tangential derivatives. Let $k \in [1, s]$ be a fixed integer. Applying the tangential derivative $\partial^\alpha = \partial_t^{\alpha_0} \partial_1^{\alpha_1}$ with $|\alpha| = k$ to system (4.18) yields the equations for $\partial^\alpha W^\pm$ that involve the linear terms of the derivatives, $\partial^{\alpha-\beta} \partial_t W^\pm$ and $\partial^{\alpha-\beta} \partial_1 W^\pm$, with $|\beta| = 1$. These terms cannot be treated simply as source terms, owing to the loss of derivatives in the energy estimate (6.7). To overcome this difficulty, we adopt the idea of [19] and deal with a boundary value problem for all the tangential derivatives of order equal to k , *i.e.* for $W^{(k)} := \{\partial_t^{\alpha_0} \partial_1^{\alpha_1} W^\pm, \alpha_0 + \alpha_1 = k\}$. Such a problem satisfies the same regularity and stability properties as the original problem (4.18) and (4.20). Repeating the derivation in § 5, we find that $W^{(k)}$ obeys an energy estimate similar to (6.7) with new source terms $\mathcal{F}^{(k)}$ and $\mathcal{G}^{(k)}$. Then we can employ the Gagliardo–Nirenberg and Moser-type inequalities (*cf.* [19, Theorems 8–10]) to derive the following estimate for tangential derivatives (see [19, Proposition 1] for the detailed proof).

Lemma 6.2 (Estimate of tangential derivatives) *Assume that the hypotheses of Theorem 6.1 hold. Then there exist constants $C_s > 0$ and $\gamma_s \geq 1$, independent of T , such that, for all $\gamma \geq \gamma_s$ and for all $(W, \psi) \in H^{s+2}_\gamma(\Omega_T) \times H^{s+2}_\gamma(\omega_T)$ that are solutions of problem (4.18) and (4.20), the following estimate holds:*

$$\begin{aligned} & \sqrt{\gamma} \|W\|_{L^2(H^s_\gamma(\omega_T))} + \|W^{\text{nc}}|_{x_2=0}\|_{H^s_\gamma(\omega_T)} + \|\psi\|_{H^{s+1}_\gamma(\omega_T)} \\ & \leq C_s \left\{ \gamma^{-1} \|g\|_{H^{s+1}_\gamma(\omega_T)} + \gamma^{-3/2} \|F^\pm\|_{L^2(H^{s+1}_\gamma(\omega_T))} \right. \\ & \quad + \gamma^{-1} \left(\|W^{\text{nc}}|_{x_2=0}\|_{L^\infty(\omega_T)} + \|\psi\|_{W^{1,\infty}(\omega_T)} \right) \|(\mathring{V}, \partial_2 \mathring{V}, \nabla \mathring{\Psi})|_{x_2=0}\|_{H^{s+1}_\gamma(\omega_T)} \\ & \quad \left. + \gamma^{-3/2} \|W\|_{W^{1,\infty}(\Omega_T)} \|(\mathring{V}, \nabla \mathring{\Psi})\|_{H^{s+2}_\gamma(\Omega_T)} \right\}. \end{aligned} \quad (6.8)$$

We recall that the boundary matrix for our problem (4.18) and (4.20) is not invertible. Thus, there is no hope to estimate all the normal derivatives of W directly from (4.18) by employing the standard argument for noncharacteristic boundary problems as in [41]. Nevertheless, for our problem (4.12), we can obtain the estimate of the missing normal derivatives through the equation of the “linearized vorticity”.

In view of the original equations (2.14), taking into account the change of variables and linearization, we define the “linearized vorticity” as:

$$\xi^\pm := \left(\partial_1 - \frac{\partial_1 \mathring{\Phi}^\pm}{\partial_2 \mathring{\Phi}^\pm} \partial_2 \right) \dot{V}_3^\pm - \frac{1}{\partial_2 \mathring{\Phi}^\pm} \partial_2 \dot{V}_2^\pm, \quad (6.9)$$

where \dot{V}_2^\pm and \dot{V}_3^\pm are the second and third components of the good unknown (4.10), respectively. We notice that multiplying (A.6) by $S_1(U)$ leads to system (2.18), where $S_1(U)$ is defined by (A.9). Thus, we multiply (4.12a) with matrices $S_1(\dot{U}^\pm)^{-1}$ to obtain

$$(B_0(\dot{U}^\pm)\partial_t + B_1(\dot{U}^\pm)\partial_1 + \tilde{B}_2(\dot{U}^\pm, \dot{\Phi}^\pm)\partial_2)\dot{V}^\pm + \tilde{C}(\dot{U}^\pm, \dot{\Phi}^\pm)\dot{V}^\pm = \tilde{f}^\pm, \quad (6.10)$$

where $\tilde{C}(\dot{U}^\pm, \dot{\Phi}^\pm) := S_1(\dot{U}^\pm)^{-1}C(\dot{U}^\pm, \dot{\Phi}^\pm)$, $\tilde{f}^\pm := S_1(\dot{U}^\pm)^{-1}f^\pm$, and matrices B_j are defined by (A.7)–(A.8). In light of (4.3a), we have

$$\begin{aligned} \tilde{B}_2(\dot{U}, \dot{\Phi}) &:= \frac{1}{\partial_2 \dot{\Phi}} (B_2(\dot{U}) - \partial_t \dot{\Phi} B_0(\dot{U}) - \partial_2 \dot{\Phi} B_1(\dot{U})) \\ &= \frac{1}{\partial_2 \dot{\Phi}} \begin{pmatrix} \tilde{B}_2^{11} & \tilde{B}_2^{12} & \tilde{B}_2^{13} \\ -\dot{N}^{-1} \partial_1 \dot{\Phi} & 0 & 0 \\ \dot{N}^{-1} & 0 & 0 \end{pmatrix}, \end{aligned}$$

where the explicit form of \tilde{B}_2^{1j} is of no interest. Using the second and third components of (6.10) yields the following equations for $\dot{\xi}^\pm$:

$$\begin{aligned} &(\partial_t + \dot{v}_1^\pm \partial_1)\dot{\xi}^\pm \\ &= \partial_1 \mathcal{F}_2^\pm - \frac{1}{\partial_2 \dot{\Phi}^\pm} (\partial_1 \dot{\Phi}^\pm \partial_2 \mathcal{F}_2^\pm + \partial_2 \mathcal{F}_1^\pm) + \dot{\Lambda}_1^\pm \cdot \partial_1 \dot{V}^\pm + \dot{\Lambda}_2^\pm \cdot \partial_2 \dot{V}^\pm, \end{aligned} \quad (6.11)$$

where vectors $\dot{\Lambda}_1^\pm$ and $\dot{\Lambda}_2^\pm$ are C^∞ -functions of $(\dot{V}^\pm, \nabla \dot{V}^\pm, \nabla \dot{\Psi}^\pm, \nabla^2 \dot{\Psi}^\pm)$ that vanish at the origin, and the source terms \mathcal{F}_1^\pm and \mathcal{F}_2^\pm are given by

$$\mathcal{F}_1^\pm := \dot{\Gamma}_\pm^{-1}(\tilde{f}^\pm - \tilde{C}(\dot{U}^\pm, \dot{\Phi}^\pm)\dot{V}^\pm)_2, \quad \mathcal{F}_2^\pm := \dot{\Gamma}_\pm^{-1}(\tilde{f}^\pm - \tilde{C}(\dot{U}^\pm, \dot{\Phi}^\pm)\dot{V}^\pm)_3.$$

Employing a standard energy estimate to the transport equations (6.11), we can apply the Gagliardo–Nirenberg and Moser-type inequalities to derive the following estimate of $\dot{\xi}^\pm$:

Lemma 6.3 (Estimate of vorticity) *Assume that the hypotheses of Theorem 6.1 hold. Then there exist constants $C_s > 0$ and $\gamma_s \geq 1$, independent of T , such that, for all $\gamma \geq \gamma_s$ and for all $(W, \psi) \in H_\gamma^{s+2}(\Omega_T) \times H_\gamma^{s+2}(\omega_T)$ that are solutions of problem (4.18) and (4.20), functions $\dot{\xi}^\pm$ defined by (6.9) satisfy the following estimate:*

$$\begin{aligned} &\gamma \|\dot{\xi}^\pm\|_{H_\gamma^{s-1}(\Omega_T)} \\ &\leq C_s \left\{ \|f^\pm\|_{H_\gamma^s(\Omega_T)} + \|f^\pm\|_{L^\infty(\Omega_T)} \|(\dot{V}, \nabla \dot{\Psi})\|_{H_\gamma^s(\Omega_T)} + \|\dot{V}^\pm\|_{H_\gamma^s(\Omega_T)} \right. \\ &\quad \left. + \|\dot{V}^\pm\|_{W^{1,\infty}(\Omega_T)} (\|\dot{V}\|_{H_\gamma^{s+1}(\Omega_T)} + \|\nabla \dot{\Psi}\|_{H_\gamma^s(\Omega_T)}) \right\}. \end{aligned} \quad (6.12)$$

We are going to make the estimate for all the normal derivatives by means of estimates (6.8) and (6.12) for the tangential derivatives and the linearized vorticity. To this end, we need to express the normal derivatives $\partial_2 W^\pm$ in terms of the tangential derivatives $\partial_t W^\pm$, $\partial_1 W^\pm$, and vorticity $\dot{\xi}^\pm$. Since $\mathbf{I}_2 =$

diag(0, 1, 1), the normal derivatives $\partial_2 W_2^\pm$ and $\partial_2 W_3^\pm$ are directly given by (4.18) as:

$$\mathbf{I}_2 \partial_2 W^\pm = \mathbf{I}_2 (F^\pm - \mathbf{A}_0^\pm \partial_t W^\pm - \mathbf{A}_1^\pm \partial_1 W^\pm - \mathbf{C}^\pm W^\pm). \quad (6.13)$$

The ‘‘missing’’ normal derivatives $\partial_2 W_1^\pm$ can be expressed by ξ^\pm and equations (4.18). From transformation (4.17) and definition (4.16), we have

$$\begin{aligned} \partial_2 \dot{V}_2 &= \dot{\varsigma} \partial_2 W_1 + \frac{\dot{\varrho}}{N \dot{c}} \partial_2 (W_2 - W_3) + \partial_2 \dot{\varsigma} W_1 + \partial_2 \left(\frac{\dot{\varrho}}{N \dot{c}} \right) (W_2 - W_3), \\ \partial_2 \dot{V}_3 &= \dot{\varrho} \partial_2 W_1 - \frac{\dot{\varsigma}}{N \dot{c}} \partial_2 (W_2 - W_3) + \partial_2 \dot{\varrho} W_1 - \partial_2 \left(\frac{\dot{\varsigma}}{N \dot{c}} \right) (W_2 - W_3), \end{aligned}$$

where we have omitted indices ‘‘ \pm ’’. By definition (6.9), we obtain

$$\begin{aligned} &(\partial_1 \dot{\Phi} \dot{\varrho} + \dot{\varsigma}) \partial_2 W_1 \\ &= \partial_2 \dot{\Phi} (\partial_1 \dot{V}_3 - \dot{\xi}) + \frac{\partial_1 \dot{\Phi} \dot{\varsigma} - \dot{\varrho}}{N \dot{c}} \partial_2 (W_2 - W_3) + C(\dot{U}, \dot{\Phi}) W, \end{aligned} \quad (6.14)$$

where $C(\dot{U}, \dot{\Phi})$ is a C^∞ -function of $(\dot{V}, \nabla \dot{V}, \nabla \dot{\Psi}, \nabla^2 \dot{\Psi})$ that vanishes at the origin. According to (4.9), we see that $\partial_1 \dot{\Phi} \dot{\varrho} + \dot{\varsigma} \gtrsim 1$ by taking $K_0 > 0$ sufficiently small. Then we find from (6.13)–(6.14) that

$$\begin{aligned} \partial_2 W^\pm &= \tilde{\mathbf{A}}_3^\pm F^\pm + \tilde{\mathbf{A}}_0^\pm \partial_t W^\pm \\ &\quad + \tilde{\mathbf{A}}_1^\pm \partial_1 W^\pm + \tilde{\mathbf{C}}^\pm W^\pm - \frac{\partial_2 \dot{\Phi}^\pm}{\partial_1 \dot{\Phi}^\pm \dot{\varrho}^\pm + \dot{\varsigma}^\pm} \begin{pmatrix} \xi^\pm \\ 0 \\ 0 \end{pmatrix}, \end{aligned} \quad (6.15)$$

where $\tilde{\mathbf{A}}_{0,1,3}^\pm$ are C^∞ -functions of $(\dot{V}, \nabla \dot{\Psi})$, and $\tilde{\mathbf{C}}^\pm$ are C^∞ -functions of $(\dot{V}, \nabla \dot{V}, \nabla \dot{\Psi}, \nabla^2 \dot{\Psi})$ that vanish at the origin. Although the linearized problem (4.18) and (4.20) is characteristic, we manage to express all the normal derivatives $\partial_2 W^\pm$ by (6.15) as a linear combination of the tangential derivatives, vorticity, zero-th order terms, and source terms. Then we can prove the following result similar to [19, Proposition 3], so we omit its proof.

Lemma 6.4 (Estimate of normal derivatives) *Assume that the hypotheses of Theorem 6.1 hold. Then there exist constants $C_s > 0$ and $\gamma_s \geq 1$, which are independent of T , such that, for all $\gamma \geq \gamma_s$ and solutions $(W, \psi) \in H_\gamma^{s+2}(\Omega_T) \times H_\gamma^{s+2}(\omega_T)$ of problem (4.18) and (4.20), the following estimate holds for all integer $k \in [1, m]$:*

$$\begin{aligned} \|\partial_2^k W^\pm\|_{L^2(H_\gamma^{s-k}(\omega_T))} &\leq C_s \{ \|(F^\pm, W^\pm, \xi^\pm)\|_{H_\gamma^{s-1}(\Omega_T)} + \|W^\pm\|_{L^2(H_\gamma^s(\omega_T))} \\ &\quad + \|\xi^\pm\|_{L^\infty(\Omega_T)} \|(\dot{V}, \nabla \dot{\Psi})\|_{H_\gamma^{s-1}(\Omega_T)} \\ &\quad + \|(F^\pm, W^\pm)\|_{L^\infty(\Omega_T)} \|(\dot{V}, \nabla \dot{\Psi})\|_{H_\gamma^s(\Omega_T)} \}. \end{aligned} \quad (6.16)$$

In light of definition (2.27), we see that, for all $s \in \mathbb{N}$ and $\theta \in H_\gamma^s(\Omega_T)$,

$$\|\theta\|_{H_\gamma^s(\Omega_T)} = \sum_{k=0}^s \|\partial_2^k \theta\|_{L^2(H_\gamma^{s-k}(\omega_T))}, \quad \gamma \|\theta\|_{H_\gamma^{s-1}(\Omega_T)} \leq \|\theta\|_{H_\gamma^s(\Omega_T)}.$$

By virtue of these identities, we combine Lemmas 6.2–6.4 and employ the Moser-type inequalities to obtain the following *a priori* estimates on the H_γ^{s-} norm of solution \dot{V}^\pm to the linearized problem (4.12).

Proposition 6.1 *Assume that the hypotheses of Theorem 6.1 hold. Then there exists a constant $K_0 > 0$ (independent of s and T) and constants $C_s > 0$ and $\gamma_s \geq 1$ (depending on s , but independent of T) such that, if $K \leq K_0$, then, for all $\gamma \geq \gamma_s$ and solutions $(\dot{V}, \psi) \in H_\gamma^{s+2}(\Omega_T) \times H_\gamma^{s+2}(\omega_T)$ of problem (4.12), the following estimate holds:*

$$\begin{aligned} & \sqrt{\gamma} \|\dot{V}^\pm\|_{H_\gamma^s(\Omega_T)} + \|\mathbb{P}^\pm(\dot{\varphi})\dot{V}^\pm|_{x_2=0}\|_{H_\gamma^s(\omega_T)} + \|\psi\|_{H_\gamma^{s+1}(\omega_T)} \\ & \leq C_s \{ \gamma^{-1/2} \|f^\pm\|_{H_\gamma^s(\Omega_T)} + \gamma^{-3/2} \|f^\pm\|_{L^2(H_\gamma^{s+1}(\omega_T))} + \gamma^{-1} \|g\|_{H_\gamma^{s+1}(\omega_T)} \\ & \quad + \gamma^{-1} (\|\mathbb{P}^\pm(\dot{\varphi})\dot{V}^\pm\|_{L^\infty(\omega_T)} + \|\psi\|_{W^{1,\infty}(\omega_T)}) \|(\dot{V}, \partial_2 \dot{V}, \nabla \dot{\Psi})\|_{H_\gamma^{s+1}(\omega_T)} \\ & \quad + (\gamma^{-1/2} \|f^\pm\|_{L^\infty(\Omega_T)} + \gamma^{-3/2} \|\dot{V}^\pm\|_{W^{1,\infty}(\Omega_T)}) \|(\dot{V}, \nabla \dot{\Psi})\|_{H_\gamma^{s+2}(\Omega_T)} \}. \end{aligned} \quad (6.17)$$

6.3 Proof of Theorem 6.1

Theorem 6.2 shows that the linearized problem (4.12) is well-posed for sources terms $(f^\pm, g) \in L^2(H^1(\omega_T)) \times H^1(\omega_T)$ that vanish in the past. Following [6, 41], we can use Proposition 6.1 to convert Theorem 6.2 into a well-posedness result of (4.12) in H^s . More precisely, we can prove that, under the assumptions of Theorem 6.1, if $(f^\pm, g) \in H^{s+1}(\Omega_T) \times H^{s+1}(\omega_T)$ vanish in the past, then there exists a unique solution $(\dot{V}^\pm, \psi) \in H^s(\Omega_T) \times H^{s+1}(\omega_T)$ that vanishes in the past and satisfies (6.17) for all $\gamma \geq \gamma_s$.

It remains to prove the tame estimate (6.2). To this end, we first fix the value of γ such that γ is greater than the maximum of $\gamma_3, \dots, \gamma_{\bar{\alpha}}$. Using (6.17) with $s = 3$ and (6.1), we have

$$\begin{aligned} & \|\dot{V}^\pm\|_{H_\gamma^3(\Omega_T)} + \|\mathbb{P}^\pm(\dot{\varphi})\dot{V}^\pm|_{x_2=0}\|_{H_\gamma^3(\omega_T)} + \|\psi\|_{H_\gamma^4(\omega_T)} \\ & \lesssim K (\|f^\pm\|_{L^\infty(\Omega_T)} + \|\dot{V}^\pm\|_{W^{1,\infty}(\Omega_T)} + \|\mathbb{P}^\pm(\dot{\varphi})\dot{V}^\pm|_{x_2=0}\|_{L^\infty(\omega_T)} + \|\psi\|_{W^{1,\infty}(\omega_T)}) \\ & \quad + \|f^\pm\|_{H_\gamma^4(\Omega_T)} + \|g\|_{H_\gamma^4(\omega_T)}. \end{aligned} \quad (6.18)$$

Note that $T > 0$ and γ have been fixed. Thanks to the classical Sobolev inequalities that $\|\theta\|_{L^\infty(\Omega_T)} \lesssim \|\theta\|_{H^2(\Omega_T)}$ and $\|\theta\|_{L^\infty(\omega_T)} \lesssim \|\theta\|_{H^2(\omega_T)}$, we utilize (6.18) and take $K > 0$ sufficiently small to obtain that $\|f^\pm\|_{L^\infty(\Omega_T)} \lesssim$

$\|f^\pm\|_{H_\gamma^4(\Omega_T)}$ and

$$\begin{aligned} & \|\dot{V}^\pm\|_{W^{1,\infty}(\Omega_T)} + \|\mathbb{P}^\pm(\dot{\varphi})\dot{V}^\pm|_{x_2=0}\|_{L^\infty(\omega_T)} + \|\psi\|_{W^{1,\infty}(\omega_T)} \\ & \lesssim \|\dot{V}^\pm\|_{H_\gamma^3(\Omega_T)} + \|\mathbb{P}^\pm(\dot{\varphi})\dot{V}^\pm|_{x_2=0}\|_{H_\gamma^3(\omega_T)} + \|\psi\|_{H_\gamma^4(\omega_T)} \\ & \lesssim \|f^\pm\|_{H_\gamma^4(\Omega_T)} + \|g\|_{H_\gamma^4(\omega_T)}. \end{aligned}$$

Plugging these estimates into (6.17) yields (6.2), which completes the proof of Theorem 6.1.

7 Construction of the Approximate Solution

In this section, we introduce the ‘‘approximate’’ solution (U^a, Φ^a) in order to reduce the original problem (2.20) and (2.23) into a nonlinear problem with zero initial data. We naturally expect to solve this reformulated problem in the space of functions vanishing in the past, so that Theorem 6.1, which is the well-posedness result in the same function space for the linearized problem, can be applied. We need to impose the necessary compatibility conditions on the initial data (U_0^\pm, φ_0) for the existence of smooth approximate solutions (U^a, Φ^a) , which are solutions of problem (2.20) and (2.23) in the sense of Taylor’s series at time $t = 0$.

Let $s \geq 3$ be an integer. Assume that $\tilde{U}_0^\pm := U_0^\pm - \bar{U}^\pm \in H^{s+1/2}(\mathbb{R}_+^2)$ and $\varphi_0 \in H^{s+1}(\mathbb{R})$. We also assume without loss of generality that $(\tilde{U}_0^\pm, \varphi_0)$ has a compact support:

$$\text{supp } \tilde{U}_0^\pm \subset \{x_2 \geq 0, x_1^2 + x_2^2 \leq 1\}, \quad \text{supp } \varphi_0 \subset [-1, 1]. \quad (7.1)$$

We extend φ_0 from \mathbb{R} to \mathbb{R}_+^2 by constructing $\tilde{\Phi}_0^\pm = \tilde{\Phi}_0^\pm \in H^{s+3/2}(\mathbb{R}_+^2)$, which satisfies

$$\tilde{\Phi}_0^\pm|_{x_2=0} = \varphi_0, \quad \text{supp } \tilde{\Phi}_0^\pm \subset \{x_2 \geq 0, x_1^2 + x_2^2 \leq 2\}, \quad (7.2)$$

and the estimate:

$$\|\tilde{\Phi}_0^\pm\|_{H^{s+3/2}(\mathbb{R}_+^2)} \leq C\|\varphi_0\|_{H^{s+1}(\mathbb{R})}. \quad (7.3)$$

By virtue of (7.3) and the Sobolev embedding theorem, we infer that, for φ_0 small enough in $H^{s+1}(\mathbb{R})$, the following estimates hold for $\Phi_0^\pm := \tilde{\Phi}_0^\pm + \bar{\Phi}_0^\pm$:

$$\pm \partial_2 \Phi_0^\pm \geq \frac{7}{8} \quad \text{for all } x \in \mathbb{R}_+^2. \quad (7.4)$$

For problem (2.23), we prescribe the initial data:

$$\Phi^\pm|_{t=0} = \Phi_0^\pm. \quad (7.5)$$

Let us denote the perturbation by $(\tilde{U}^\pm, \tilde{\Phi}^\pm) := (U^\pm - \bar{U}^\pm, \Phi^\pm - \bar{\Phi}^\pm)$, and the traces of the ℓ -th order time derivatives on $\{t = 0\}$ by

$$\tilde{U}_\ell^\pm := \partial_t^\ell \tilde{U}^\pm|_{t=0}, \quad \tilde{\Phi}_\ell^\pm := \partial_t^\ell \tilde{\Phi}^\pm|_{t=0}, \quad \ell \in \mathbb{N}.$$

To introduce the compatibility conditions, we need to determine traces \tilde{U}_ℓ^\pm and $\tilde{\Phi}_\ell^\pm$ in terms of the initial data \tilde{U}_0^\pm and $\tilde{\Phi}_0^\pm$ through equations (2.20a) and (2.23a). For this purpose, we set $\mathcal{W}^\pm := (\tilde{U}^\pm, \nabla_x \tilde{U}^\pm, \nabla_x \tilde{\Phi}^\pm) \in \mathbb{R}^{11}$, and rewrite (2.20a) and (2.23a) as

$$\partial_t \tilde{U}^\pm = \mathbf{F}_1(\mathcal{W}^\pm), \quad \partial_t \tilde{\Phi}^\pm = \mathbf{F}_2(\mathcal{W}^\pm), \quad (7.6)$$

where \mathbf{F}_1 and \mathbf{F}_2 are suitable C^∞ -functions that vanish at the origin. After applying operator ∂_t^ℓ to (7.6), we take the traces at time $t = 0$. One can employ the generalized Faà di Bruno's formula (cf. [37, Theorem 2.1]) to derive

$$\tilde{U}_{\ell+1}^\pm = \sum_{\alpha_i \in \mathbb{N}^{11}, |\alpha_1| + \dots + |\alpha_\ell| = \ell} D^{\alpha_1 + \dots + \alpha_\ell} \mathbf{F}_1(\mathcal{W}_0^\pm) \prod_{i=1}^{\ell} \frac{\ell!}{\alpha_i!} \left(\frac{\mathcal{W}_i^\pm}{i!} \right)^{\alpha_i}, \quad (7.7)$$

$$\tilde{\Phi}_{\ell+1}^\pm = \sum_{\alpha_i \in \mathbb{N}^{11}, |\alpha_1| + \dots + |\alpha_\ell| = \ell} D^{\alpha_1 + \dots + \alpha_\ell} \mathbf{F}_2(\mathcal{W}_0^\pm) \prod_{i=1}^{\ell} \frac{\ell!}{\alpha_i!} \left(\frac{\mathcal{W}_i^\pm}{i!} \right)^{\alpha_i}, \quad (7.8)$$

where \mathcal{W}_i^\pm denotes trace $(\tilde{U}_i^\pm, \nabla_x \tilde{U}_i^\pm, \nabla_x \tilde{\Phi}_i^\pm)$ at $t = 0$. From (7.7)–(7.8), one can determine $(\tilde{U}_\ell^\pm, \tilde{\Phi}_\ell^\pm)_{\ell \geq 0}$ inductively as functions of the initial data \tilde{U}_0^\pm and $\tilde{\Phi}_0^\pm$. Furthermore, we have the following lemma (see [34, Lemma 4.2.1] for the proof):

Lemma 7.1 *Assume that (7.1)–(7.4) hold. Then the equations (7.7)–(7.8) determine $\tilde{U}_\ell^\pm \in H^{s+1/2-\ell}(\mathbb{R}_+^2)$ for $\ell = 1, \dots, s$, and $\tilde{\Phi}_\ell^\pm \in H^{s+3/2-\ell}(\mathbb{R}_+^2)$ for $\ell = 1, \dots, s+1$, such that*

$$\text{supp } \tilde{U}_\ell^\pm \subset \{x_2 \geq 0, x_1^2 + x_2^2 \leq 1\}, \quad \text{supp } \tilde{\Phi}_\ell^\pm \subset \{x_2 \geq 0, x_1^2 + x_2^2 \leq 2\}.$$

In addition,

$$\begin{aligned} & \sum_{\ell=0}^s \|\tilde{U}_\ell^\pm\|_{H^{s+1/2-\ell}(\mathbb{R}_+^2)} + \sum_{\ell=0}^{s+1} \|\tilde{\Phi}_\ell^\pm\|_{H^{s+3/2-\ell}(\mathbb{R}_+^2)} \\ & \leq C(\|\tilde{U}_0^\pm\|_{H^{s+1/2}(\mathbb{R}_+^2)} + \|\varphi_0\|_{H^{s+1}(\mathbb{R})}), \end{aligned}$$

where constant $C > 0$ depends only on s and $\|(\tilde{U}_0^\pm, \tilde{\Phi}_0^\pm)\|_{W^{1,\infty}(\mathbb{R}_+^2)}$.

To construct a *smooth* approximate solution, one has to impose certain assumptions on traces \tilde{U}_ℓ^\pm and $\tilde{\Phi}_\ell^\pm$. We are now ready to introduce the following terminology.

Definition 7.1 (Compatibility conditions) Let $s \geq 3$ be an integer. Let $\tilde{U}_0^\pm := U_0^\pm - \bar{U}_0^\pm \in H^{s+1/2}(\mathbb{R}_+^2)$ and $\varphi_0 \in H^{s+1}(\mathbb{R})$ satisfy (7.1). The initial data U_0^\pm and φ_0 are said to be compatible up to order s if there exist functions $\tilde{\Phi}_0^\pm \in H^{s+3/2}(\mathbb{R}_+^2)$ satisfying (7.2)–(7.4) such that functions $\tilde{U}_1^\pm, \dots, \tilde{U}_s^\pm, \tilde{\Phi}_1^\pm, \dots, \tilde{\Phi}_{s+1}^\pm$ determined by (7.7)–(7.8) satisfy:

$$\partial_2^j (\tilde{\Phi}_\ell^+ - \tilde{\Phi}_\ell^-)|_{x_2=0} = 0 \quad \text{for } j, \ell \in \mathbb{N} \text{ with } j + \ell < s + 1, \quad (7.9a)$$

$$\partial_2^j (\tilde{p}_\ell^+ - \tilde{p}_\ell^-)|_{x_2=0} = 0 \quad \text{for } j, \ell \in \mathbb{N} \text{ with } j + \ell < s, \quad (7.9b)$$

and

$$\int_{\mathbb{R}_+^2} |\partial_2^{s+1-\ell} (\tilde{\Phi}_\ell^+ - \tilde{\Phi}_\ell^-)|^2 dx_1 \frac{dx_2}{x_2} < \infty \quad \text{for } \ell = 0, \dots, s + 1, \quad (7.10a)$$

$$\int_{\mathbb{R}_+^2} |\partial_2^{s-\ell} (\tilde{p}_\ell^+ - \tilde{p}_\ell^-)|^2 dx_1 \frac{dx_2}{x_2} < \infty \quad \text{for } \ell = 0, \dots, s. \quad (7.10b)$$

It follows from Lemma 7.1 that $\tilde{p}_0^\pm, \dots, \tilde{p}_{s-2}^\pm, \tilde{\Phi}_0^\pm, \dots, \tilde{\Phi}_{s-1}^\pm \in H^{5/2}(\mathbb{R}_+^2) \subset W^{1,\infty}(\mathbb{R}_+^2)$. Then we can take the j -th order derivatives of the traces in (7.9). In what follows, we employ $\varepsilon_0(\cdot)$ to denote a function that tends to 0 when its argument tends to 0. Relations (7.9)–(7.10) enable us to utilize the lifting result in [31, Theorem 2.3] to construct the approximate solution in the following lemma. We refer to [19, Lemma 3] for the proof.

Lemma 7.2 *Let $s \geq 3$ be an integer. Assume that $\tilde{U}_0^\pm := U_0^\pm - \bar{U}_0^\pm \in H^{s+1/2}(\mathbb{R}_+^2)$ and $\varphi_0 \in H^{s+1}(\mathbb{R})$ satisfy (7.1), and that the initial data U_0^\pm and φ_0 are compatible up to order s . If \tilde{U}_0^\pm and φ_0 are sufficiently small, then there exist functions $U^{a\pm}, \Phi^{a\pm}$, and φ^a such that $\tilde{U}^{a\pm} := U^{a\pm} - \bar{U}^\pm \in H^{s+1}(\Omega)$, $\tilde{\Phi}^{a\pm} := \Phi^{a\pm} - \bar{\Phi}^\pm \in H^{s+2}(\Omega)$, $\varphi^a \in H^{s+3/2}(\partial\Omega)$, and*

$$\partial_t \Phi^{a\pm} + v_1^{a\pm} \partial_1 \Phi^{a\pm} - v_2^{a\pm} = 0 \quad \text{in } \Omega, \quad (7.11a)$$

$$\partial_t^j \mathbb{L}(U^{a\pm}, \Phi^{a\pm})|_{t=0} = 0 \quad \text{for } j = 0, \dots, s - 1, \quad (7.11b)$$

$$\Phi^{a+} = \Phi^{a-} = \varphi^a \quad \text{on } \partial\Omega, \quad (7.11c)$$

$$\mathbb{B}(U^{a+}, U^{a-}, \varphi^a) = 0 \quad \text{on } \partial\Omega. \quad (7.11d)$$

Furthermore, we have

$$\pm \partial_2 \Phi^{a\pm} \geq \frac{3}{4} \quad \text{for all } (t, x) \in \Omega, \quad (7.12)$$

$$\text{supp}(\tilde{U}^{a\pm}, \tilde{\Phi}^{a\pm}) \subset \{t \in [-T, T], x_2 \geq 0, x_1^2 + x_2^2 \leq 3\}, \quad (7.13)$$

and

$$\begin{aligned} & \|\tilde{U}^{a\pm}\|_{H^{s+1}(\Omega)} + \|\tilde{\Phi}^{a\pm}\|_{H^{s+2}(\Omega)} + \|\varphi^a\|_{H^{s+3/2}(\partial\Omega)} \\ & \leq \varepsilon_0 (\|\tilde{U}_0^\pm\|_{H^{s+1/2}(\mathbb{R}_+^2)} + \|\varphi_0\|_{H^{s+1}(\mathbb{R})}). \end{aligned} \quad (7.14)$$

Let us denote $U^a := (U^{a+}, U^{a-})^\top$ and $\Phi^a := (\Phi^{a+}, \Phi^{a-})^\top$. Vector (U^a, Φ^a) in Lemma 7.2 is called the approximate solution to problem (2.20) and (2.23). Relations (7.11c) and (7.13) immediately imply that φ^a is supported within $\{-T \leq t \leq T, x_1^2 \leq 3\}$. Since $s \geq 3$, it follows from (7.14) and the Sobolev embedding theorem that

$$\|\tilde{U}^{a\pm}\|_{W^{2,\infty}(\Omega)} + \|\tilde{\Phi}^{a\pm}\|_{W^{3,\infty}(\Omega)} \leq \varepsilon_0 (\|\tilde{U}_0^\pm\|_{H^{s+1/2}(\mathbb{R}_+^2)} + \|\varphi_0\|_{H^{s+1}(\mathbb{R})}).$$

We are going to reformulate the original problem into that with zero initial data by using the approximate solution (U^a, Φ^a) . Let us introduce

$$f^a := \begin{cases} -\mathbb{L}(U^a, \Phi^a) & \text{if } t > 0, \\ 0 & \text{if } t < 0. \end{cases} \quad (7.15)$$

Since $(\tilde{U}^{a\pm}, \nabla \tilde{\Phi}^{a\pm}) \in H^{s+1}(\Omega)$, taking into account (7.11b) and (7.13), we obtain that $f^a \in H^s(\Omega)$ and

$$\text{supp } f^a \subset \{0 \leq t \leq T, x_2 \geq 0, x_1^2 + x_2^2 \leq 3\}.$$

Using the Moser-type inequalities and the fact that f^a vanishes in the past, we obtain from (7.14) the estimate:

$$\|f^a\|_{H^s(\Omega)} \leq \varepsilon_0 (\|\tilde{U}_0^\pm\|_{H^{s+1/2}(\mathbb{R}_+^2)} + \|\varphi_0\|_{H^{s+1}(\mathbb{R})}). \quad (7.16)$$

Let (U^a, Φ^a) be the approximate solution defined in Lemma 7.2. By virtue of (7.11) and (7.15), we see that $(U, \Phi) = (U^a, \Phi^a) + (V, \Psi)$ is a solution of the original problem (2.20) and (2.23) on $[0, T] \times \mathbb{R}_+^2$, if $V = (V^+, V^-)^\top$ and $\Psi = (\Psi^+, \Psi^-)^\top$ solve the following problem:

$$\begin{cases} \mathcal{L}(V, \Psi) := \mathbb{L}(U^a + V, \Phi^a + \Psi) - \mathbb{L}(U^a, \Phi^a) = f^a & \text{in } \Omega_T, \\ \mathcal{E}(V, \Psi) := \partial_t \Psi + (v_1^a + v_1) \partial_1 \Psi + v_1 \partial_1 \Phi^a - v_2 = 0 & \text{in } \Omega_T, \\ \mathcal{B}(V, \psi) := \mathbb{B}(U^a + V, \varphi^a + \psi) = 0 & \text{on } \omega_T, \\ \Psi^+ = \Psi^- = \psi & \text{on } \omega_T, \\ (V, \Psi) = 0, & \text{for } t < 0. \end{cases} \quad (7.17)$$

The initial data (2.20c) and (7.5) have been absorbed into the interior equations. From (7.11a) and (7.11d), we observe that $(V, \Psi) = 0$ satisfies (7.17) for $t < 0$. Therefore, the original nonlinear problem on $[0, T] \times \mathbb{R}_+^2$ is now reformulated as a problem on Ω_T whose solutions vanish in the past.

8 Nash–Moser Iteration

In this section, we prove the existence of solutions to problem (7.17) by a suitable iteration scheme of Nash–Moser type (*cf.* Hörmander [24]). First, we introduce the smoothing operators S_θ and describe the iterative scheme for problem (7.17).

8.1 The Iterative Scheme

We first state the following result from [19, Proposition 4].

Proposition 8.1 *Let $T > 0$ and $\gamma \geq 1$, and let $m \geq 4$ be an integer. Then there exists a family $\{S_\theta\}_{\theta \geq 1}$ of smoothing operators:*

$$S_\theta : \mathcal{F}_\gamma^3(\Omega_T) \times \mathcal{F}_\gamma^3(\Omega_T) \longrightarrow \bigcap_{\beta \geq 3} \mathcal{F}_\gamma^\beta(\Omega_T) \times \mathcal{F}_\gamma^\beta(\Omega_T),$$

where $\mathcal{F}_\gamma^\beta(\Omega_T) := \{u \in H_\gamma^\beta(\Omega_T) : u = 0 \text{ for } t < 0\}$ is a closed subspace of $H_\gamma^\beta(\Omega_T)$ such that

$$\|S_\theta u\|_{H_\gamma^\beta(\Omega_T)} \leq C\theta^{(\beta-\alpha)_+} \|u\|_{H_\gamma^\alpha(\Omega_T)} \quad \text{for all } \alpha, \beta \in [1, m], \quad (8.1a)$$

$$\|S_\theta u - u\|_{H_\gamma^\beta(\Omega_T)} \leq C\theta^{\beta-\alpha} \|u\|_{H_\gamma^\alpha(\Omega_T)} \quad \text{for all } 1 \leq \beta \leq \alpha \leq m, \quad (8.1b)$$

$$\left\| \frac{d}{d\theta} S_\theta u \right\|_{H_\gamma^\beta(\Omega_T)} \leq C\theta^{\beta-\alpha-1} \|u\|_{H_\gamma^\alpha(\Omega_T)} \quad \text{for all } \alpha, \beta \in [1, m], \quad (8.1c)$$

and

$$\begin{aligned} & \| (S_\theta u - S_\theta v)|_{x_2=0} \|_{H_\gamma^\beta(\omega_T)} \\ & \leq C\theta^{(\beta+1-\alpha)_+} \| (u - v)|_{x_2=0} \|_{H_\gamma^\alpha(\omega_T)} \quad \text{for all } \alpha, \beta \in [1, m], \end{aligned} \quad (8.2)$$

where $\alpha, \beta \in \mathbb{N}$, $(\beta - \alpha)_+ := \max\{0, \beta - \alpha\}$, and $C > 0$ is a constant depending only on m . In particular, if $u = v$ on ω_T , then $S_\theta u = S_\theta v$ on ω_T . Furthermore, there exists another family of smoothing operators (still denoted by S_θ) acting on the functions defined on the boundary ω_T and satisfying the properties in (8.1) with norms $\|\cdot\|_{H_\gamma^\alpha(\omega_T)}$.

The proof of (8.2) is based on the following lifting operator (see [22, Chapter 5] and [19]).

Lemma 8.1 *Let $T > 0$ and $\gamma \geq 1$, and let $m \geq 1$ be an integer. Then there exists an operator \mathcal{R}_T , which is continuous from $\mathcal{F}_\gamma^s(\omega_T)$ to $\mathcal{F}_\gamma^{s+1/2}(\Omega_T)$ for all $s \in [1, m]$, such that, if $s \geq 1$ and $u \in \mathcal{F}_\gamma^s(\omega_T)$, then $(\mathcal{R}_T u)|_{x_2=0} = u$.*

Now, following [19], we describe the iteration scheme for problem (7.17).

Assumption (A-1): $(V_0, \Psi_0, \psi_0) = 0$ and, for $k = 0, \dots, n$, (V_k, Ψ_k, ψ_k) are already given and satisfy

$$(V_k, \Psi_k, \psi_k)|_{t < 0} = 0, \quad \Psi_k^+|_{x_2=0} = \Psi_k^-|_{x_2=0} = \psi_k. \quad (8.3)$$

Let us consider

$$V_{n+1} = V_n + \delta V_n, \quad \Psi_{n+1} = \Psi_n + \delta \Psi_n, \quad \psi_{n+1} = \psi_n + \delta \psi_n, \quad (8.4)$$

where these differences δV_n , $\delta \Psi_n$, and $\delta \psi_n$ will be specified below.

First we are going to find $(\delta\dot{V}_n, \delta\psi_n)$ by solving the effective linear problem:

$$\begin{cases} \mathbb{L}'_e(U^a + V_{n+1/2}, \Phi^a + \Psi_{n+1/2})\delta\dot{V}_n = f_n & \text{in } \Omega_T, \\ \mathbb{B}'_e(U^a + V_{n+1/2}, \Phi^a + \Psi_{n+1/2})(\delta\dot{V}_n, \delta\psi_n) = g_n & \text{on } \omega_T, \\ (\delta\dot{V}_n, \delta\psi_n) = 0 & \text{for } t < 0, \end{cases} \quad (8.5)$$

where operators \mathbb{L}'_e and \mathbb{B}'_e are defined in (4.12a)–(4.12b),

$$\delta\dot{V}_n := \delta V_n - \frac{\partial_2(U^a + V_{n+1/2})}{\partial_2(\Phi^a + \Psi_{n+1/2})} \delta\Psi_n \quad (8.6)$$

is the “good unknown” (*cf.* (4.10)), and $(V_{n+1/2}, \Psi_{n+1/2})$ is a smooth modified state such that $(U^a + V_{n+1/2}, \Phi^a + \Psi_{n+1/2})$ satisfies constraints (4.1)–(4.3). The source term (f_n, g_n) will be defined through the accumulated error terms at Step n later on.

We define the modified state as

$$\begin{cases} \Psi_{n+1/2}^\pm := S_{\theta_n} \Psi_n^\pm, & v_1(V_{n+1/2}^\pm) := S_{\theta_n} v_1(V_n^\pm), \\ p_{n+1/2}^\pm := S_{\theta_n} p_n^\pm \mp \frac{1}{2} \mathcal{R}_T(S_{\theta_n} p_n^+|_{x_2=0} - S_{\theta_n} p_n^-|_{x_2=0}), \\ v_2(V_{n+1/2}^\pm) := \partial_t \Psi_{n+1/2}^\pm + (v_1^{a\pm} + v_1(V_{n+1/2}^\pm)) \partial_1 \Psi_{n+1/2}^\pm \\ \quad + v_1(V_{n+1/2}^\pm) \partial_1 \Phi^{a\pm}, \end{cases} \quad (8.7)$$

where S_{θ_n} are the smoothing operators defined in Proposition 8.1 with sequence $\{\theta_n\}$ given by

$$\theta_0 \geq 1, \quad \theta_n = \sqrt{\theta_0^2 + n}, \quad (8.8)$$

and \mathcal{R}_T is the lifting operator given in Lemma 8.1. Thanks to (8.3), we have

$$\begin{cases} \Psi_{n+1/2}^+|_{x_2=0} = \Psi_{n+1/2}^-|_{x_2=0} =: \psi_{n+1/2}, \\ p_{n+1/2}^+|_{x_2=0} = p_{n+1/2}^-|_{x_2=0}, \\ \mathcal{E}(V_{n+1/2}, \Psi_{n+1/2}) = 0, \\ (V_{n+1/2}, \Psi_{n+1/2}, \psi_{n+1/2})|_{t<0} = 0. \end{cases} \quad (8.9)$$

It then follows from (7.11) that $(U^a + V_{n+1/2}, \Phi^a + \Psi_{n+1/2})$ satisfies (4.3a) and (4.3c)–(4.3d). The additional constraint (4.3b) will be obtained by taking the initial data small enough.

The error terms at Step n are defined from the following decompositions:

$$\begin{aligned} & \mathcal{L}(V_{n+1}, \Psi_{n+1}) - \mathcal{L}(V_n, \Psi_n) \\ &= \mathbb{L}'(U^a + V_n, \Phi^a + \Psi_n)(\delta V_n, \delta\Psi_n) + e'_n \\ &= \mathbb{L}'(U^a + S_{\theta_n} V_n, \Phi^a + S_{\theta_n} \Psi_n)(\delta V_n, \delta\Psi_n) + e'_n + e''_n \\ &= \mathbb{L}'(U^a + V_{n+1/2}, \Phi^a + \Psi_{n+1/2})(\delta V_n, \delta\Psi_n) + e'_n + e''_n + e'''_n \\ &= \mathbb{L}'_e(U^a + V_{n+1/2}, \Phi^a + \Psi_{n+1/2})\delta\dot{V}_n + e'_n + e''_n + e'''_n + D_{n+1/2}\delta\Psi_n \end{aligned} \quad (8.10)$$

and

$$\begin{aligned}
& \mathcal{B}(V_{n+1}|_{x_2=0}, \psi_{n+1}) - \mathcal{B}(V_n|_{x_2=0}, \psi_n) \\
&= \mathbb{B}'(U^a + V_n, \Phi^a + \Psi_n)(\delta V_n|_{x_2=0}, \delta\psi_n) + \tilde{e}'_n \\
&= \mathbb{B}'(U^a + S_{\theta_n} V_n, \Phi^a + S_{\theta_n} \Psi_n)(\delta V_n|_{x_2=0}, \delta\psi_n) + \tilde{e}'_n + \tilde{e}''_n \\
&= \mathbb{B}'_e(U^a + V_{n+1/2}, \Phi^a + \Psi_{n+1/2})(\delta \dot{V}_n|_{x_2=0}, \delta\psi_n) + \tilde{e}'_n + \tilde{e}''_n + \tilde{e}'''_n, \quad (8.11)
\end{aligned}$$

where we have set

$$D_{n+1/2} := \frac{1}{\partial_2(\Phi^a + \Psi_{n+1/2})} \partial_2 \mathbb{L}(U^a + V_{n+1/2}, \Phi^a + \Psi_{n+1/2}), \quad (8.12)$$

and have used (4.11) to obtain the last identity in (8.10). Let us set

$$e_n := e'_n + e''_n + e'''_n + D_{n+1/2} \delta \Psi_n, \quad \tilde{e}_n := \tilde{e}'_n + \tilde{e}''_n + \tilde{e}'''_n. \quad (8.13)$$

Assumption (A-2): $f_0 := S_{\theta_0} f^a$, $(e_0, \tilde{e}_0, g_0) := 0$, and $(f_k, g_k, e_k, \tilde{e}_k)$ are already given and vanish in the past for $k = 0, \dots, n-1$.

We compute the accumulated error terms at Step n , $n \geq 1$, by

$$E_n := \sum_{k=0}^{n-1} e_k, \quad \tilde{E}_n := \sum_{k=0}^{n-1} \tilde{e}_k. \quad (8.14)$$

Then we compute f_n and g_n for $n \geq 1$ from the equations:

$$\sum_{k=0}^n f_k + S_{\theta_n} E_n = S_{\theta_n} f^a, \quad \sum_{k=0}^n g_k + S_{\theta_n} \tilde{E}_n = 0. \quad (8.15)$$

Under assumptions (A-1)–(A-2), $(V_{n+1/2}, \Psi_{n+1/2})$ and (f_n, g_n) have been specified from (8.7) and (8.15). Then we can obtain $(\delta \dot{V}_n, \delta\psi_n)$ as the solution of the linear problem (8.5) by applying Theorem 6.1.

Next we need to construct $\delta \Psi_n = (\delta \Psi_n^+, \delta \Psi_n^-)^T$ satisfying $\delta \Psi_n^\pm|_{x_2=0} = \delta\psi_n$. We use the boundary conditions in (8.5) (cf. (4.7)–(4.8), (4.13), and (8.6)) to derive that $\delta\psi_n$ satisfies

$$\begin{aligned}
& \left(\frac{\epsilon^2 \partial_t \Phi_{n+1/2}^+}{N_{n+1/2}^+ h_{n+1/2}^+ (\Gamma_{n+1/2}^+)^2}, \frac{\varrho_{n+1/2}^+}{h_{n+1/2}^+ \Gamma_{n+1/2}^+}, \frac{-\varsigma_{n+1/2}^+}{h_{n+1/2}^+ \Gamma_{n+1/2}^+} \right) \Big|_{x_2=0} \\
& \times \left(\delta \dot{V}_n^+ + \delta\psi_n \frac{\partial_2 U_{n+1/2}^+}{\partial_2 \Phi_{n+1/2}^+} \right) \Big|_{x_2=0} + \partial_t(\delta\psi_n) + v_1(U_{n+1/2}^+) \Big|_{x_2=0} \partial_1(\delta\psi_n) \\
&= g_{n,2}, \\
& \left(\frac{\epsilon^2 \partial_t \Phi_{n+1/2}^-}{N_{n+1/2}^- h_{n+1/2}^- (\Gamma_{n+1/2}^-)^2}, \frac{\varrho_{n+1/2}^-}{h_{n+1/2}^- \Gamma_{n+1/2}^-}, \frac{-\varsigma_{n+1/2}^-}{h_{n+1/2}^- \Gamma_{n+1/2}^-} \right) \Big|_{x_2=0} \\
& \times \left(\delta \dot{V}_n^- + \delta\psi_n \frac{\partial_2 U_{n+1/2}^-}{\partial_2 \Phi_{n+1/2}^-} \right) \Big|_{x_2=0} + \partial_t(\delta\psi_n) + v_1(U_{n+1/2}^-) \Big|_{x_2=0} \partial_1(\delta\psi_n) \\
&= g_{n,2} - g_{n,1},
\end{aligned}$$

where we have set $U_{n+1/2}^\pm := U^{a\pm} + V_{n+1/2}^\pm$, $\Phi_{n+1/2}^\pm := \Phi^{a\pm} + \Psi_{n+1/2}^\pm$, and

$$(N_{n+1/2}^\pm, h_{n+1/2}^\pm, \Gamma_{n+1/2}^\pm, \varrho_{n+1/2}^\pm, \varsigma_{n+1/2}^\pm) := (N, h, \Gamma, \varrho, \varsigma)(U_{n+1/2}^\pm, \Phi_{n+1/2}^\pm),$$

with $(N, h, \Gamma, \varrho, \varsigma)$ defined in (2.11), (3.3), and (4.9). Then we define $\delta\Psi_n^+$ and $\delta\Psi_n^-$ as solutions to the following equations:

$$\begin{aligned} & \left(\frac{\epsilon^2 \partial_t \Phi_{n+1/2}^+}{N_{n+1/2}^+ h_{n+1/2}^+ (\Gamma_{n+1/2}^+)^2}, \frac{\varrho_{n+1/2}^+}{h_{n+1/2}^+ \Gamma_{n+1/2}^+}, \frac{-\varsigma_{n+1/2}^+}{h_{n+1/2}^+ \Gamma_{n+1/2}^+} \right) \\ & \times \left(\delta\Psi_n^+ \frac{\partial_2 U_{n+1/2}^+}{\partial_2 \Phi_{n+1/2}^+} + \delta\dot{V}_n^+ \right) + \partial_t(\delta\Psi_n^+) + v_1(U_{n+1/2}^+) \partial_1(\delta\Psi_n^+) \\ & = \mathcal{R}_T g_{n,2} + G_n^+, \end{aligned} \quad (8.16)$$

$$\begin{aligned} & \left(\frac{\epsilon^2 \partial_t \Phi_{n+1/2}^-}{N_{n+1/2}^- h_{n+1/2}^- (\Gamma_{n+1/2}^-)^2}, \frac{\varrho_{n+1/2}^-}{h_{n+1/2}^- \Gamma_{n+1/2}^-}, \frac{-\varsigma_{n+1/2}^-}{h_{n+1/2}^- \Gamma_{n+1/2}^-} \right) \\ & \times \left(\delta\Psi_n^- \frac{\partial_2 U_{n+1/2}^-}{\partial_2 \Phi_{n+1/2}^-} + \delta\dot{V}_n^- \right) + \partial_t(\delta\Psi_n^-) + v_1(U_{n+1/2}^-) \partial_1(\delta\Psi_n^-) \\ & = \mathcal{R}_T(g_{n,2} - g_{n,1}) + G_n^-, \end{aligned} \quad (8.17)$$

where the source terms G_n^\pm will be chosen by using a decomposition for operator \mathcal{E} .

We define the error terms: \hat{e}'_n , \hat{e}''_n , and \hat{e}'''_n by

$$\begin{aligned} \mathcal{E}(V_{n+1}, \Psi_{n+1}) - \mathcal{E}(V_n, \Psi_n) &= \mathcal{E}'(V_n, \Psi_n)(\delta V_n, \delta \Psi_n) + \hat{e}'_n \\ &= \mathcal{E}'(S_{\theta_n} V_n, S_{\theta_n} \Psi_n)(\delta V_n, \delta \Psi_n) + \hat{e}'_n + \hat{e}''_n \\ &= \mathcal{E}'(V_{n+1/2}, \Psi_{n+1/2})(\delta V_n, \delta \Psi_n) + \hat{e}'_n + \hat{e}''_n + \hat{e}'''_n, \end{aligned} \quad (8.18)$$

and denote

$$\hat{e}_n := \hat{e}'_n + \hat{e}''_n + \hat{e}'''_n, \quad \hat{E}_n := \sum_{k=0}^{n-1} \hat{e}_k. \quad (8.19)$$

From (7.11a), we have

$$\mathcal{E}(V, \Psi) = \partial_t(U^a + V) + v_1(U^a + V) \partial_1(\Phi^a + \Psi) - v_2(U^a + V).$$

Similar to the derivation of (4.6) and (4.12b), we deduce that

$$\mathcal{E}'(V_{n+1/2}^\pm, \Psi_{n+1/2}^\pm)(\delta V_n^\pm, \delta \Psi_n^\pm)$$

are equal to the left-hand sides of (8.16)–(8.17), respectively. Then it follows from (8.16)–(8.18) that

$$\mathcal{E}(V_{n+1}, \Psi_{n+1}) - \mathcal{E}(V_n, \Psi_n) = \begin{pmatrix} \mathcal{R}_T g_{n,2} + G_n^+ + \hat{e}_n^+ \\ \mathcal{R}_T(g_{n,2} - g_{n,1}) + G_n^- + \hat{e}_n^- \end{pmatrix}.$$

Summing these relations and using $\mathcal{E}(V_0, \Psi_0) = 0$ yield

$$\mathcal{E}(V_{n+1}^-, \Psi_{n+1}^-) = \mathcal{R}_T \left(\sum_{k=0}^n (g_{k,2} - g_{k,1}) \right) + \sum_{k=0}^n G_k^- + \hat{E}_{n+1}^-.$$

On the other hand, we obtain from (8.5) and (8.11) that

$$g_n = \mathcal{B}(V_{n+1}|_{x_2=0}, \psi_{n+1}) - \mathcal{B}(V_n|_{x_2=0}, \psi_n) - \tilde{e}_n. \quad (8.20)$$

In view of (7.17) and (2.22), one obtains the relations:

$$\begin{aligned} & (\mathcal{B}(V_{n+1}|_{x_2=0}, \psi_{n+1}))_2 \\ &= \mathcal{E}(V_{n+1}^+|_{x_2=0}, \psi_{n+1}) \\ &= \mathcal{E}(V_{n+1}^-|_{x_2=0}, \psi_{n+1}) + (\mathcal{B}(V_{n+1}|_{x_2=0}, \psi_{n+1}))_1. \end{aligned} \quad (8.21)$$

Summing (8.20) and using $\mathcal{B}(V_0|_{x_2=0}, \psi_0) = 0$, we have

$$\begin{aligned} & \mathcal{E}(V_{n+1}^-, \Psi_{n+1}^-) \\ &= \mathcal{R}_T \left(\mathcal{E}(V_{n+1}^-|_{x_2=0}, \psi_{n+1}) - \tilde{E}_{n+1,2} + \tilde{E}_{n+1,1} \right) + \sum_{k=0}^n G_k^- + \hat{E}_{n+1}^-. \end{aligned} \quad (8.22)$$

Similarly, we can also obtain

$$\begin{aligned} & \mathcal{E}(V_{n+1}^+, \Psi_{n+1}^+) \\ &= \mathcal{R}_T \left(\mathcal{E}(V_{n+1}^+|_{x_2=0}, \psi_{n+1}) - \tilde{E}_{n+1,2} \right) + \sum_{k=0}^n G_k^+ + \hat{E}_{n+1}^+. \end{aligned} \quad (8.23)$$

Assumption (A-3): $(G_0^+, G_0^-, \hat{e}_0) = 0$, and $(G_k^+, G_k^-, \hat{e}_k)$ are already given and vanish in the past for $k = 0, \dots, n-1$.

Under assumptions (A-1)–(A-3), taking into account (8.22)–(8.23) and the property of \mathcal{R}_T , we compute the source terms G_n^\pm from

$$S_{\theta_n} (\hat{E}_n^+ - \mathcal{R}_T \tilde{E}_{n,2}) + \sum_{k=0}^n G_k^+ = 0, \quad (8.24a)$$

$$S_{\theta_n} (\hat{E}_n^- - \mathcal{R}_T \tilde{E}_{n,2} + \mathcal{R}_T \tilde{E}_{n,1}) + \sum_{k=0}^n G_k^- = 0. \quad (8.24b)$$

From assumption (A-3) and the properties of S_θ , it is clear that G_n^\pm vanish in the past. As in [22], one can also check that the trace of G_n^\pm on ω_T vanishes. Hence, we can find $\delta\Psi_n^\pm$, vanishing in the past and satisfying $\delta\Psi_n^\pm|_{x_2=0} = \delta\psi_n$, as the unique smooth solutions to the transport equations (8.16)–(8.17).

Once $\delta\Psi_n$ is specified, we can obtain δV_n from (8.6) and $(V_{n+1}, \Psi_{n+1}, \psi_{n+1})$ from (8.4). The error terms: $e'_n, e''_n, e'''_n, \tilde{e}'_n, \tilde{e}''_n, \tilde{e}'''_n, \hat{e}'_n, \hat{e}''_n$, and \hat{e}'''_n are computed from (8.10)–(8.11) and (8.18). Then e_n, \tilde{e}_n , and \hat{e}_n are obtained from (8.13) and (8.19).

Using (8.5) and (8.15), we sum (8.10) and (8.20) from $n = 0$ to m , respectively, to obtain

$$\mathcal{L}(V_{m+1}, \Psi_{m+1}) = \sum_{n=0}^m f_n + E_{m+1} = S_{\theta_m} f^a + (I - S_{\theta_m}) E_m + e_m, \quad (8.25)$$

$$\mathcal{B}(V_{m+1}|_{x_2=0}, \psi_{m+1}) = \sum_{n=0}^m g_n + \tilde{E}_{m+1} = (I - S_{\theta_m}) \tilde{E}_m + \tilde{e}_m. \quad (8.26)$$

Plugging (8.24) into (8.22)–(8.23), we utilize (8.21) to deduce

$$\left\{ \begin{array}{l} \mathcal{E}(V_{m+1}^-, \Psi_{m+1}^-) = \mathcal{R}_T((\mathcal{B}(V_{m+1}|_{x_2=0}, \psi_{m+1}))_2 - (\mathcal{B}(V_{m+1}|_{x_2=0}, \psi_{m+1}))_1) \\ \quad + (I - S_{\theta_m})(\hat{E}_m^- - \mathcal{R}_T(\tilde{E}_{m,2} - \tilde{E}_{m,1})) \\ \quad + \hat{e}_m^- - \mathcal{R}_T(\tilde{e}_{m,2} - \tilde{e}_{m,1}), \\ \mathcal{E}(V_{m+1}^+, \Psi_{m+1}^+) = \mathcal{R}_T((\mathcal{B}(V_{m+1}|_{x_2=0}, \psi_{m+1}))_2) \\ \quad + (I - S_{\theta_m})(\hat{E}_m^+ - \mathcal{R}_T \tilde{E}_{m,2}) + \hat{e}_m^+ - \mathcal{R}_T \tilde{e}_{m,2}. \end{array} \right. \quad (8.27)$$

Since $S_{\theta_m} \rightarrow I$ as $m \rightarrow \infty$, we can formally obtain the solution to problem (7.17) from

$$\mathcal{L}(V_{m+1}, \Psi_{m+1}) \rightarrow f^a, \quad \mathcal{B}(V_{m+1}|_{x_2=0}, \psi_{m+1}) \rightarrow 0, \quad \mathcal{E}(V_{m+1}, \Psi_{m+1}) \rightarrow 0,$$

provided that the error terms: $(e_m, \tilde{e}_m, \hat{e}_m) \rightarrow 0$.

8.2 Inductive Hypothesis

Given a constant $\varepsilon > 0$ and an integer $\tilde{\alpha}$ that will be chosen later on, we assume that **(A-1)**–**(A-3)** are satisfied and that the following estimate holds:

$$\|\tilde{U}^a\|_{H_{\tilde{\gamma}}^{\tilde{\alpha}+3}(\Omega_T)} + \|\tilde{\Phi}^a\|_{H_{\tilde{\gamma}}^{\tilde{\alpha}+4}(\Omega_T)} + \|\varphi^a\|_{H_{\tilde{\gamma}}^{\tilde{\alpha}+7/2}(\Omega_T)} + \|f^a\|_{H_{\tilde{\gamma}}^{\tilde{\alpha}+2}(\Omega_T)} \leq \varepsilon. \quad (8.28)$$

Given another integer α , our inductive hypothesis reads:

$$(H_{n-1}) \left\{ \begin{array}{l} \text{(a)} \quad \|(\delta V_k, \delta \Psi_k)\|_{H_{\tilde{\gamma}}^s(\Omega_T)} + \|\delta \psi_k\|_{H_{\tilde{\gamma}}^{s+1}(\omega_T)} \leq \varepsilon \theta_k^{s-\alpha-1} \Delta_k \\ \quad \text{for all } k = 0, \dots, n-1 \text{ and } s \in [3, \tilde{\alpha}] \cap \mathbb{N}; \\ \text{(b)} \quad \|\mathcal{L}(V_k, \Psi_k) - f^a\|_{H_{\tilde{\gamma}}^s(\Omega_T)} \leq 2\varepsilon \theta_k^{s-\alpha-1} \\ \quad \text{for all } k = 0, \dots, n-1 \text{ and } s \in [3, \tilde{\alpha} - 2] \cap \mathbb{N}; \\ \text{(c)} \quad \|\mathcal{B}(V_k|_{x_2=0}, \psi_k)\|_{H_{\tilde{\gamma}}^s(\omega_T)} \leq \varepsilon \theta_k^{s-\alpha-1} \\ \quad \text{for all } k = 0, \dots, n-1 \text{ and } s \in [4, \alpha] \cap \mathbb{N}; \\ \text{(d)} \quad \|\mathcal{E}(V_k, \Psi_k)\|_{H_{\tilde{\gamma}}^3(\Omega_T)} \leq \varepsilon \theta_k^{2-\alpha} \quad \text{for all } k = 0, \dots, n-1, \end{array} \right.$$

where $\Delta_k := \theta_{k+1} - \theta_k$ with θ_k defined by (8.8). Notice that

$$\frac{1}{3\theta_k} \leq \Delta_k = \sqrt{\theta_k^2 + 1} - \theta_k \leq \frac{1}{2\theta_k} \quad \text{for all } k \in \mathbb{N}.$$

In particular, sequence (Δ_k) is decreasing and tends to 0. Our goal is to show that, for a suitable choice of parameters $\theta_0 \geq 1$ and $\varepsilon > 0$, and for f^a small enough, (H_{n-1}) implies (H_n) and that (H_0) holds. Once this goal is achieved, we infer that (H_n) holds for all $n \in \mathbb{N}$, which enables us to conclude the proof of Theorem 2.1.

From now on, we assume that (H_{n-1}) holds. As in [19], hypothesis (H_{n-1}) yields the following consequences:

Lemma 8.2 *If θ_0 is large enough, then, for each $k = 0, \dots, n$ and each integer $s \in [3, \tilde{\alpha}]$,*

$$\|(V_k, \Psi_k)\|_{H_\gamma^s(\Omega_T)} + \|\psi_k\|_{H_\gamma^{s+1}(\omega_T)} \leq \begin{cases} \varepsilon \theta_k^{(s-\alpha)_+} & \text{if } s \neq \alpha, \\ \varepsilon \log \theta_k & \text{if } s = \alpha, \end{cases} \quad (8.29)$$

$$\|((I - S_{\theta_k})V_k, (I - S_{\theta_k})\Psi_k)\|_{H_\gamma^s(\Omega_T)} \leq C\varepsilon \theta_k^{s-\alpha}. \quad (8.30)$$

Furthermore, for each $k = 0, \dots, n$, and each integer $s \in [3, \tilde{\alpha} + 4]$,

$$\|(S_{\theta_k}V_k, S_{\theta_k}\Psi_k)\|_{H_\gamma^s(\Omega_T)} \leq \begin{cases} C\varepsilon \theta_k^{(s-\alpha)_+} & \text{if } s \neq \alpha, \\ C\varepsilon \log \theta_k & \text{if } s = \alpha. \end{cases} \quad (8.31)$$

Estimates (8.30)–(8.31) follow directly from (8.1) and (8.29).

8.3 Estimate of the Error Terms

To deduce (H_n) from (H_{n-1}) , we need to estimate the quadratic error terms e'_k , \tilde{e}'_k , and \hat{e}'_k , the first substitution error terms e''_k , \tilde{e}''_k , and \hat{e}''_k , the second substitution error terms e'''_k , \tilde{e}'''_k , and \hat{e}'''_k , and the last error term $D_{k+1/2}\delta\Psi_k$. Recall from (8.10) that

$$e'_k = \mathcal{L}(V_{k+1}, \Psi_{k+1}) - \mathcal{L}(V_k, \Psi_k) - \mathbb{L}'(U^a + V_k, \Phi^a + \Psi_k)(\delta V_k, \delta \Psi_k),$$

which can be rewritten as

$$e'_k = \int_0^1 \mathbb{L}''(U^a + V_k + \tau\delta V_k, \Phi^a + \Psi_k + \tau\delta\Psi_k)((\delta V_k, \delta\Psi_k), (\delta V_k, \delta\Psi_k))(1 - \tau) d\tau, \quad (8.32)$$

where operator \mathbb{L}'' is defined by

$$\mathbb{L}''(U, \Phi)((V', \Psi'), (V'', \Psi'')) := \left. \frac{d}{d\tau} \mathbb{L}'(U + \tau V'', \Phi + \tau \Psi'')(V', \Psi') \right|_{\tau=0},$$

with operator \mathbb{L}' given in (4.4). We can also obtain a similar expression for \tilde{e}'_k (resp. \hat{e}'_k) defined by (8.11) (resp. (8.18)) in terms of the second derivative operator \mathbb{B}'' (resp. \mathcal{E}'').

To control the quadratic error terms, we need the estimates for operators \mathbb{L}'' , \mathbb{B}'' , and \mathcal{E}'' (see (8.32)). They can be obtained from the explicit forms

of \mathbb{L}'' , \mathbb{B}'' , and \mathcal{E}'' by applying the Moser-type and Sobolev embedding inequalities. Omitting detailed calculations, we find that the explicit forms of operators $\mathcal{E}''(U, \Phi)$ and $\mathbb{B}''(U, \Phi)$ depend on state (U, Φ) , which make the following estimates for \mathcal{E}'' and \mathbb{B}'' different from those obtained in [19, Proposition 5]. This difference is caused by the introduction of new primary unknowns (p, hw_1, hw_2) .

Proposition 8.2 *Let $T > 0$ and $s \in \mathbb{N}$ with $s \geq 3$. Assume that $(\tilde{U}, \tilde{\Phi}) \in H_\gamma^{s+1}(\Omega_T)$, $\tilde{\Phi}|_{x_2=0} \in H_\gamma^{s+1}(\omega_T)$, and $\|(\tilde{U}, \tilde{\Phi})\|_{H_\gamma^3(\Omega_T)} \leq \tilde{K}$ for all $\gamma \geq 1$. Then there exist two positive constants \tilde{K}_0 and C , which are independent of T and γ , such that, if $\tilde{K} + \varepsilon \leq \tilde{K}_0$, $\gamma \geq 1$, $(V_1, \Psi_1), (V_2, \Psi_2) \in H_\gamma^{s+1}(\Omega_T)$, and $(W_1, \psi_1), (W_2, \psi_2) \in H_\gamma^{s+1}(\omega_T) \times H_\gamma^{s+1}(\omega_T)$, then*

$$\begin{aligned} & \|\mathbb{L}''(U^a + \tilde{U}, \Phi^a + \tilde{\Phi})((V_1, \Psi_1), (V_2, \Psi_2))\|_{H_\gamma^s(\Omega_T)} \\ & + \|\mathcal{E}''(\tilde{U}, \tilde{\Phi})((V_1, \Psi_1), (V_2, \Psi_2))\|_{H_\gamma^s(\Omega_T)} \\ & \leq C \left\{ \|(V_1, \Psi_1)\|_{W^{1,\infty}(\Omega_T)} \|(V_2, \Psi_2)\|_{W^{1,\infty}(\Omega_T)} \|(\tilde{U} + \tilde{U}^a, \tilde{\Phi} + \tilde{\Phi}^a)\|_{H_\gamma^{s+1}(\Omega_T)} \right. \\ & \quad \left. + \sum_{i \neq j} \|(V_i, \Psi_i)\|_{H_\gamma^{s+1}(\Omega_T)} \|(V_j, \Psi_j)\|_{W^{1,\infty}(\Omega_T)} \right\}, \\ & \|\mathbb{B}''(U^a + \tilde{U}, \Phi^a + \tilde{\Phi})((W_1, \psi_1), (W_2, \psi_2))\|_{H_\gamma^s(\omega_T)} \\ & \leq C \left\{ \|(W_1, \psi_1)\|_{W^{1,\infty}(\omega_T)} \|(W_2, \psi_2)\|_{W^{1,\infty}(\omega_T)} \|(\tilde{U} + \tilde{U}^a, \tilde{\Phi} + \tilde{\Phi}^a)\|_{H_\gamma^{s+1}(\omega_T)} \right. \\ & \quad \left. + \sum_{i \neq j} \|(W_i, \psi_i)\|_{H_\gamma^{s+1}(\omega_T)} \|(W_j, \psi_j)\|_{W^{1,\infty}(\omega_T)} \right\}. \end{aligned}$$

Using Proposition 8.2, we obtain the following result.

Lemma 8.3 (Estimate of the quadratic error terms) *Let $\alpha \geq 4$. Then there exist $\varepsilon > 0$ sufficiently small and $\theta_0 \geq 1$ sufficiently large such that, for all $k = 0, \dots, n-1$, and all integers $s \in [3, \tilde{\alpha} - 1]$,*

$$\|(e'_k, \hat{e}'_k)\|_{H_\gamma^s(\Omega_T)} + \|\tilde{e}'_k\|_{H_\gamma^s(\omega_T)} \leq C \varepsilon^2 \theta_k^{L_1(s)-1} \Delta_k,$$

where $L_1(s) := \max\{(s+2-\alpha)_+ + 4 - 2\alpha, s+3-2\alpha\}$.

Proof In light of (H_{n-1}) and (8.29), we have

$$\sup_{0 \leq \tau \leq 1} \|(V_k + \tau \delta V_k, \Psi_k + \tau \delta \Psi_k)\|_{H_\gamma^3(\Omega_T)} \leq C \varepsilon.$$

For ε small enough, we can apply Proposition 8.2 and use the Sobolev embedding inequality, (8.28), and (H_{n-1}) to obtain

$$\begin{aligned} & \|e'_k\|_{H_\gamma^s(\Omega_T)} \\ & \leq C \left\{ \varepsilon^2 \theta_k^{4-2\alpha} \Delta_k^2 (\varepsilon + \|(V_k, \Psi_k, \delta V_k, \delta \Psi_k)\|_{H_\gamma^{s+1}(\Omega_T)}) + \varepsilon^2 \theta_k^{s+2-2\alpha} \Delta_k^2 \right\} \end{aligned}$$

for $s \in [3, \tilde{\alpha} - 1]$. If $s + 1 \neq \alpha$, then we obtain from (8.29) and the inequality $2\theta_k \Delta_k \leq 1$ that

$$\|e'_k\|_{H_\gamma^s(\Omega_T)} \leq C\varepsilon^2 \Delta_k^2 (\theta_k^{(s+1-\alpha)_+ + 4 - 2\alpha} + \theta_k^{s+2-2\alpha}) \leq C\varepsilon^2 \theta_k^{L_1(s)-1} \Delta_k.$$

If $s + 1 = \alpha$, then using (8.29) and $\alpha \geq 4$ yields

$$\begin{aligned} \|e'_k\|_{H_\gamma^{\alpha-1}(\Omega_T)} &\leq C\varepsilon^2 \Delta_k^2 \{(\varepsilon + \varepsilon \log \theta_k + \varepsilon \theta_k^{-1} \Delta_k) + \theta_k^{1-\alpha}\} \\ &\leq C\varepsilon^2 \Delta_k^2 \theta_k^{1-\alpha} \leq C\varepsilon^2 \theta_k^{L_1(\alpha-1)-1} \Delta_k. \end{aligned}$$

The estimates for \hat{e}'_k and \tilde{e}'_k are similar and follow by applying Proposition 8.2 and the trace theorem. This completes the proof. \square

Now we estimate the first substitution error terms e''_k , \tilde{e}''_k , and \hat{e}''_k given in (8.10)–(8.11), and (8.18) by rewriting them in terms of \mathbb{L}'' , \mathbb{B}'' , and \mathcal{E}'' . For instance, \tilde{e}''_k can be rewritten as

$$\begin{aligned} \tilde{e}''_k &= \int_0^1 \mathbb{B}''(U^a + S_{\theta_k} V_k + \tau(I - S_{\theta_k}) V_k, \Phi^a + S_{\theta_k} \Psi_k + \tau(I - S_{\theta_k}) \Psi_k) \\ &\quad ((\delta V_k|_{x_2=0}, \delta \psi_k), ((I - S_{\theta_k}) V_k|_{x_2=0}, (I - S_{\theta_k}) \Psi_k|_{x_2=0})) d\tau. \end{aligned} \quad (8.33)$$

Then we have the following lemma.

Lemma 8.4 (Estimate of the first substitution error terms) *Let $\alpha \geq 4$. Then there exist $\varepsilon > 0$ sufficiently small and $\theta_0 \geq 1$ sufficiently large such that, for all $k = 0, \dots, n - 1$, and all integers $s \in [3, \tilde{\alpha} - 2]$,*

$$\|(e''_k, \hat{e}''_k)\|_{H_\gamma^s(\Omega_T)} + \|\tilde{e}''_k\|_{H_\gamma^s(\omega_T)} \leq C\varepsilon^2 \theta_k^{L_2(s)-1} \Delta_k,$$

where $L_2(s) := \max\{(s + 2 - \alpha)_+ + 6 - 2\alpha, s + 5 - 2\alpha\}$.

Proof It follows from (8.30) and (8.31) that

$$\sup_{0 \leq \tau \leq 1} \|(S_{\theta_k} V_k + \tau(I - S_{\theta_k}) V_k, S_{\theta_k} \Psi_k + \tau(I - S_{\theta_k}) \Psi_k)\|_{H_\gamma^3(\Omega_T)} \leq C\varepsilon.$$

For ε sufficiently small, we can apply Proposition 8.2 to estimate \mathbb{B}'' in (8.33). Employ the trace and embedding theorems to obtain

$$\begin{aligned} &\|\tilde{e}''_k\|_{H_\gamma^s(\omega_T)} \\ &\lesssim \|(\delta V_k, \delta \Psi_k)\|_{H_\gamma^3(\Omega_T)} \|((I - S_{\theta_k}) V_k, (I - S_{\theta_k}) \Psi_k)\|_{H_\gamma^3(\Omega_T)} \\ &\quad \times \|(\tilde{U}^a + S_{\theta_k} V_k + \tau(I - S_{\theta_k}) V_k, \tilde{\Phi}^a + S_{\theta_k} \Psi_k + \tau(I - S_{\theta_k}) \Psi_k)\|_{H_\gamma^{s+2}(\Omega_T)} \\ &\quad + \|(\delta V_k, \delta \Psi_k)\|_{H_\gamma^{s+2}(\Omega_T)} \|((I - S_{\theta_k}) V_k, (I - S_{\theta_k}) \Psi_k)\|_{H_\gamma^3(\Omega_T)} \\ &\quad + \|((I - S_{\theta_k}) V_k, (I - S_{\theta_k}) \Psi_k)\|_{H_\gamma^{s+2}(\Omega_T)} \|(\delta V_k, \delta \Psi_k)\|_{H_\gamma^3(\Omega_T)}. \end{aligned}$$

Using estimates (8.28), (H_{n-1}) , and (8.30)–(8.31), we obtain that, for $s+2 \neq \alpha$ and $s+2 \leq \tilde{\alpha}$,

$$\begin{aligned} & \|\tilde{e}_k''\|_{H_\gamma^s(\omega_T)} \\ & \leq C \left\{ \varepsilon^2 \theta_k^{5-2\alpha} \Delta_k (\varepsilon + \varepsilon \theta_k^{(s+2-\alpha)_+}) + \varepsilon^2 \theta_k^{s+4-2\alpha} \Delta_k \right\} \leq C \varepsilon^2 \theta_k^{L_2(s)-1} \Delta_k. \end{aligned}$$

For $s+2 = \alpha$, we obtain

$$\begin{aligned} \|\tilde{e}_k''\|_{H_\gamma^s(\omega_T)} & \leq C \left\{ \varepsilon^2 \theta_k^{5-2\alpha} \Delta_k (\varepsilon + \varepsilon \log \theta_k) + \varepsilon^2 \theta_k^{2-\alpha} \Delta_k \right\} \\ & \leq C \varepsilon^2 \Delta_k (\theta_k^{5-2\alpha} \log \theta_k + \theta_k^{2-\alpha}) \leq C \varepsilon^2 \theta_k^{L_2(\alpha-2)-1} \Delta_k, \end{aligned}$$

owing to $\alpha \geq 4$. The estimate for e_k'' and \hat{e}_k'' can be deduced in the same way. \square

Now we estimate the second substitution error terms e_k''' , \tilde{e}_k''' , and \hat{e}_k''' given in (8.10)–(8.11) and (8.18) by rewriting them in terms of \mathbb{L}'' , \mathbb{B}'' , and \mathcal{E}'' . For instance, \hat{e}_k''' can be rewritten as

$$\hat{e}_k''' = \int_0^1 \mathcal{E}''(V_{k+1/2} + \tau(S_{\theta_k} V_k - V_{k+1/2}), \Psi_{k+1/2})((\delta V_k, \delta \Psi_k), (S_{\theta_k} V_k - V_{k+1/2}, 0)) \, d\tau.$$

Here we have used relation $\Psi_{k+1/2} = S_{\theta_k} \Psi_k$ (cf. (8.7)). Then we have the following result.

Lemma 8.5 (Estimate of the second substitution error terms) *Let $\alpha \geq 4$. Then there exist $\varepsilon > 0$ sufficiently small and $\theta_0 \geq 1$ sufficiently large such that, for all $k = 0, \dots, n-1$, and all integers $s \in [3, \tilde{\alpha}-1]$,*

$$\|(e_k''', \hat{e}_k''')\|_{H_\gamma^s(\Omega_T)} + \|\tilde{e}_k'''\|_{H_\gamma^s(\omega_T)} \leq C \varepsilon^2 \theta_k^{L_3(s)-1} \Delta_k,$$

where $L_3(s) := \max\{(s+2-\alpha)_+ + 8 - 2\alpha, s+6-2\alpha\}$.

Proof Omitting detailed derivation, we can use the inductive assumption (H_{n-1}) , definition (8.7), and the properties of S_θ and \mathcal{R}_T to obtain

$$\|S_{\theta_k} V_k - V_{k+1/2}\|_{H_\gamma^s(\Omega_T)} \leq C \varepsilon \theta_k^{s+1-\alpha} \quad (8.34)$$

for all $k = 0, \dots, n-1$ and all integers $s \in [3, \tilde{\alpha}+3]$. We refer to [19, Proposition 7] for the proof of (8.34). It follows from (8.31) and (8.34) that

$$\|V_{k+1/2}\|_{H_\gamma^s(\Omega_T)} \leq C \varepsilon \theta_k^{(s-\alpha)_++1} \quad \text{for } s \in [3, \tilde{\alpha}+3]. \quad (8.35)$$

Thus, we have

$$\begin{aligned} & \|(\tilde{U}^a + V_{k+1/2} + \tau(S_{\theta_k} V_k - V_{k+1/2}), \tilde{\Phi}^a + \Psi_{k+1/2})\|_{H_\gamma^{s+1}(\Omega_T)} \\ & \leq C \varepsilon \theta_k^{(s+1-\alpha)_++1}. \end{aligned} \quad (8.36)$$

For ε small enough, one may apply Proposition 8.2 and use (H_{n-1}) , (8.34), and (8.36) to deduce

$$\|\hat{e}_k'''\|_{H_\gamma^s(\Omega_T)} \leq C \left\{ \varepsilon^2 \theta_k^{6-2\alpha} \Delta_k \varepsilon \theta_k^{(s+1-\alpha)_++1} + \varepsilon^2 \theta_k^{s+4-2\alpha} \Delta_k \right\} \leq C \varepsilon^2 \theta_k^{L_3(s)-1} \Delta_k.$$

The estimate for e_k''' and \tilde{e}_k''' can be deduced in a similar way by using the trace theorem. \square

We now estimate the last error term (8.12):

$$D_{k+1/2} \delta \Psi_k = \frac{\delta \Psi_k}{\partial_2(\Phi^a + \Psi_{k+1/2})} R_k,$$

where $R_k := \partial_2 \mathbb{L}(U^a + V_{k+1/2}, \Phi^a + \Psi_{k+1/2})$. This error term results from the introduction of the good unknown in decomposition (8.10). Note from (8.7), (8.28), and (8.31) that

$$|\partial_2(\Phi^a + \Psi_{k+1/2})| = |\partial_2 \bar{\Phi} + \partial_2(\tilde{\Phi}^a + \Psi_{k+1/2})| \geq \frac{1}{2},$$

provided that ε is small enough. Then we have the following estimate.

Lemma 8.6 *Let $\alpha \geq 5$ and $\tilde{\alpha} \geq \alpha + 2$. Then there exist $\varepsilon > 0$ sufficiently small and $\theta_0 \geq 1$ sufficiently large such that, for all $k = 0, \dots, n-1$, and for all integers $s \in [3, \tilde{\alpha} - 2]$, we have*

$$\|D_{k+1/2} \delta \Psi_k\|_{H_\gamma^s(\Omega_T)} \leq C \varepsilon^2 \theta_k^{L(s)-1} \Delta_k, \quad (8.37)$$

where $L(s) := \max\{(s+2-\alpha)_+ + 8 - 2\alpha, (s+1-\alpha)_+ + 9 - 2\alpha, s+6-2\alpha\}$.

Proof The proof follows from the arguments as in [1, 19]. Let $\Omega_T^+ := (0, T) \times \mathbb{R}_+^2$. Since $\delta \Psi_k$ vanishes in the past, using the Moser-type inequality, we obtain

$$\begin{aligned} \|D_{k+1/2} \delta \Psi_k\|_{H_\gamma^s(\Omega_T)} &= \|D_{k+1/2} \delta \Psi_k\|_{H_\gamma^s(\Omega_T^+)} \\ &\leq C \left\{ \|\delta \Psi_k\|_{L^\infty(\Omega_T^+)} (\|R_k\|_{H_\gamma^s(\Omega_T^+)} + \|R_k\|_{L^\infty(\Omega_T^+)}) \|\tilde{\Phi}^a + \Psi_{k+1/2}\|_{H_\gamma^{s+1}(\Omega_T^+)} \right. \\ &\quad \left. + \|\delta \Psi_k\|_{H_\gamma^s(\Omega_T^+)} \|R_k\|_{L^\infty(\Omega_T^+)} \right\}. \end{aligned} \quad (8.38)$$

To estimate R_k , we introduce the following decomposition for $t > 0$:

$$\begin{aligned} &\mathbb{L}(U^a + V_{k+1/2}, \Phi^a + \Psi_{k+1/2}) - \mathcal{L}(V_k, \Psi_k) + f^a \\ &= \mathbb{L}(U^a + V_{k+1/2}, \Phi^a + \Psi_{k+1/2}) - \mathbb{L}(U^a + V_k, \Phi^a + \Psi_k) \\ &= \int_0^1 \mathbb{L}'(U^a + V_k + \tau(V_{k+1/2} - V_k), \\ &\quad \Phi^a + \Psi_k + \tau(\Psi_{k+1/2} - \Psi_k))(V_{k+1/2} - V_k, \Psi_{k+1/2} - \Psi_k) \, d\tau. \end{aligned} \quad (8.39)$$

If $s \leq \tilde{\alpha} - 3$, the inductive assumption (H_{n-1}) implies

$$\|\mathcal{L}(V_k, \Psi_k) - f^a\|_{H_\gamma^{s+1}(\Omega_T)} \leq 2\varepsilon \theta_k^{s-\alpha}. \quad (8.40)$$

Since we can obtain an estimate for \mathbb{L}' similar to that for \mathbb{L}'' (see Proposition 8.2), using Lemma 8.2 and (8.34) leads to

$$\begin{aligned} & \|\mathbb{L}(U^a + V_{k+1/2}, \Phi^a + \Psi_{k+1/2}) - \mathbb{L}(U^a + V_k, \Phi^a + \Psi_k)\|_{H_\gamma^{s+1}(\Omega_T)} \\ & \leq C\varepsilon(\theta_k^{s+3-\alpha} + \theta_k^{(s+2-\alpha)_+ + 5-\alpha}). \end{aligned} \quad (8.41)$$

Plugging (8.40)–(8.41) into (8.39) yields

$$\|R_k\|_{H_\gamma^s(\Omega_T)} \leq C\varepsilon(\theta_k^{s+3-\alpha} + \theta_k^{(s+2-\alpha)_+ + 5-\alpha}) \quad \text{for } s \in [3, \tilde{\alpha} - 3]. \quad (8.42)$$

If $s = \tilde{\alpha} - 2 \geq \alpha$, then we use (8.31) and (8.35) to obtain

$$\begin{aligned} \|R_k\|_{H_\gamma^s(\Omega_T)} & \leq \|\mathbb{L}(U^a + V_{k+1/2}, \Phi^a + \Psi_{k+1/2})\|_{H_\gamma^{s+1}(\Omega_T)} \\ & \leq C\|(\tilde{U}^a + V_{k+1/2}, \tilde{\Phi}^a + \Psi_{k+1/2})\|_{H_\gamma^{s+2}(\Omega_T)} \\ & \leq C\varepsilon\theta_k^{s+3-\alpha}. \end{aligned}$$

Thus, we obtain estimate (8.42) for $s \in [3, \tilde{\alpha} - 2]$. Thanks to (H_{n-1}) and (8.42), we utilize the embedding inequality to find

$$\|\delta\Psi_k\|_{L^\infty(\Omega_T)} \leq C\varepsilon\theta_k^{2-\alpha}\Delta_k, \quad \|R_k\|_{L^\infty(\Omega_T)} \leq C\varepsilon\theta_k^{6-\alpha}.$$

Using these bounds and plugging (8.42), (H_{n-1}) , and (8.31) into (8.38) yield (8.37). \square

From Lemmas 8.3–8.6, we can immediately obtain the following estimate for e_k , \tilde{e}_k , and \hat{e}_k defined in (8.13) and (8.19).

Lemma 8.7 *Let $\alpha \geq 5$. Then there exist $\varepsilon > 0$ sufficiently small and $\theta_0 \geq 1$ sufficiently large such that, for all $k = 0, \dots, n-1$, and for all integers $s \in [3, \tilde{\alpha} - 2]$, we have*

$$\|e_k\|_{H_\gamma^s(\Omega_T)} + \|\hat{e}_k\|_{H_\gamma^s(\Omega_T)} + \|\tilde{e}_k\|_{H_\gamma^s(\omega_T)} \leq C\varepsilon^2\theta_k^{L(s)-1}\Delta_k, \quad (8.43)$$

where $L(s)$ is defined in Lemma 8.6.

Lemma 8.7 yields the estimate of the accumulated error terms E_k , \tilde{E}_k , and \hat{E}_k that are defined in (8.14) and (8.19).

Lemma 8.8 *Let $\alpha \geq 7$ and $\tilde{\alpha} = \alpha + 4$. Then there exist $\varepsilon > 0$ sufficiently small and $\theta_0 \geq 1$ sufficiently large such that*

$$\|(E_n, \hat{E}_n)\|_{H_\gamma^{\alpha+2}(\Omega_T)} + \|\tilde{E}_n\|_{H_\gamma^{\alpha+2}(\omega_T)} \leq C\varepsilon^2\theta_n. \quad (8.44)$$

Proof Notice that $L(\alpha + 2) \leq 1$ if $\alpha \geq 7$. From (8.43), we have

$$\begin{aligned} & \|(E_n, \hat{E}_n)\|_{H_\gamma^{\alpha+2}(\Omega_T)} + \|\tilde{E}_n\|_{H_\gamma^{\alpha+2}(\omega_T)} \\ & \leq \sum_{k=0}^{n-1} \{ \|(e_k, \hat{e}_k)\|_{H_\gamma^{\alpha+2}(\Omega_T)} + \|\tilde{e}_k\|_{H_\gamma^{\alpha+2}(\omega_T)} \} \\ & \leq \sum_{k=0}^{n-1} C\varepsilon^2\Delta_k \leq C\varepsilon^2\theta_n, \end{aligned}$$

provided that $\alpha \geq 7$ and $\alpha + 2 \in [3, \tilde{\alpha} - 2]$. Thus, the minimal possible $\tilde{\alpha}$ is $\alpha + 4$. \square

8.4 Proof of Theorem 2.1

To prove our main result, we first derive the estimates of the source terms f_n , g_n , and G_n^\pm defined in (8.15) and (8.24).

Lemma 8.9 *Let $\alpha \geq 7$ and $\tilde{\alpha} = \alpha + 4$. Then there exist $\varepsilon > 0$ sufficiently small and $\theta_0 \geq 1$ sufficiently large such that, for all integers $s \in [3, \tilde{\alpha} + 1]$,*

$$\|f_n\|_{H_\gamma^s(\Omega_T)} \leq C \Delta_n \{ \theta_n^{s-\alpha-2} (\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)} + \varepsilon^2) + \varepsilon^2 \theta_n^{L(s)-1} \}, \quad (8.45)$$

$$\|g_n\|_{H_\gamma^s(\omega_T)} \leq C \varepsilon^2 \Delta_n (\theta_n^{s-\alpha-2} + \theta_n^{L(s)-1}), \quad (8.46)$$

and for all integers $s \in [3, \tilde{\alpha}]$,

$$\|G_n^\pm\|_{H_\gamma^s(\Omega_T)} \leq C \varepsilon^2 \Delta_n (\theta_n^{s-\alpha-2} + \theta_n^{L(s)-1}). \quad (8.47)$$

Proof It follows from (8.15) that

$$f_n = (S_{\theta_n} - S_{\theta_{n-1}})f^a - (S_{\theta_n} - S_{\theta_{n-1}})E_{n-1} - S_{\theta_n}e_{n-1}.$$

Using (8.1a), (8.1c), (8.43), and (8.44), we obtain the estimates:

$$\begin{aligned} \|(S_{\theta_n} - S_{\theta_{n-1}})f^a\|_{H_\gamma^s(\Omega_T)} &\leq C \theta_{n-1}^{s-\alpha-2} \|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)} \Delta_{n-1}, \\ \|(S_{\theta_n} - S_{\theta_{n-1}})E_{n-1}\|_{H_\gamma^s(\Omega_T)} &\leq C \theta_{n-1}^{s-\alpha-3} \|E_{n-1}\|_{H_\gamma^{\alpha+2}(\Omega_T)} \Delta_{n-1} \\ &\leq C \varepsilon^2 \theta_{n-1}^{s-\alpha-2} \Delta_{n-1}, \\ \|S_{\theta_n}e_{n-1}\|_{H_\gamma^s(\Omega_T)} &\leq C \varepsilon^2 \theta_{n-1}^{L(s)-1} \Delta_{n-1}. \end{aligned}$$

Combining the above estimates with the inequalities: $\theta_{n-1} \leq \theta_n \leq \sqrt{2}\theta_{n-1}$ and $\Delta_{n-1} \leq 3\Delta_n$, we derive (8.45). Similarly, we obtain (8.46). To prove (8.47), we use (8.24) to find

$$G_n^+ = (S_{\theta_n} - S_{\theta_{n-1}})(\mathcal{R}_T \tilde{E}_{n-1,2} - \hat{E}_{n-1}^+) + S_{\theta_n}(\mathcal{R}_T \tilde{e}_{n-1,2} - \hat{e}_{n-1}^+).$$

Then we obtain the estimate for G_n^+ by using (8.43)–(8.44) as above. The estimate of G_n^- is the same. \square

We are going to obtain the estimate of the solution to problem (8.5) by employing the tame estimate (6.2).

Lemma 8.10 *Let $\alpha \geq 7$. If $\varepsilon > 0$ and $\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}/\varepsilon$ are sufficiently small, and if $\theta_0 \geq 1$ is sufficiently large, then, for all integers $s \in [3, \tilde{\alpha}]$,*

$$\|(\delta V_n, \delta \Psi_n)\|_{H_\gamma^s(\Omega_T)} + \|\delta \psi_n\|_{H_\gamma^{s+1}(\omega_T)} \leq \varepsilon \theta_n^{s-\alpha-1} \Delta_n. \quad (8.48)$$

Proof Let us consider problem (8.5), which can be solved, since $U^a + V_{n+1/2}$ and $\Phi^a + \Psi_{n+1/2}$ satisfy the required constraints (8.9). Constraint (4.1) can be obtained by truncating the coefficients, $U^a + V_{n+1/2}$ and $\Phi^a + \Psi_{n+1/2}$, by a suitable cut-off function, while (4.3b) can be obtained by taking $\varepsilon > 0$ small enough. We can consider the coefficients with a fixed compact support. In order

to apply Theorem 6.1, we obtain (6.1), by using the classical trace estimate, (8.28), (8.31), (8.34), and $\alpha \geq 7$. Thus, we can employ the tame estimate (6.2) to obtain

$$\begin{aligned} & \|\delta \dot{V}_n\|_{H_\gamma^s(\Omega_T)} + \|\delta \psi_n\|_{H_\gamma^{s+1}(\omega_T)} \\ & \leq C \left\{ (\|f_n\|_{H_\gamma^4(\Omega_T)} + \|g_n\|_{H_\gamma^4(\omega_T)}) (\|\tilde{U}^a + V_{n+1/2}, \tilde{\Phi}^a + \Psi_{n+1/2}\|_{H_\gamma^{s+3}(\Omega_T)} \right. \\ & \quad \left. + \|f_n\|_{H_\gamma^{s+1}(\Omega_T)} + \|g_n\|_{H_\gamma^{s+1}(\omega_T)}) \right\}. \end{aligned} \quad (8.49)$$

The particular case $s = 3$ implies

$$\|\delta \dot{V}_n\|_{H_\gamma^3(\Omega_T)} \leq C (\|f_n\|_{H_\gamma^4(\Omega_T)} + \|g_n\|_{H_\gamma^4(\omega_T)}). \quad (8.50)$$

Given $\delta \psi_n$, we can compute $\delta \Psi_n$ from equations (8.16)–(8.17). Performing the energy estimates for $\delta \Psi_n$, and using Lemma 8.2, (8.34), and the Sobolev embedding theorem, we derive

$$\begin{aligned} & \gamma \|\delta \Psi_n\|_{H_\gamma^s(\Omega_T)} \\ & \leq C \left\{ \|g_n\|_{H_\gamma^s(\omega_T)} + \|G_n\|_{H_\gamma^s(\Omega_T)} + \|\delta \dot{V}_n\|_{H_\gamma^s(\Omega_T)} + \|\delta \dot{V}_n\|_{H_\gamma^3(\Omega_T)} \right. \\ & \quad \left. \times \|\tilde{\Phi}^a + S_{\theta_n} \Psi_n\|_{H_\gamma^{s+1}(\Omega_T)} + \varepsilon \theta_n^{(s+2-\alpha)_+} \|\delta \Psi_n\|_{H_\gamma^3(\Omega_T)} \right\} \end{aligned} \quad (8.51)$$

for all integers $s \in [3, \tilde{\alpha}]$ and ε small enough. For $s = 3$, using (8.50), we have

$$\|\delta \Psi_n\|_{H_\gamma^3(\Omega_T)} \leq C (\|f_n\|_{H_\gamma^4(\Omega_T)} + \|g_n\|_{H_\gamma^4(\omega_T)} + \|G_n\|_{H_\gamma^3(\Omega_T)}). \quad (8.52)$$

In view of (8.6), using estimates (8.49), (8.51)–(8.52), and the Moser-type inequality, we obtain

$$\begin{aligned} & \|(\delta V_n, \delta \Psi_n)\|_{H_\gamma^s(\Omega_T)} + \|\delta \psi_n\|_{H_\gamma^{s+1}(\omega_T)} \\ & \leq C \left\{ \|f_n\|_{H_\gamma^{s+1}(\Omega_T)} + \|g_n\|_{H_\gamma^{s+1}(\omega_T)} + \|G_n\|_{H_\gamma^s(\Omega_T)} \right. \\ & \quad + (\|f_n\|_{H_\gamma^4(\Omega_T)} + \|g_n\|_{H_\gamma^4(\omega_T)} + \|G_n\|_{H_\gamma^3(\Omega_T)}) \\ & \quad \left. \times (\|\tilde{U}^a + V_{n+1/2}, \tilde{\Phi}^a + \Psi_{n+1/2}\|_{H_\gamma^{s+3}(\Omega_T)} + \varepsilon \theta_n^{(s+2-\alpha)_+}) \right\} \end{aligned} \quad (8.53)$$

for all integers $s \in [3, \tilde{\alpha}]$. Using Lemma 8.9, (8.31), and (8.34), we obtain from (8.53) that

$$\begin{aligned} & \|(\delta V_n, \delta \Psi_n)\|_{H_\gamma^s(\Omega_T)} + \|\delta \psi_n\|_{H_\gamma^{s+1}(\omega_T)} \\ & \leq C \Delta_n \left\{ \theta_n^{2-\alpha} (\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)} + \varepsilon^2) + \varepsilon^2 \theta_n^{9-2\alpha} \right\} (\varepsilon \theta_n^{(s+3-\alpha)_+} + \varepsilon \theta_n^{s+4-\alpha}) \\ & \quad + C \Delta_n \left\{ \theta_n^{s-\alpha-1} (\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)} + \varepsilon^2) + \varepsilon^2 \theta_n^{L(s+1)-1} \right\}. \end{aligned} \quad (8.54)$$

Exactly as in [19], we can obtain the following inequalities:

$$\begin{cases} L(s+1) \leq s - \alpha, \\ (s+3-\alpha)_+ + 2 - \alpha \leq s - \alpha - 1, \\ (s+3-\alpha)_+ + 9 - 2\alpha \leq s - \alpha - 1, \\ s+6-2\alpha \leq s - \alpha - 1, \\ s+13-3\alpha \leq s - \alpha - 1, \end{cases}$$

for $\alpha \geq 7$ and $s \in [3, \tilde{\alpha}]$. Thus, (8.54) yields

$$\|(\delta V_n, \delta \Psi_n)\|_{H_\gamma^s(\Omega_T)} + \|\delta \psi_n\|_{H_\gamma^{s+1}(\omega_T)} \leq C(\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)} + \varepsilon^2)\theta_n^{s-\alpha-1}\Delta_n,$$

and (8.48) follows by taking $\varepsilon + \|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}/\varepsilon$ small enough. \square

Estimate (8.48) is inequality (a) of (H_n) . We now prove the other inequalities in (H_n) .

Lemma 8.11 *Let $\alpha \geq 7$. If $\varepsilon > 0$ and $\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}/\varepsilon$ are sufficiently small, and if $\theta_0 \geq 1$ is sufficiently large, then, for all integers $s \in [3, \tilde{\alpha} - 2]$,*

$$\|\mathcal{L}(V_n, \Psi_n) - f^a\|_{H_\gamma^s(\Omega_T)} \leq 2\varepsilon\theta_n^{s-\alpha-1}. \quad (8.55)$$

Moreover, for all integers $s \in [4, \alpha]$,

$$\|\mathcal{B}(V_n|_{x_2=0}, \psi_n)\|_{H_\gamma^s(\omega_T)} \leq \varepsilon\theta_n^{s-\alpha-1} \quad (8.56)$$

and

$$\|\mathcal{E}(V_n, \Psi_n)\|_{H_\gamma^3(\Omega_T)} \leq \varepsilon\theta_n^{2-\alpha}. \quad (8.57)$$

Proof From (8.25), we have

$$\begin{aligned} \|\mathcal{L}(V_n, \Psi_n) - f^a\|_{H_\gamma^s(\Omega_T)} &\leq \|(I - S_{\theta_{n-1}})f^a\|_{H_\gamma^s(\Omega_T)} \\ &\quad + \|(S_{\theta_{n-1}} - I)E_{n-1}\|_{H_\gamma^s(\Omega_T)} + \|e_{n-1}\|_{H_\gamma^s(\Omega_T)}. \end{aligned}$$

For $s \in [\alpha + 1, \tilde{\alpha} - 2]$, using (8.1a) and (8.28), we obtain

$$\begin{aligned} \|(I - S_{\theta_{n-1}})f^a\|_{H_\gamma^s(\Omega_T)} &\leq \theta_{n-1}^{s-\alpha-1}(\|f^a\|_{H_\gamma^s(\Omega_T)} + C\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}) \\ &\leq \varepsilon\theta_n^{s-\alpha-1}\left(1 + \frac{C\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}}{\varepsilon}\right), \end{aligned}$$

while, for $s \in [3, \alpha + 1]$, applying (8.1b), we have

$$\|(I - S_{\theta_{n-1}})f^a\|_{H_\gamma^s(\Omega_T)} \leq C\theta_{n-1}^{s-\alpha-1}\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)} \leq C\theta_n^{s-\alpha-1}\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}.$$

Lemma 8.8 and (8.1b) imply

$$\|(I - S_{\theta_{n-1}})E_{n-1}\|_{H_\gamma^s(\Omega_T)} \leq C\theta_{n-1}^{s-\alpha-2}\|E_{n-1}\|_{H_\gamma^{\alpha+2}(\Omega_T)} \leq C\varepsilon^2\theta_n^{s-\alpha-1}$$

for $3 \leq s \leq \alpha + 2 = \tilde{\alpha} - 2$. It follows from (8.43) that

$$\|e_{n-1}\|_{H_\gamma^s(\Omega_T)} \leq C\varepsilon^2\theta_{n-1}^{L(s)-1}\Delta_{n-1} \leq C\varepsilon^2\theta_n^{L(s)-2} \leq C\varepsilon^2\theta_n^{s-\alpha-1}.$$

By virtue of the above estimates, we choose ε and $\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}/\varepsilon$ sufficiently small to obtain (8.55). Similarly, using decompositions (8.26)–(8.27), we can prove estimates (8.56)–(8.57). \square

In view of Lemmas 8.10–8.11, we have obtained (H_n) from (H_{n-1}) , provided that $\alpha \geq 7$, $\tilde{\alpha} = \alpha + 4$, (8.28) holds, $\varepsilon > 0$ and $\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}/\varepsilon$ are sufficiently small, and $\theta_0 \geq 1$ is large enough. Fixing constants $\alpha, \tilde{\alpha}, \varepsilon > 0$ and $\theta_0 \geq 1$, we now prove (H_0) .

Lemma 8.12 *If $\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}/\varepsilon$ is sufficiently small, then (H_0) holds.*

Proof Recall from assumptions (A-1)–(A-3) that $(V_0, \Psi_0, \psi_0, g_0, G_0^\pm) = 0$ and $f_0 = S_{\theta_0} f^a$. Then it follows from (8.7) that $(V_{1/2}, \Psi_{1/2}) = 0$. Thanks to (8.28) and the properties of the approximate solution in Lemma 7.2, we may apply Theorem 6.1 to obtain $(\delta\dot{V}_0, \delta\psi_0)$ as the unique solution of (8.5) for $n = 0$, which satisfies

$$\|\delta\dot{V}_0\|_{H_\gamma^s(\Omega_T)} + \|\delta\psi_0\|_{H_\gamma^{s+1}(\Omega_T)} \leq C \|S_{\theta_0} f^a\|_{H_\gamma^{s+1}(\Omega_T)}.$$

Then we find $\delta\Psi_0^\pm$ from equations (8.16)–(8.17) with $n = 0$. The standard energy estimates yield

$$\|\delta\Psi_0\|_{H_\gamma^s(\Omega_T)} \leq C \|\delta\dot{V}_0\|_{H_\gamma^s(\Omega_T)} \quad \text{for } s \in [3, \tilde{\alpha}],$$

which, combined with (8.6) and (8.28), implies

$$\begin{aligned} & \|(\delta V_0, \delta\Psi_0)\|_{H_\gamma^s(\Omega_T)} + \|\delta\psi_0\|_{H_\gamma^{s+1}(\Omega_T)} \\ & \leq C \|S_{\theta_0} f^a\|_{H_\gamma^{s+1}(\Omega_T)} \leq C \theta_0^{(s-\alpha)^+} \|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}. \end{aligned}$$

If $\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}/\varepsilon$ is suitably small, then we can obtain inequality (a) of (H_0) . The other inequalities of (H_0) can be shown to hold by taking $\|f^a\|_{H_\gamma^{\alpha+1}(\Omega_T)}$ small enough. \square

From (8.12) and Lemmas 8.10–8.11, we derive that (H_n) holds for every $n \in \mathbb{N}$, provided that the parameters are well-chosen and that f^a is sufficiently small. We are now in a position to conclude the proof of Theorem 2.1.

Proof of Theorem 2.1 We consider the initial data (U_0^\pm, φ_0) satisfying all the assumptions of Theorem 2.1. Let $\tilde{\alpha} = \mu - 2$ and $\alpha = \tilde{\alpha} - 4 \geq 7$. Then the initial data U_0^\pm and φ_0 are compatible up to order $\mu = \tilde{\alpha} + 2$. From (7.14) and (7.16), we obtain (8.28) and all the requirements of Lemmas 8.10–8.12, provided that $(\tilde{U}_0^\pm, \varphi_0)$ is sufficiently small in $H^{\mu+1/2}(\mathbb{R}_+^2) \times H^{\mu+1}(\mathbb{R})$ with $\tilde{U}_0^\pm := U_0^\pm - \bar{U}^\pm$. Hence, for small initial data, property (H_n) holds for all integers n . In particular, we have

$$\sum_{k=0}^{\infty} \left(\|(\delta V_k, \delta\Psi_k)\|_{H_\gamma^s(\Omega_T)} + \|\delta\psi_k\|_{H_\gamma^{s+1}(\omega_T)} \right) \leq C \sum_{k=0}^{\infty} \theta_k^{s-\alpha-2} < \infty$$

for $s \in [3, \alpha - 1]$. Thus, sequence (V_k, Ψ_k) converges to some limit (V, Ψ) in $H_\gamma^{\alpha-1}(\Omega_T)$, and sequence ψ_k converges to some limit ψ in $H_\gamma^\alpha(\Omega_T)$. Passing to the limit in (8.55)–(8.56) for $s = \alpha - 1 = \mu - 7$, and in (8.57), we obtain (7.17). Therefore, $(U, \Phi) = (U^a + V, \Phi^a + \Psi)$ is a solution on Ω_T^\pm of the original problem (2.20) and (2.23). This completes the proof. \square

Appendix A Symmetrization of the Relativistic Euler Equations

Under assumption (2.6), Makino–Ukai [33] showed that there exists a strictly convex entropy function for the relativistic Euler equations (2.4), which yields a symmetrizer for (2.4) by following Godunov’s symmetrization procedure in [23]. By contrast, the symmetrizable hyperbolic system (2.18) is deduced by using a purely algebraic symmetrization of the relativistic Euler equations (2.4); see Trakhinin [48] for a different algebraic symmetrization.

In order to derive (2.18), we need to recover another conservation law (that is, conservation of particle number) from equations (2.4). Denoting by N the particle number density and by e the specific internal energy, then

$$\rho = N(1 + \epsilon^2 e). \quad (\text{A.1})$$

The particle number density N was introduced by Taub [46]. For a perfect fluid, N and e are functions of the two thermodynamic variables ρ and S (specific entropy). According to the first law of thermodynamics, the following differential relation holds:

$$TdS = de + p dN^{-1}, \quad (\text{A.2})$$

where T is the absolute temperature. By virtue of (A.1)–(A.2), we have

$$\frac{\partial(\ln N)}{\partial \rho} = \frac{1}{\rho + \epsilon^2 p}, \quad \frac{\partial(\ln N)}{\partial S} = -\epsilon^2 NT. \quad (\text{A.3})$$

In the case of barotropic fluids where pressure p depends solely on ρ , it is natural to introduce the “mathematical” particle number density N as a function of ρ only such that the first relation in (A.3) holds. This motivates us to define $N = N(\rho)$ as (2.11).

Let (ρ, v) be a C^1 -solution to (2.4). It follows from (2.4a) and $h = (\rho + \epsilon^2 p)/N$ that

$$-h\Gamma \{ \partial_t(N\Gamma) + \partial_k(N\Gamma v_k) \} = N\Gamma \{ \Gamma(\partial_t + v_k \partial_k)h + h(\partial_t + v_k \partial_k)\Gamma \} - \epsilon^2 \partial_t p.$$

In view of (2.13) and (2.4b), we obtain

$$\begin{aligned} Nh\Gamma(\partial_t + v_k \partial_k)\Gamma &= Nh\Gamma \epsilon^2 v_j (\partial_t + v_k \partial_k) w_j \\ &= -\epsilon^2 N |w|^2 (\partial_t + v_k \partial_k)h - \epsilon^2 h \Gamma^{-1} |w|^2 \{ \partial_t(N\Gamma) + \partial_k(N\Gamma v_k) \} - \epsilon^2 v_j \partial_j p, \end{aligned}$$

which implies

$$\begin{aligned} h\Gamma^{-1}(\epsilon^2 |w|^2 - \Gamma^2) \{ \partial_t(N\Gamma) + \partial_k(N\Gamma v_k) \} \\ = N(\Gamma^2 - \epsilon^2 |w|^2)(\partial_t + v_k \partial_k)h - \epsilon^2 (\partial_t + v_k \partial_k)p. \end{aligned}$$

Thanks to (2.11) and (2.13), $N(\rho)h'(\rho) = \epsilon^2 p'(\rho)$ and $\Gamma^2 - \epsilon^2 |w|^2 = 1$. Then we obtain the conservation of particle number:

$$\partial_t(N\Gamma) + \partial_k(N\Gamma v_k) = 0. \quad (\text{A.4})$$

Equations (2.14) then follow from (2.4b) and (A.4). Using the identities: $\Gamma_t = \epsilon^2 v \cdot \partial_t w$ and $N'(p) = \frac{N'(\rho)}{p'(\rho)} = \frac{1}{h(\rho)c^2(\rho)}$, we see from (A.4) that

$$\Gamma(\partial_t + v \cdot \nabla_x)p + Nh c^2 (\epsilon^2 v \cdot \partial_t w + \nabla_x \cdot w) = 0.$$

We use the relations: $h'(p) = \epsilon^2/N$ and $w = \Gamma v$ to deduce

$$\begin{aligned} & \Gamma(1 - \epsilon^4 c^2 |v|^2) \partial_t p + \Gamma(1 - \epsilon^2 c^2) v \cdot \nabla_x p \\ & + N c^2 (\epsilon^2 v \cdot \partial_t (hw) + \nabla_x \cdot (hw)) = 0. \end{aligned} \quad (\text{A.5})$$

Set $U := (p, hw_1, hw_2)^\top$. Then equations (2.14) and (A.5) can be written as

$$B_0(U) \partial_t U + B_1(U) \partial_1 U + B_2(U) \partial_2 U = 0, \quad (\text{A.6})$$

where the coefficient matrices are given by

$$B_0(U) := \begin{pmatrix} \Gamma(1 - \epsilon^4 c^2 |v|^2) & \epsilon^2 c^2 N v^\top \\ 0 & \Gamma I_2 \end{pmatrix}, \quad (\text{A.7})$$

$$B_j(U) := \begin{pmatrix} \Gamma v_j (1 - \epsilon^2 c^2) & N c^2 e_j^\top \\ N^{-1} e_j & \Gamma v_j I_2 \end{pmatrix}, \quad j = 1, 2, \quad (\text{A.8})$$

Here we have set $e_j := (\delta_{1j}, \delta_{2j})^\top$ and $I_2 := (\delta_{ij})_{2 \times 2}$ with δ_{ij} being the Kronecker symbol. Let us define

$$S_1(U) := \begin{pmatrix} 1 & \epsilon^2 N c^2 \Gamma^{-1} v^\top \\ 0 & I_2 - \epsilon^2 v \otimes v \end{pmatrix}. \quad (\text{A.9})$$

Multiplying (A.6) by $S_1(U)$ and using the identity: $\Gamma^2 - \epsilon^2 c^2 \Gamma^2 + \epsilon^2 c^2 = \Gamma^2(1 - \epsilon^4 c^2 |v|^2)$, we obtain system (2.18). Conversely, we can also deduce system (2.4) from (2.18) so that we derive the equivalence of these two systems in the region where the solutions are in C^1 .

It remains to show that system (2.18) is symmetrizable hyperbolic in region $\{\rho_* < \rho < \rho^*, |v| < \epsilon^{-1}\}$. Let us set the *Friedrichs symmetrizer*:

$$S_2(U) := \begin{pmatrix} 1 & -2\epsilon^2 N c^2 \Gamma v^\top \\ 0 & N^2 c^2 I_2 \end{pmatrix}. \quad (\text{A.10})$$

After straightforward calculations, we derive that all matrices $S_2(U)A_j(U)$ are symmetric, and the eigenvalues of $S_2(U)A_0(U)$ are

$$\lambda_1 = \Gamma(1 - \epsilon^4 c^2 |v|^2), \quad \lambda_2 = \Gamma N^2 c^2, \quad \lambda_3 = \Gamma N^2 c^2 (1 - \epsilon^2 |v|^2).$$

Assumption (2.6) yields that λ_1 , λ_2 , and λ_3 are all positive. Consequently, $S_2(U)A_0(U)$ is positive definite and system (2.18) is symmetrizable hyperbolic.

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References

1. Alinhac, S.: Existence d'ondes de raréfaction pour des systèmes quasi-linéaires hyperboliques multidimensionnels. *Comm. Partial Diff. Eqs.* **14**, 173–230 (1989). <http://dx.doi.org/10.1080/03605308908820595>
2. Anile, A.M.: *Relativistic Fluids and Magneto-fluids, With Application in Astrophysics and Plasma Physics*. Cambridge University Press, Cambridge (1990). <http://dx.doi.org/10.1017/CBO9780511564130>
3. Artola, M., Majda, A.J.: Nonlinear development of instabilities in supersonic vortex sheets: I. The basic kink modes. *Phys. D* **28**, 253–281 (1987). [http://dx.doi.org/10.1016/0167-2789\(87\)90019-4](http://dx.doi.org/10.1016/0167-2789(87)90019-4)
4. Benzoni-Gavage, S., Serre, D.: *Multidimensional Hyperbolic Partial Differential Equations: First-Order Systems and Applications*. Oxford University Press, Oxford (2007). <http://dx.doi.org/10.1093/acprof:oso/9780199211234.001.0001>
5. Catania, D., D'Abicco, M., Secchi, P.: Weak stability of the plasma–vacuum interface problem. *J. Differential Equations* **261**, 3169–3219 (2016). <http://dx.doi.org/10.1016/j.jde.2016.05.023>
6. Chazarain, J., Piriou, A.: *Introduction to the Theory of Linear Partial Differential Equations*. North-Holland Publishing Co., New York (1982)
7. Chen, G.-Q., Feldman, M.: *The Mathematics of Shock Reflection-Diffraction and von Neumann's Conjectures*. Princeton University Press, Princeton, NJ (2018)
8. Chen, G.-Q., Wang, Y.-G.: Existence and stability of compressible current-vortex sheets in three-dimensional magnetohydrodynamics. *Arch. Ration. Mech. Anal.* **187**, 369–408 (2008). <http://dx.doi.org/10.1007/s00205-007-0070-8>
9. Chen, G.-Q., Wang, Y.-G.: Characteristic discontinuities and free boundary problems for hyperbolic conservation laws. In: Holden, H., Karlsen, K.H. (eds.) *Nonlinear Partial Differential Equations*, pp. 53–81. Springer, Heidelberg, (2012). <http://dx.doi.org/10.1007/978-3-642-25361-4>
10. Chen, R.M., Hu, J., Wang, D.: Linear stability of compressible vortex sheets in two-dimensional elastodynamics. *Adv. Math.* **311**, 18–60 (2017). <http://dx.doi.org/10.1016/j.aim.2017.02.014>

11. Chen, R.M., Hu, J., Wang, D.: Linear stability of compressible vortex sheets in 2d elastodynamics: Variable coefficients. *arxiv:1804.07850* (2018). <https://arxiv.org/abs/1804.07850>
12. Christodoulou, D.: *The Formation of Shocks in 3-Dimensional Fluids*. European Mathematical Society (EMS), Zürich (2007). <https://doi.org/10.4171/031>
13. Christodoulou, D.: *The Shock Development Problem*. European Mathematical Society (EMS), Zürich (2018). <https://arxiv.org/abs/1705.00828>
14. Choquet-Bruhat, Y.: *General Relativity and the Einstein Equations*. Oxford University Press, Oxford (2009). <http://dx.doi.org/10.1093/acprof:oso/9780199230723.001.0001>
15. Coulombel, J.-F.: Weakly stable multidimensional shocks. *Ann. Inst. H. Poincaré Anal. Non Linéaire* **21**, 401–443 (2004). <http://dx.doi.org/10.1016/j.anihpc.2003.04.001>
16. Coulombel, J.-F.: Well-posedness of hyperbolic initial boundary value problems. *J. Math. Pures Appl. (9)* **84**, 786–818 (2005). <http://dx.doi.org/10.1016/j.matpur.2004.10.005>
17. Coulombel, J.-F., Secchi, P.: The stability of compressible vortex sheets in two space dimensions. *Indiana Univ. Math. J.* **53**, 941–1012 (2004). <http://dx.doi.org/10.1512/iumj.2004.53.2526>
18. Coulombel, J.-F., Secchi, P.: On the transition to instability for compressible vortex sheets. *Proc. Roy. Soc. Edinburgh Sect. A* **134**, 885–892 (2004). <http://dx.doi.org/10.1017/S0308210500003528>
19. Coulombel, J.-F., Secchi, P.: Nonlinear compressible vortex sheets in two space dimensions. *Ann. Sci. Éc. Norm. Supér. (4)* **41**, 85–139 (2008). <https://eudml.org/doc/272142>
20. Coulombel, J.-F., Secchi, P.: Uniqueness of 2-d compressible vortex sheets. *Commun. Pure Appl. Anal.* **8**, 1439–1450 (2009). <http://dx.doi.org/10.3934/cpaa.2009.8.1439>
21. Fejer, J.A., Miles, J.W.: On the stability of a plane vortex sheet with respect to three-dimensional disturbances. *J. Fluid Mech.* **15**, 335–336 (1963). <https://doi.org/10.1017/S002211206300029X>
22. Francheteau, J., Métivier, G.: Existence de chocs faibles pour des systèmes quasi-linéaires hyperboliques multidimensionnels. *Astérisque* **268** (2000)
23. Godunov, S. K.: An interesting class of quasi-linear systems. *Dokl. Akad. Nauk SSSR* **139**, 521–523 (1961).
24. Hörmander, L.: The boundary problems of physical geodesy. *Arch. Ration. Mech. Anal.* **62**, 1–52 (1976). <http://dx.doi.org/10.1007/BF00251855>
25. Huang, F., Wang, D., Yuan, D.: Nonlinear stability and existence of vortex sheets for inviscid liquid-gas two-phase flow. *arxiv:1808.05905* (2018). <https://arxiv.org/abs/1808.05905>
26. Jang, J., LeFloch, P.G., Masmoudi, N.: Lagrangian formulation and a priori estimates for relativistic fluid flows with vacuum. *J. Diff. Eqs.* **260**, 5481–5509 (2016). <http://dx.doi.org/10.1016/j.jde.2015.12.004>
27. Kreiss, H.-O.: Initial boundary value problems for hyperbolic systems. *Comm. Pure Appl. Math.* **23**, 277–298 (1970).

- <http://dx.doi.org/10.1002/cpa.3160230304>
28. Lax, P.D.: Hyperbolic systems of conservation laws, II. *Comm. Pure Appl. Math.* **10**, 537–566 (1957). <http://dx.doi.org/10.1002/cpa.3160100406>
 29. Lax, P.D., Phillips, R.S.: Local boundary conditions for dissipative symmetric linear differential operators. *Comm. Pure Appl. Math.* **13**, 427–455 (1960). <http://dx.doi.org/10.1002/cpa.3160130307>
 30. Lichnerowicz, A.: *Relativistic Hydrodynamics and Magnetohydrodynamics*. WA Benjamin, New York (1967)
 31. Lions, J.-L., Magenes, E.: *Non-homogeneous Boundary Value Problems and Applications, Vol. II*. Springer-Verlag, New York-Heidelberg (1972)
 32. Majda, A., Osher, S.: Initial-boundary value problems for hyperbolic equations with uniformly characteristic boundary. *Comm. Pure Appl. Math.* **28**, 607–675 (1975). <http://dx.doi.org/10.1002/cpa.3160280504>
 33. Makino, T., Ukai, S.: Local smooth solutions of the relativistic Euler equation, II. *Kodai Math. J.* **18**, 365–375 (1995). <https://projecteuclid.org/euclid.kmj/1138043432>
 34. Métivier, G.: Stability of multidimensional shocks. In: Freistühler, H., Szepessy, A. (eds.) *Advances in the Theory of Shock Waves*, pp. 25–103. Birkhäuser Boston, Boston, MA, 2001. <http://dx.doi.org/10.1007/978-1-4612-0193-9>
 35. Métivier, G., Zumbrun, K.: Large viscous boundary layers for noncharacteristic nonlinear hyperbolic problems. *Mem. Amer. Math. Soc.* **175** (826) (2005). <http://dx.doi.org/10.1090/memo/0826>
 36. Miles, J.W.: On the disturbed motion of a plane vortex sheet. *J. Fluid Mech.* **4**, 538–552 (1958). <https://doi.org/10.1017/S0022112058000653>
 37. Mishkov, R.L.: Generalization of the formula of Faa di Bruno for a composite function with a vector argument. *Int. J. Math. Math. Sci.* **24**, 481–491 (2000). <http://dx.doi.org/10.1155/S0161171200002970>
 38. Morando, A., Trebeschi, P.: Two-dimensional vortex sheets for the non-isentropic Euler equations: linear stability. *J. Hyper. Diff. Eqs.* **5**, 487–518 (2008). <http://dx.doi.org/10.1142/S021989160800157X>
 39. Morando, A., Trakhinin, Y., Trebeschi, P.: Local existence of MHD contact discontinuities. *Arch. Ration. Mech. Anal.* **228**, 691–742 (2018). <https://doi.org/10.1007/s00205-017-1203-3>
 40. Pan, R., Smoller, J.A.: Blowup of smooth solutions for relativistic Euler equations. *Comm. Math. Phys.* **262**, 729–755 (2006). <http://dx.doi.org/10.1007/s00220-005-1464-9>
 41. Rauch, J.B., Massey, III, F.J.: Differentiability of solutions to hyperbolic initial-boundary value problems. *Trans. Amer. Math. Soc.* **189**, 303–318 (1974). <https://doi.org/10.1090/S0002-9947-1974-0340832-0>
 42. Ruan, L., Wang, D., Weng, S., Zhu, C.: Rectilinear vortex sheets of inviscid liquid-gas two-phase flow: linear stability. *Commun. Math. Sci.* **14**, 735–776 (2016). <http://dx.doi.org/10.4310/CMS.2016.v14.n3.a7>
 43. Secchi, P.: Well-posedness of characteristic symmetric hyperbolic systems. *Arch. Ration. Mech. Anal.* **134**, 155–197, (1996).

- <http://dx.doi.org/10.1007/BF00379552>
44. Secchi, P., Trakhinin Y.: Well-posedness of the plasma-vacuum interface problem. *Nonlinearity* **27**, 105–169 (2014).
<http://dx.doi.org/10.1088/0951-7715/27/1/105>
 45. Smoller J., Temple, B.: Global solutions of the relativistic Euler equations. *Comm. Math. Phys.* **156**, 67–99 (1993).
<http://projecteuclid.org/euclid.cmp/1104253517>
 46. Taub, A.H.: On circulation in relativistic hydrodynamics. *Arch. Ration. Mech. Anal.* **3**, 312–324 (1959). <http://dx.doi.org/10.1007/BF00284183>
 47. Trakhinin, Y.: The existence of current-vortex sheets in ideal compressible magnetohydrodynamics. *Arch. Ration. Mech. Anal.* **191**, 245–310 (2009).
<http://dx.doi.org/10.1007/s00205-008-0124-6>
 48. Trakhinin, Y.: Local existence for the free boundary problem for nonrelativistic and relativistic compressible Euler equations with a vacuum boundary condition. *Comm. Pure Appl. Math.* **62**, 1551–1594 (2009).
<http://dx.doi.org/10.1002/cpa.20282>
 49. Wang, Y.-G., Yu, F.: Stabilization effect of magnetic fields on two-dimensional compressible current-vortex sheets. *Arch. Ration. Mech. Anal.* **208**, 341–389 (2013). <http://dx.doi.org/10.1007/s00205-012-0601-9>
 50. Wang, Y.-G., Yu, F.: Stability of contact discontinuities in three-dimensional compressible steady flows. *J. Differential Equations* **255**, 1278–1356 (2013). <http://dx.doi.org/10.1016/j.jde.2013.05.014>
 51. Wang, Y.-G., Yu, F.: Structural stability of supersonic contact discontinuities in three-dimensional compressible steady flows. *SIAM J. Math. Anal.* **47**, 1291–1329 (2015). <http://dx.doi.org/10.1137/140976169>
 52. Wang, Y.-G., Yuan, H.: Weak stability of transonic contact discontinuities in three-dimensional steady non-isentropic compressible Euler flows. *Z. Angew. Math. Phys.* **66**, 341–388 (2015).
<http://dx.doi.org/10.1007/s00033-014-0404-y>